## roceedings



HIGH TEMPERATURE RESISTOF.


Polytechnic Research and Derelopmext Co.

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# Military Components FOR EVERY APPLICATION 

A HUNDRED STOCK UNITS in our catalog B... 30,000 special designs

## POWER <br> COMPONENTS

The scope of military power com ponents produced at UIC ranges from 500 lb . plate transformers to miniaturized 2 oz units. .. hermetically sealed and encapsulated... molded types.

## ENCAPSULATED UNITS

8 years of encapsulation experience assure maximum reliability in this class of UTC material.

## MOLDED UNITS

ITC molded units range from $1 / 30$ oz miniatures to the $100 \mathrm{lb}, 3$ phase unit illustrated.

PULSE
TRANSFORMERS
UTC pulse transformers cover the range from molded siructures weighing a fraction of an ounce to h gh power modulator applications.


## FILTERS

UTC filters, equalizers and discriminators are produced in designs from 1 cycles to 400 mc . Carrier, aircraft, and telemetering types available in standard designs.

## MINIATURIZED COMPONENTS

UIC H-30 series audios are the smail-
est hermetic types made. Class A, B, and H power components of maximum miniaturization are regular production at UTC.


## MAGNETIC AMPLIFIERS

In addition to a stock line of servo motor magnetic amplifiers, UTC manufactures a wide variety to customer specificaticns. Si turable reactors are supplied for frequencies from 1 cycle to 4C mc.

Operate at temperatures to $125^{\circ} \mathrm{C}$ without voltage derating

Hithstand dielectric test of lwice rated voltage

Insulation resistance higher than any other metallized paper capacitor

Self healing dielectric

Here are the finest capacitors which the present state of the art can produce.

In the application of stringent quality controls, Sprague has gone so far as to metallize its own paper ... the only commercial manufacturer to do this. Thus Sprague is the only capacitor manufacturer with complete control over the end product. And in no other type of capacitor does quality in manufacture play so important a part in performance.

## metallized paper capacitors

A complete rangle of ratings and sizes, hermetically sealed with glass-to-metal solderseals in corrosionresistant cases, is available in numerous mounting and terminal styles. Write for Engineering Bullefin 224 on your letterhead.


## SPRAGUE'

## choose from this complete line of

Sprague, on request, will provide you with complete application engineering service for optimum results in the use of pulse transformers.
NOW YOU CAN CIIOOSE from eighteen sandard pule tranformers in tour major consruction seves, all in guantey production at Sprague. The stondard transformers cotered in the talle helow offer a complete range of chatacterintion for computer circuits, boching oncillator circuis, memory array driving circuis, etc.
These hermedically sated unith will meet such stringent mititary specifatations as MII.-T27, and operate at temperature up to $85^{\circ} \mathrm{C}$. Special desigh are availabie for high acceleration and high ambient semperature operation. In addition, the electrical counterparts of eath tramsormer can be obtatined in lower cost housings designed for typical commerdial enviromment requirements.
Complete information on this high-reliability pulse tranformer line is provided in Engineering Bulletin 502A, available on letterhead request to the Technical Literature Section, Sprague Flectric Company, 235 Marhall Street, North Adams, Massachusetes.

Type $15 z$
miniature bathtub pulse
transformer

| $\begin{gathered} \text { Type } \\ \text { No. } \end{gathered}$ | $\begin{aligned} & \text { Turns } \\ & \text { Ratio } \end{aligned}$ | Pulse Width $\mu$ seconds | Rise Time $\mu$ seconds | $\begin{aligned} & \text { Primary } \\ & \text { Inductance } \end{aligned}$ | $\begin{aligned} & \text { Leakage } \\ & \text { Inductance } \end{aligned}$ | $\begin{gathered} \text { Repenition } \\ \text { Rote } \end{gathered}$ | $\begin{gathered} \text { Lood and } \\ \text { Output } \end{gathered}$ | Typical Applications |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 1021 | 5:1 | 0.1 | 0.04 | $200 \mu \mathrm{H}$ | $5 \mu \mathrm{H}$ | 1 to 2 MC | 15 volts 100 ohms |  |
| 1022 | $4: 1$ | 0.07 | 0.03 | $200 \mu \mathrm{H}$ | $20 \mu \mathrm{H}$ | 1 to 2 MC | 20 volts 100 ohms | Used in digital computer circuity for |
| 1023 | 1:1 | 0.07 | 0.03 | $125 \mu \mathrm{H}$ | $12 \mu \mathrm{H}$ | 1 to 2 MC | 20 volts 200 ohms | impeciance matching and inter- |
| 1024 | 3:1 | 0.07 | 0.03 | $160 \mu \mathrm{H}$ | $15 \mu \mathrm{H}$ | 1 to 2 MC | 20 volts 100 ohms | stage coupling. <br> Pulses are of |
| 1026 | 4:1 | 0.1 | 0.04 | $200 \mu \mathrm{H}$ | $6 \mu \mathrm{H}$ | 1 to 2 MC | 17 volts 100 ohms | sine wave type. |
| 10212 | $1: 1$ | 0.25 | 0.02 | $200 \mu \mathrm{H}$ | $2 \mu \mathrm{H}$ | 12KC | 100 volts | Blocking Oscillator |
| 10213 | 1:1 | 0.33 | 0.07 | $240 \mu \mathrm{H}$ | $2 \mu \mathrm{H}$ | 2KC | 50 volts | Blocking Oscillator |
| 10214 | 7:1:1 | 0.50 | 0.05 | 1.2 mH | $20 \mu \mathrm{H}$ | 1MC | 25 volts | Impedance Matching |
| 1521 | 3:1 | 5.0 | 0.04 | 7.5 mH | $22 \mu \mathrm{H}$ | 10 KC | 10 volts 100 ohms | Impedance Matching and Pulse Inversion |
| 1572 | 2:1 | 0.5 | 0.07 | 6 mH | $15 \mu \mathrm{H}$ |  | 40 volts | Blockiwg Oscillator |
| 1573 | 5:1 | 10.0 | 0.04 | 12 mH | $70 \mu \mathrm{H}$ | 10 KC | 10 volts | Impedance Matching |
| 1574 | 1:1.4 | 6.0 | 0.1 | 16 mH | $15 \mu \mathrm{H}$ | 0.4 KC | 15 volts | Blockirg Oscillator |
| 2021 | $\begin{gathered} 5: 5: 1 \\ \text { Push-Pull } \end{gathered}$ | 1.5 | 0.25 | 4.0 mH | 0.3 MH |  | 5 volts 10 ohms | Memory Core Current Driver |
| 2073 | 6:1 | 1104 | 0.22 | 18 mH | 0.8 MH | $\begin{aligned} & 250 \mathrm{KC} \\ & \text { (max.) } \end{aligned}$ | $\begin{aligned} & 21 \text { volts } \\ & 200 \text { ohms } \end{aligned}$ | Current Driver |
| 2024 | 6:1:1 | 1 to 7 | 0.25 | 55 mH | 0.3 MH | $\begin{aligned} & 50 \mathrm{KC} \\ & \text { (max.) } \end{aligned}$ | $\begin{aligned} & 22 \text { volts } \\ & 400 \text { ohms } \end{aligned}$ | Current Driver and Pulse I iversion |
| 2025 | $\begin{aligned} & 3.3: 3.3: 1 \\ & \text { Push-Pull } \end{aligned}$ | 2.4 | 0.2 | 2.8 mH | 0.2 MH |  | 2.5 volts 6 ohms | Memory C:ore Current Crive ${ }^{*}$ |
| 2026 | 11:1 | 6.0 | 0.2 | 90 mH | 0.2 MH | $\begin{aligned} & 50 \mathrm{KC} \\ & (\text { max. }) \end{aligned}$ | $\begin{aligned} & 10 \text { volts } \\ & 75 \text { ohms } \end{aligned}$ | Current <br> Transformer |
| 4121 | 7:1:1 | 0.50 | 0.05 | 1.2 mH | $20 \mu \mathrm{H}$ | 1 MC | 25 volts | Impedance Matchin |

[^0]This power-type wire wound axial-lead Blue Jacket is hardly larger than a match head but it performs like a giant! It's a rugged vitreous-enamel coated job-and like the entire Blue Jacket family, it is built to withstand severest humidity performance requirements.

Blue Jackets are ideal for dip-soldered sub-assemblies . . for point-to-point wiring . . . for terminal board mounting and processed wiring boards. They're low in
cost, eliminate extra hardware, save time and labor in mounting!

Axial-lead Blue Jackets in 3,5 and 10 watt ratings are available without delay in any quantity you require. $\star \star \star \star$

| SPRAGUE <br> TYPE NO, | WATTAGE <br> RATING | DIMENSIONS <br> L.(inches) | MAXIMUM <br> RESISTANCE |  |
| :---: | :---: | :---: | :---: | :---: |
| $151 E$ | 3 | $17 / 2$ | $13 / 4$ | $10,000 \Omega$ |
| $27 E$ | 5 | $11 / 3$ | $5 / 6$ | $30,000 \Omega$ |
| $28 E$ | 10 | $17 / 6$ | $5 / 6$ | $50,000 \Omega$ |

Standard Resistance Tolerance: $\pm 5 \%$

SPRAGUE ELECTRIC COMPANY • 235 MARSHALL ST. • NORTH ADAMS, MASS.


Above, Bell Laboratories microchemist applies plastic disc in heated clamp to relay contact. Imprint reveals contours of surface and picks up contaminants, if any. Part of portable test set is shown on table. Contacts, shown in small sketches, are of precious metal fused to base metal.

# He's "fingerprinting" 

## a relay contact

Bell Laboratories microchemists have perfected an ingenious new technique for "fingerprinting" relay contacts, the tiny switches on which a dial telephone system critically depends.

Using a portable test set, a chemist makes a plastic print of a contact. On-the-spot examination of the print with a microscope and chemical reagents quickly reveals the effects, if any, of arcing, friction, dust or corrosive vapors. While the chemist studies the print, urgently needed contacts continue in service. Findings point the way to improve relay performance.

This is another example of how Bell Telephone Laboratories research helps to keep your telephone system the world's best.


Preparing dise for microscopic examination. On-the-spot examination may reveal acid, alkali, sulfur, soot or other polluting agents peculiar to an area.


A microscopic look at disc often provides lead to nature of trouble. Unlike actual contact, print can be examined with transmitted light and high magnification.


Here the plastic disc has picked up microscopic lint that insulates contact, stops current. (Picture enlarged 200 times.) Traces of contaminants are identified in microgram quant ties. Inert plastic resists test chemicals that wouid damage contact.

Improving telephone service for America provides careers for creative men in scientific and technical fields

## NEW VARIAN KLYSTRONS ADD SEVEN LEAGUE BOOTS

to microwave transmission...
 tional line-of-sight limits. Designed to exacting Varian quality and performence standards, applications for these versctile klystrons include long rarige communication and cw radar or illuminator service . . . available in the following types and frequency ranges:

| VARIAN TUBE TYPE | FRE QUENCY <br> RANGE (MC) | VARIAN TUBE TYPE | FREQUENCY RANGE (MC) |
| :---: | :---: | :---: | :---: |
| A | 1700-1930 | VA.802 (1 Kwi D | 2450-2700 |
| $\begin{aligned} & V A-3 C 0 \\ & 10 \mathrm{~K} \cdot \mathrm{~N}) \end{aligned}$ | 1935-2160 | VA-803 (1) Kw B | 3700-4200 |
| (10 KW ${ }^{\text {C }}$ | 2160-2400 | VA 804 (1 Kw ) B | 4400.5000 |
| VA-302 | 1700-1930 | VA. 805 (1 Kw) B | 5925-6425 |
| (1 K.w) ${ }^{\text {c }}$ | 1930-2160 | VA-805 (1 Kw) D | 6575.6875 |
| (1) $\bar{C}$ | 2160-2400 | VA-805 (1 KW) D | 6575.6875 |

## Built for long, trouble-free service...

Varian 1 Kw amplifier klystrons offer many advantages for commercial transmitter operation. Rugged, integral-cavity design, air-cooled operation, wide range tuners and conservatively rated, thoriated tungsten buttons provide a life expectancy in excess of 10,000 hours. One power supply design can be used for the entire frequency range . . . no special r.f. equipment is needed. Other outstanding features include:

- Low roise, negligible microphonics.
- High goin - over 50 db ... no intermediate amplifiers required.
- Standard waveguide output - permits direct coupling.
* High efficiency and simplicity of installation.

EXTEND YOUR MICROWAVE HORIZONS . . . Wrise lodoy for complete specifications and technical information on the new Varian 1 Kw and 10 Kw amplifier klystrons . . . dala on the Varion V-42 and other high power klystrons is aiso availoble. Address our Applications Engineering Depariment or contact your nearest Varian representative





- Confinuously Tunable Thru Video VHF and UHF Frequencies, 50KC-950MC Range
- Sweep Widths to 40 MC
- Single Dial Tuning

Used with a standard cathode ray oscilloscope, the Kay Calibrated Mega-Sucep will display the response characteristic of wide band circuits over the frequency range of approximately 50 kc to 950 mc . It features: a calibrated dial indication of the approximate output frequency. The center frequency of the sweeping output voltage may thus be set to an accuracy of abont $10 \%$. The calibrated. Mega-Sureer is the ideal instrument for use in alignment of amplifiers and filters...also as an FM source of wide range for instructional and lab purposes.

## SPECIFICATIONS

Freq. Range: 50 kc to 950 me.
Freq. Sweep: Sawtooth, adjustable to 40 mc .
Repetition rate, 50 to $100 \mathrm{c} / \mathrm{s}$.
RF Output: High, approx. 100 mv max. into open circuit. Low, 5 mv into open circuit.
RF Ouppuf Coniral: Microwave atteruator continuously variable to 26 db .
Output Wavefarm: Less than $5 \%$ harmonic distortion at max. output.
Merer: Provides crystal detector current for peak output.
Regulated Pawer Supply: $105-125$ v., 50 to 60 cps. Power Input, 100 watts.
Send for Catalog 110.A
$\$ 495$ f.c.b. factary

Widest range of the kiay line of sweeping oscillators. Yrovides continuous frequency coverage up through UHF-TV bands50 ke to 1000 me . Widely used in radar systen development and in alignment and testing of TV and F M systems and components, as well as wide band $1 F$ and BF amplifiers and filters. Write for Catalog 100-A. Price, $\$ 465$ f.o.b. factory.


Higher output model calibrated Mega-Sweep, with zero levei baseline. Iligher output facilitates frequency response testing of UIIF converters or tuncrs. Wider sweep width permits multi-channel response viewing. Zero level baseline is convenient means of measuring gain of test circuit.

SPECIFICATIONS
frequency Eange Outpui Impedance Output Valtage

1. $10 \mathrm{mc}-950 \mathrm{mc}$
70 ohms unbalanced
(Into Load)
. 10 . 450 mc 0.3 Vols
5weep Width: Continuously variable to approx. 40 mc max.
Write for Catalog 111.A Price, $\$ \mathbf{S 7 5}$ f.o.b. factory

## kay 112-a calibrated Mega-Suecp

Same as $111-A$, except total frequency range is 800 mc to 1200 mc . Catalog 112-A. Price, $\$ 575$ f.o.b. factory.


HIGH FREQUENCY TRANSISTORS - HERMETICALIY SEALED CASE

| TYPE | Collector |  | Emitter$\mathrm{mA}$ | Extrin. Base Resis. ohms | Base Current Ampl. Factor | Alpha Freq. Cutoff mc . | Max. Junc. Temp ${ }^{\circ} \mathrm{C}$ | Temp. Rise ${ }^{\circ} \mathrm{C} / \mathrm{mW}$ | Coll. Capac. $\mu \mu{ }^{\dagger}$ | Gain |  | Rise time* $\mu$ Secs | Decay time* $\mu$ Secs |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | Volts | Cutoff $\mu \mathrm{A}$ |  |  |  |  |  |  |  | $\begin{gathered} \text { at } \\ 455 \mathrm{kc} \\ \mathrm{db} \end{gathered}$ | $\begin{gathered} \text { at } \\ 2 \mathrm{mc} \\ \mathrm{db} \end{gathered}$ |  |  |
| 2N112 (CN760) | $-6$ | 1 | $-1.0$ | 75 | 40 | 5 | - 85 | 0.62 | 14 | 32 | 18 | 0.05 | 0.06 |
| 2N113 (CR761) | $-6$ | 1 | $-1.0$ | 75 | 45 | 10 | 85 | 0.62 | 14 | 33 | 20 | 0.04 | 0.05 |

Nole: above characteristics are average except where noted
There are more - several times more

Nike, as graceful as the Greek goddess for which she is named, locates, pursues and destroys hostile aircraft. Nike reaches far beyond conventional antiaircraft weapons; outmaneuvers fighters or bombers alike - actually thinks her way to the kill

## From bottom to top:

Nike blasts off
Nike reaches full flight speed in seconds Unerringly, the Nike system's electronic "brain" takes her to the target.


Selection of Western Electric Company as prime contractor for the U. S. Army's Nike guided missile systems was logically based on the necessity for supreme reliability of manufacture and of consequent performance.

Selection of Raytheon Subminiature Tubes by Western Electric was dictated by that same necessity. A number of the subminiature tubes that go into the Nike system's superhuman "brain" are Raytheon Tubes.
No pains were spared, no tests overlooked in securing the very finest, most dependable tubes for the Army's Nike. Think, then, of your own tube applications and their needs whether they be for low microphonics, low power, long life, extreme reliability under severe service conditions or a combination of requirements. Will you be satisfied with anything less than the best? Specify Raytheon Quality Subminiature Tubes.

## RAYTHEON Flat Press Subminiature Tubes ...the tubes with the SEAL of RELIABILITY

The long, flat press glass to metal seal is a Raytheon development that reduces glass strain, button cracking, lead burning, lead corrosion and lead breakage. Its in-line leads permit easier socketing and easier wiring. It is ideal for printed circuitry.





[^1]
a METER TYPE 190-A


Q METER TYPE 260-A


FM-AM SIGNAL GENERATOR TYPE 202-B


SWEEP SIGNAL GENERATOR TYPE 240-A


RX METER TYPE 250-A


RF VOLTAGE STANDARD TYPE 245-A


RX METER: Wide Frequency Band RF Bridge
Q METERS: Low, Medium, High and Very High Frequencies UNIVERTERS: Low, Medium and High Frequency Converters
SIGNAL GENERATORS: Frequency and Amplitude Modulated For Aircraft Navigation, Mobile and TV Receivers. Precision broad band sweeps with markers
SIGNAL GENERATOR CALIBRATORS: RF Voltage Standard in the low microvolt range over a wide frequency range

| a meters |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
| Type | Freq. Range | Q Range | Tuning Capacity Range | Q Aceuracy | Price |
| $\begin{aligned} & 260-A \\ & 190-A \end{aligned}$ | 50 ke to 50 mc 20 mc to 260 me | $\begin{array}{r} 10 \text { to } 625 \\ 5 \text { to } 1200 \end{array}$ | $\begin{gathered} 30-450 \mathrm{mmf} \\ 7.5 \mathrm{to} 100 \mathrm{mmf} \end{gathered}$ | $\begin{aligned} & 5 \% \text { to } 30 \mathrm{me} \\ & 7 \% \text { to } 100 \mathrm{mc} \end{aligned}$ | $\begin{aligned} & \$ 725 . \\ & \$ 625 . \end{aligned}$ |
| FM-AM SIGNAL GENERATORS |  |  |  |  |  |
| Type | Freq. Range | Output Range | Modulation FM AM | Application | Price |
| $\begin{aligned} & 202-B \\ & 202-C \\ & 202-D \end{aligned}$ | $\begin{array}{r} 54-216 \mathrm{mc} \\ 54-216 \mathrm{mc} \\ 175-250 \mathrm{mc} \end{array}$ | 0.1 to 200,000 $\mu v$ <br> 0.1 to $200,000 \mu v$ <br> 0.1 to $200,000 \mu \mathrm{v}$ | $0-240 \mathrm{ke}$ $0-50 \%$ <br> $0-240 \mathrm{kc}$ $0-50 \%$ <br> $0-240 \mathrm{ke}$ $0-100 \%$ | General Mobile Telemetering | $\begin{array}{r} \$ 975 . \\ \$ 1090 . \\ \$ 980 . \end{array}$ |
| SWEEP SIGNAL GENERATOR |  |  |  |  |  |
| Type | Freq. Range | Output Range | Modulation FM AM | Markers | Price |
| 240-A | 4.5 to 120 mc | 1.0 to 300,000 uv | $\begin{aligned} & =1 \% \text { to } \pm 30 \% 30 \% \\ & \text { Center Freq. } \end{aligned}$ | Crystal \& Pip. | \$1375. |
| OMNI-RANGE SIGNAL GENERATOR (Crystal Monitored) |  |  |  |  |  |
| Type | Freq. Range | Output Range | Modulation | Applicetion | Price |
| 211-A | $88-140 \mathrm{mc}$ | 0.1 10 200,000 $\mu \mathrm{v}$ | 0-100\% am Om | ni-Range Revrs. | \$1800. |
| GLIDE SLOPE SIGNAL GENERATOR |  |  |  |  |  |
| Type | Freq. Range | Output Range | Modulation | Application | Price |
| 232-A | 329-335 mc | 1.0 to 200,000 uv | 0-100\% am Glid | de Slope Revrs. | \$1500. |
| WIDE BAND IMPEDANCE MEASURING EQUIPMENT-RX Meter |  |  |  |  |  |
| Type | Freq. Range | R Range | C Range | 1 Range | Price |
| 250-A | 0.5 to 250 mc | 15 to 100,000 ohms | 0-20 $\mu \mathrm{mf}$ ( 0.00 | , Hh - 100 mh | \$1250. |
| RF VOLTAGE StANDARD |  |  |  |  |  |
| Type | Freq. Range | Calibrated Out. | Output Impedance | Application | Price |
| 245-A | 0.1 to $1,000 \mathrm{mc}$ | 0.5, 1.0,2.0 mv | 50 ohms ${ }^{\text {Sig }}$ | Calibrates nal Generators | \$315. |
| UNIVERTERS |  |  |  |  |  |
| Type | Freq. Range | Output Range | Modulation | Accessory to | Price |
| $\begin{aligned} & 207-A \\ & 207-B \\ & 203-B \end{aligned}$ | 0.1 to 55 mc 0.1 to 55 mc 0.11 to 25 mc | $\begin{aligned} & 0.1 \text { to } 100,000 \mu \mathrm{\nu v} \\ & 0.1 \text { to } 100,000 \mu \mathrm{uv} \\ & 1.0 \text { to } 100,000 \mu \mathrm{\mu} \end{aligned}$ | $0-240 \mathrm{ke}$ $0-50 \%$ <br> $0-240 \mathrm{ke}$ $0-50 \%$ <br> 1.5 to 30 me $30 \%$ | $\begin{gathered} 202-B \text { and 202-C } \\ 202-\mathrm{D} \\ 240 \mathrm{~A} \end{gathered}$ | $\begin{aligned} & \$ 345 . \\ & \$ 345 . \\ & \$ 345 . \end{aligned}$ |



BOONTON RADIO
вооптон-n.J.U.S.A. oyporativer



The Tektronix Type 531 Oscilloscope is far ahead in performance characteristics, and is capable of a much wider range of applications than the ordinary general-purpose laboratory oscilloscope.

0THE TYPE 531 EXCELS in vertical-amplifier characteristics - with the Type 531 BPlug -in Preamplifier it offers accurately calibrated sensitivity to $0.05 \mathrm{v} / \mathrm{cm}$ from de to $10 \mathrm{mc}, 0.035-\mu \mathrm{sec}$ risetime... to $0.005 \mathrm{v} / \mathrm{cm}$ from 5 cycles to $9 \mathrm{mc}, 0.04-\mu \mathrm{sec}$ risetime.
(2) THE TYPE 531 EXCELS in sweep characteristics - Miller-runup circuitry generates linear sweeps in the extremely wide range of $0.02 \mu \mathrm{sec} / \mathrm{cm}$ to 12 $\mathrm{sec} / \mathrm{cm}$ ( $6100,0000,000-$-00-1 ratio), with 24 accurately calibrated sweeps from $0.1 \mu \mathrm{sec} / \mathrm{cm}$ to $5 \mathrm{sec} / \mathrm{cm}$. $5 x$ magnifer is accurate on all ranges.
3 THE TYPE 531 EXCELS in triggering facilities -offering amplitude-level selection, automatic rrigkering, and 30 -nce sync in addition to all standard triggering modes.


## (4) THE TYPE 531 EXCELS in writing character.

istics-new Tektronixalecision metallized crt with 10-kvaccelerating potential powides high brighness, improved focus. and excellent linearity. (Recorded writing rate exceeds $175 \mathrm{~cm} / \mu \mathrm{sec}$ ).
5 THE TYPE 531 EXCELS in versatility-Quick change plug-in preamplifiers and inherent oxilloscope capabilities combine to convert the Type 531 to applications normally requiring separate highlyspecialized instruments. Available plug-in units pro. vide for dual-trace ... low-leved differential . . . wide. hand differential...and micro-sensitive applications in addition to wide-band high-gain applications. Current development work promises greatly-extended capabilities through new designs in plug-in units.

> Type 531 Oscilloscope $-\$ 995$
> Type 53B Plug-in Unit - $\$ 125$
> prices f.0.b. Portland (Beaverton), Oregon

## Tektronix, Inc.

P. O. BOX 831 B, PORTLAND 7, OREGON

CYpress 2-2611 . Cable: TEKTRONIX

filtered by bayreon


Atomic Submarine
U.S.S. NAUTILUS

## FROM GUIDED MISSILES

 TO ATOMIC SUBMARINES...
## filtron provides exactly the filter they meed

1 Electrical Test
2 Shielded Laborotory Measurements
3 Screen Roum
Interference Testing
4 Environmentol Testing
5 Attenuation Test
Conscles (per M1L-STD. 220)

From the best equipped Radio Interference Laboratories in the world-staffed with the most experienced Radia Interference Engineers, tomorrow's electric and electronic components and systems are made "Radio Interference Free" today.
FILTRON's exceptional facilities are available for the Radio Interference testing AND filtering of your equipment to meet Military Radio Interference Specifications.
Combining engineering facilifies, application experience, and manufacturing ability, Filfron competently handles RF interference problems from start to finish.
FILTRON's four plants, with complete production facilities-capacitor manufacturing, coil wir.ding, metal fabricating and stamping, tool and die department, assembly divisionare producing more RF interference filters than ever before.
FILTRON-the most dependable name in RF Interference Filters-is the choice of engineers, manufacturers, and mailitary and commercial laboratories the world over.

Send for our free 20 -page RADIO INTERFERENCE FHLTER CATALOG.

## Flexibility in Application

 Versatility in design... packaged analog-digital convertersShaft Position to Digital Converters features reliability, long life, non-ambiguity and speed makes these converters ideal for computers or data handling systems where serial read-out is preferred. Librascope converters transmit information at almost any rate desired up to 1 mc and in some cases above, and may be multiple timeshared, holding extra circuitry to a minimum. All units quickly adjustable, syncro-mounted. Available in Binary, Gray code or Binary decimal code as shown in chart below. Special units may be designed to your order.

Write for catalog information.



# Connect safely with Cannon under all moisture conditions 

Splash Proof • Potting Connectors • Wateright

Protection against moisture is becoming more and maore important. You'll find a wide variety of moisture-proof connectors in the Cannon line ... connectors that solve a wide range of moisture problems ... from those where only minimum protection is required to those where potting of connectors is needed to give maximum safety and performance. Certain Cannon connectors are even designed for complete submersion in such applications as underwater geophysical exploration.

Potting connectors in the AN Series include 12 designs, each in 16 sizes, in both plugs and receptacles, with pin or socket contacts, in the CA group. Potting may be applied, also, to the "K" miniatures, specials and other types. Other moisture-proof connectors in the AN Series include the popular AN-E, OA (Ordnance), special aircraft AF and F types.

For "average" moisture resistance, the XKW and BRS Series are recommended.

Where complete watertight protection is needed, you may select from the heavy-duty W Series in three AN insert sizes. They may be submerged. 2E (Signal Corps type) is available as a moisture sealed power connector..

Our engineers are available to help you with your moisture protection or potting problems. Write TODAY!


CANNON ELECIRIC COMPANY, 3209 Humboldt Street, Los Angeles 31, California.
IFactories in Los Angeles; East Haven; Toronto, Canada; Landon, England. Representatives and distributors in all princ pal cities.

Single sideband＇s spectrum and power economies offer a solution for many of today＇s communication problems．The effects of selective fading and interference due to multipath transmission are minimized through suppression of the carrier and concentrating the r－f power in the intelligence carrying sidebands．Both frequency spectrum space is conserved and the probability of adjacent channel interference reduced by the narrower bandwidth requirements of SSB communications systems．An extensive research and development program at Collins has produced equipment having optimum performance characteristics to realize the full advantages of SSB．Exciters and receivers utilize stabilized master oscillators slaved to a precise frequency standard．Transmitters are easily tuned， efficient and reliable．Channeling facilities incorporate the Collins Mechanical Filter for sideband selection．Typical of Collins equipment available for single sideband communication circuits are：


## 205G－1 TRANSMITTER

The new Collins 205G－1 Communication Transmitter has many outstanding features．Linear operation with low distortion permits multiplex RTTY and／or voice operation with minimum of interference between channels．The transmitter is manually tuned for fixed fre－ quency operation．Only four funed circuits are employed and each covers the 30 to 60 mc frequency range．All sub－units and compo－ nents are accessible from the front of the cabinet．
FREQUENCY RANGE： 30 to 60 me ．
TUNING：Manual over $\pm 2.5 \%$ range．
TYPE OF EMISSION：$A_{1}, A_{3}{ }^{b}$ ，or teleprinter signals．
POWER OUTPUT： 20 kw carrier or peak envelope power．
OUTPUT IMPEDANCE： 52 ohms with up to 2 to 1 SWR．
DRIVE REQUIREMENTS： 0.5 watt at earrier frequency at 52 ohms．
SSB DISTORTION：3rd order distortion products at least 30 db below one tone of a two－tone test signal at 20 kw P．E．P．
HARMONIC OUTPUT：2nd harmonic is at least 35 db down．

## 50P－1 RECEIVER

The 50P．I fixed frequency communication receiver consists of an RF amplifier using high $Q$ circuits，first and second i－f amplifiers and mixers．The 250 ke i－f output feeds accessory equipment for recovering the RTTY and voice signals．Maxi－ mum rejection of adjacent channel interference together with minimum inter－modulation and cross modulation is provided． When used with the Collins 708B．1 Stabilized Master Oscillator and Collins 40K．1 Frequency Standard，the total frequency error is maintained at less than one part in 100 million（ $0.000001 \%$ ） making possible better utilization of spectrum space and at． tainment of better signal－to－noise ratio by allowing bandwidth requirements to be reduced to a minimum．

FREQUENCY RANGE： 20 mc to 50 mc ． OUTPUT FREQUENCY： 250 ke i－f output．
AMBIENT TEMPERATURE RANGE： 0 to 50 degrees $C$ ．
SIZE： $31 / 2^{\prime \prime}$ high， $171 / 4^{\prime \prime}$ wide，and $7^{\prime \prime}$ deep．
MOUNTING：Relay rack．
POWER INPUT：Supplied by external type $426 \mathrm{~B}-1$ power supply － 50 watts．

Complete terminal equipment for generalized data trans－ mission including teletypewriter using synchronous detection techniques is also available to provide a fully integrated system．
Write for additional information．

## COLLINS RADIO COMPANY

# by prition ANI <br> <br> climatic condition 

 <br> <br> climatic condition}



## "Ilesigned for Application"

## Delay Lines and Networks

The James Millen Mfg. Co., Inc. Has been producing continuous delay lines and lump constant delay networks since the origination of the demand for these components in pulse formation and other eirenits requiring time delay. Thre most modern of these is the distrib)uted constant delay line designed to comply with the most stringent electrical and mechanical requirements for military, commercial and laboratory equipment.

Millen distributed constant line is available as bulk line for laboratory use and in cither flexible or metallic hermetically sealed umits adjusted to exact time delay for use in production equipment. Lump comstant delay networks may be preferred for some sperialized applications and can be furnished in open or hermetieally sealed construction. The above illustrates several typical lines of hoth types. Our engineers are available to assist you in your delay line problems.

#  MALDEN, MASSACHUSETTS, U.S.A. 

# Direct, automatic <br>  



## SPECIFICATIONS

Power Range: 5 ranges, front panel selector. Full scale readings of $.1, .3,1,3$ and 10 mw . Also continuous readings from -20 to +10 dbm . ( $0 \mathrm{dbm}=.001$ watt). Power range may be extended with attenuators or directional couplers in microwave system.
External Bolometer: Frequency range depends on bolometer mount. Bolometers can operate at resistance levels of 100 or 200 ohms and can have positive or negative temperature coefficients. Any dc bias current up to 16 ma is available for biasing positive or negative temperature coefficient bolometers. Dc bias current is continuously adjustable and independent of bolometer resistance and power level range
Suitable bolometers are:
Instrument fuses: - $b p$ - G-28A $1 / 100 \mathrm{amp}$ fuse:
Barretters: Sperry 821, Narda N821B or N610B, PRD 610A, 614, 617 or 631 C .
Thermistors: W. E. D166382 and 32A3, V. E. Co. 32A3, 32A5, Narda 333, 334.
Accuracy: $\pm 5 \%$ of full scale reading.
Power: $115 / 230 v \pm 10 \%, 50 / 1,000 \mathrm{cps}, 75$ watts.
Dimensions: Cabinet Mount: $73 / 8^{\prime \prime}$ wide, $111 / 2^{\prime \prime}$ high, $121 / 4^{\prime \prime}$ deep. Rack Mount: $19^{\prime \prime}$ wide, $7^{\prime \prime}$ high, $121 / 2^{\prime \prime}$ deep.
Weight: Net 20 lbs . Shipping 32 lbs . (cabinet mount).
Price: $\$ 250.00$.
Prices f.o.b. fatory. Data swbject to change withourl nolice.

## CW or pulsed power

## Wide frequency range

No calculations

## Assured accuracy

## Operates with wide variety of bolometers

New! -hp- 430C Microwave Power Meter

Here is the newest, finest, most dependable source of instantaneous microwave power readings available today. The new -hp- 430C gives you power readings direct in db or mw and completely eliminates tedious computations or troublesome adjustment during operation. The instrument measures either pulsed or CW power on either waveguide or coaxial systems. Operation is entirely automatic, stability is extremely high, and the meter may be used with a wide variety of bolometer mounts having either positive or negative temperature coefficients. The broad nominal measuring range can be extended to higher powers by means of directional couplers and attenuators.
For measurements of CW or pulsed power, -hp- 430C uses either an instrument fuse, barretter or thermistor as a bolometer element. Operation may be at either 100 or 200 ohms. Power is read direct in milliwatts from 0.02 to 10 mw , or in dbm from -20 to +10 dbm .

HEWLETT-PACKARD COMPANY
33410 PAGE MILL ROAD - PALO ALTO, CALIFORNIA, U. S. A. Cable "HEWPACK"
field representatives in all principal areas

NEWS and NEW PRODUCTS

## Digital-Analog Computer

Wang Laboratories, 37 Hurley St., Cambridge, Mass., annomeses design and production of the "WEDILOG," a new clectronic digital-analog differential computer, with applications in aeronautical design, trajectory problems, process control, dynamic systems analysis, etc. Athough it is a digital computer, "WEDILOG" operates on the principle of functional simulation, using numbers instead of voltages. 1 hesigned for problems in the physical sciences and engineering, this computer will solve linear and nom-linear ordinary and partial differential equations, integral equations, and simultaneous differential and algebraic equations. "WEDILOG" combines the simplicity of prohlem set-up of the 1 D analog computer with the accuracy and resolution of a full scale digital computer. The machine resolution is fise decimal digits. All mumbers are handled as true mumber with sign.


The basic "WEDILOG" computer is made up of a combination of several types of computation units operated from a central control. The problem, itself, determines the mumber and types of anits and how they are "patched" together. I'nitized constriction of these computation units enables easy expansion of the basic computer, as more comprehensive problems arise. Outputs are available as electric typewriter tabulations, multiple channel recordings, on plotting boards, or on punch cards. In addition, all variables are indicated by neon lights, readable at all times.

## Ham Tube Catalog

Tube Div. Radio Corp. of America, 415 S. Sth St., Ilarrisem, N. J., has just lrought ont a completely revised edition of the 4 page folder "Ileadliners for 1 lams.

This folder covers 45 popular RC. "I tams" types-oscillators, amplifiers, frefuency mulipliers, volage regulators, thyratrons, rectifiers, oscillograph types for test equipment, and camera tubes for use in amateur telecasting.
"Ileadliners for llams" comtains a tube

These manufacturers have invited PROCEEDINGS readers to write for literature and further technical information. Please mention your I.R.E. affiliation.
line-up chart for amateur transmitters: operating conditions for class C amplifier and oscillator, modulator, and fremencymultiplier service; single-sidetrand tube data; and latest "Ham" ratings on popmlar receiving-tube types.

Members can obtain copies of "Headliners for Hams" from local RC. 1 Tube 1)istributors, or direct from Commercial Engincering, Tube Division, Radio Corporation of America, Harrison, N., I.

## Oscilloscope Pre-Amplifier

The new IS-61 I oscilloscope preamplifier developed by Volkers \& Schaffer Mfg. Corp., Bow 990, Schenectady, N., Y', is a highly sensitive, dual input, adeling or differential amplifier, having substantially less than $1 \mu V^{\prime}$ RMS noise. It incorporates the latest results in low-moise amplifierresearch. Volkers \& Pedersen have shown (1955 1RE National Convention) that transistors, contrary to general experience and opinion, are inherently less noisy than vacoum tubes, if suitable operating parameters are selected. The resulting new "llushed "ransistor Amplifiers," having kess than $1 \mu V^{\prime}$ RMS moise over a frequency band of 60 ke , provides the input stages of the I'S-61 . . . Wditional amplification is provided by vacuum tubes.


The amplifier has a stage selecor switch which provicles a chovice of either straight vacumm tube amplification (maximum gain 10, input imperlance 100 K, frequency response 2 cps $(0250 \mathrm{kc}$ ) or combined transistor and vacum tube amplification (available maximmm gains 200 and 1000 , input impedance ! K゙, frequency response 2 (p)s to 60 kc ). All frequency responses are given for the 3 db down point. The maximam noise with straight tube operation is $5 \mu \mathrm{~V}$ R MS over the full 250 kc pass-hand and with transistor and tube operation usually much less than $0.5 \mu V^{\prime}$ R MS over the full 250 kc pass-band.

## Power Supply

The new Model lllk-240 regulated power supply developed by Krohn-Hite Instrument Co., 580 Massachusetts Ave., Cambridge 39, Mass, provides up ${ }^{7}$ to $\frac{1}{2}$ ampere of direct current at $0-500$ wolts with 0,001 per cent regulation and less that 100 microvolts of ripple. The stabilization for $\pm 10$ per cent change in line voltage is better than 0.0013 per cent.


The de and tow frequency impedance is less than 0.005 ohms. The ac imperlance is less than 0,05 ohms in series with 0.1 microhenry (4 inches of wire), Transient response is 0.001 millisecond. Typical ten hour drift is 300 ppm plus 20 millivolts.

The altra-high regulation applies over the entire operating range. For line voltages between 105 and 125 volts, the full naximum current can be drawn continuously at any output voltage.

There is an additional 0-1.50 volt, 0-5 ma : engative supply with 0.05 per cent stabilization and less than 2 millivolts of ripple. A $5-1.3$ volt, $12-2.5$ ampere de heater supply with less than 20 millivolts of ripple is included in addition to two independent 6.3 volt ac, 10 ampere heater supplies.

The two $3 \frac{1}{2}$ inch from panet meters are ruggedized and hermetically sealed.

Dimensions are $17 \frac{1}{2}$ wide, 9 high, and $1 . \frac{1}{2}$ inches deep-also availalle for rack mesunting. Price is $\$ 550.00$, f.r.b, factory:

## Nylon Screws \& Nuts

Weckesser Co., 5261 N . Tvondale Are.. Chicago 30), [ll., has awislable from inventory 10 stock items of sorews and nuts of mokded black nyon.


These non-magnetic, non-corrosive, light weight screws and muts come in sizes $\frac{6}{32}, \frac{8}{32}$ and $\frac{10}{3} \frac{0}{2}$.
(Continued on pued 20.4)

## 

ansencts
OUTPUT POWER METERS of unexcelled accuracy and neliatility have many applications


The DAVEN Output Powen Meters cre desi gned to ulignal sure the octual power olivered ber anse of the char: syatem to a givon liod. they are admirobly suitod to acteristics al wictions ticmels
other opplications iicmaly

2. Effects of Loedive Equalisotion Meoswremints: Muificehenal Miser

1. Tromirispion Lime Enaption Less in
2. Measurement mples circuith.
and other complif.ermer Mesiwremmenth.
3. Filser and Trent Measursments.
 The equipment whown on this page is bunise write for well-known stardordn Let our engineering deportment more deloslad dacific problems. help you on specific probiems.

## TYPE OP-962

Characteristics similar to Op-961. except that it can 100 watts. measure up Range: 40 seImpedance Radances belected impedances ohms. tween 2.5 and $2 \%$ over fre-
Impedance Bange ohms to 20,000 ohms. 2.5 mains essentially resistive over frequency range of $\pm 2 \%$. $10,000 \mathrm{cps}$. Accuracy

Power Range: 0.1 milli. watts to 50 watts 0.1 milliof 0.1 milliwatts. in steps
Indicatime
brated from Meter: Cali. watts and 0 io 50 milliZero level: to 17 decibels. Meter Multiplier: Extends indicating reading of the to 1.000 x beter from $0.1 x$ the db. reading value, or to +30 db . in from -10 2 db . db. in sleps of Accuracy $\pm 2$ rage 30 to 10,000 cycles. 0.1 mw to Power Range: 0.1 mw steps. 100 watts in be eitended Range may be by use of below 0.1 amplifier. extemal ampler: CaliIndicating Me watt to 1 brated hrom. 01 wato 10 watt and from -10 to db. Zero Mulliplier: Extends range of meter from 0.01
 to 100 times scale reading.

## MEASURE NOISE AND FIELD INTENSITY FROM 150 KC TO 1000 MCWITH ONE METER!

 Quickly • Accurately • Reliably
(Commercial Equivalent of AN/URM•7)
Enipire Devices Noise and Field Intensity Meter Model NF-105 permits measurements of RF interference and field intensity over the entire frequency range from 150 kilocycles to 1000 megacycles. It is merely necessary to select one of four individual plug-in tuning units, depending on the frequency range desired. Tuning units are readily interchangeable...can be used with all Empire Devices Noise and Field Intensity Melers Model NF-I05 now in the field.

Each of the four separate tuning units employs at least one RF amplifier stage with tuned input. Calibration for noise measurements is easily accomplished by means of the built-in impulse noise calibrator. With this instrument costly repetition of components common to all frequency ranges is eliminated because only the tuners need be changed. The same components... indicating circuits, calibrators, RF attenuators, detectors and audio amplifier. and power supplies... are used at all times.
Noise and Field Intensity Meter Model NF. 105 is accurate and versatile, it may be used for measuring field intensity, RF interference, or as an ultra-sensitive VTVM. A complete line of accessories is available.

Addisional information and hterature upon request
Visit us at 252 Instruments Avenue, at the I.R.E. Show
NEW YORX-019by 9.1240 - SYRACUSE-SYYocuse 2.6253 - PHILADELPHIASHerwood $1-9080$ - BOSTON-WAIHhom 5.1995. WASHINGTON. D. C. -DE Cotur 2.8000- DETROII - BRoodway 3-2900 - CLEVELANO - EVergreen 2.4114 DAYYON-FUIION 8794 . CHICAGO-COIUMBUY 1.1566 - DENVER-MAin 3.0343
 2.8103 - CANADA: MONTREAL-UNiversity 6.5149 - TORONTO-WAInU1 4-1226 halifax 4.6487-EXPORT; NEW YORK-MUrray Hill 2-3760

## EMPIRE DEVICES PRODUCTS CORPORATION

38-15 BELL 8OULEVARD BAYSIDE 61 . NEW YORK
fielo intensity meters - distortion anatyiers - impuise generaiors - coaxial altenuators - crystal mixers

These manufacturers have invited PROCEEDINGS readers to write for literature and further technical information. Please mention your I.R.E. affiliation.

## (Continued from prage 12A)

An advantage of these serews is that they are insulators in themserves, and require no insulating sleeves, bushings or washers.

The nylon is the harolest grade and it may be used to $250^{\circ} \mathrm{F}$. It resists weathering and degrading effect due to ultra-violet light because of the black dye impregnation.

## New Crystal

A new frequency comtrol in it has been introduced ly the James Kiights Co., Sandwich, III. It is mamed the "Thermystal," and represents the integrated packaging of crystal and owen to provide stability and envirommental eontrol within limited price range.


The JK Thermystal offers the following performance data: higher mert factor: vacmum enclosure increases $Q$ of crystal: calibration accuracy; 1 cysle 0.0001 per cent: temperature stability; 30 to 900 kc 0.0001 per cent, 1000 ke to 150 me 0.00005 per cent. Oven temperature varies less than $1^{\circ} \mathrm{C}$ over ambient range of $-55^{\circ} \mathrm{C}$ to $+85^{\circ} \mathrm{C}$, secular stabi.ity; Lees than (0.00t per cent per year. Crystal is specially processed and sealed in glass enclosed vacum: low oven power; 6.3 volts at 1.5 ampere maximum. The thermostat eycles less than 3 times per minute at room temperature. The unit stabilizes in less than 10 minutes when turned on at $-55^{\circ} \mathrm{C}$.
(Continued on page 115, A)
microwave and power tube operations-Waltham 54. $\rightarrow+\operatorname{sow}$


## What's New in Mnemonics?

The news is that the magnetic-core memory has emerged from the computer laboratory and has been in customer use for approximately a year, passing all tests with flying colors. This new development has been pioneered by Remington Rand with the Univac Scientific-the first installation of a commercially available computer that successfully uses magnetic-core storage.

Mnemonics, says Webster, is "the art of improving the efficiency of the memory." And, as far as electronic computers are concerned, Remington Rand has clearly established its leadership in this art.

Illustrated above is a single plane of core storage, each of which holds 4,096 binary bits of information. Planes are wafer thin, aitd a stack capable of
"remembering" 147,456 bits would measure only 13 inches in depth. The speed, economy, and reliability of this magnetic-core memory are now available in the new Univac Scientific Models 1103A and 1103B.

For the latest ioformation about the Remington Rand ragnetic-core memories or about the Umivar Scim tific, write, on your business letterhead, (i)..

## Teonimgrom Thanal

DIVISION OF SPERRY RAND COIIGAAI,ON
 Chatham to fill a specialized need, now number among the most widely used tubes in the industry. For complete information on Chatham tubes - either stock items or types built to your requirements - call or write today.


## Chatham Eiectronics



## for equipment which demands the

finest paper tubulars

Leading manufacturers specify
the industry's finest paper tubular capacifor
... The Sangame
Terechié

For critical applications such as hi-fi equipment, computers and other electronic gear... applications which require exceptionally high insulation resistance and unusual stability at high
temperatures, your best bet is a Sangamo Telechief.
It is the molded paper tubular which, tests by leading manufacturers show, outperforms all other paper tubulars in . . .


Because the Telechief outperforms other paper tubulars in all of these areas, you can be sure that here is a paper tubular which will deliver long, trouble-free capacitor life.

SAREAMO ELECTRIC COMPANY MARION, ILLINOIS


## MAGNETIC TAPE RECORDING

## helps Road-Test Timken Truck Axles



Magnetic tape recordings are now being used to duplicate rugged road-tests at the Timken-Detroit Axle Division of the Rockwell Spring and Axle Company, Detroit, Michigan. A four hour tape cycle is made of actual road surface and driving conditions. . then played back through torque and speed dynamometers repeatediy - until a test axle breaks down.
Result: more realistic and efficient testing - bet'er axles for today's trucks, buses and trailers.

## WHY TIMKEN CHOSE AMPEX

Timken engineers requ'red a recording and playback medium that could give near-perfect reproduction of the original road test phenomena . . . and would playback inciefinitely without introducing errors through wear and speed irregularities. They found that the Ampex F-M recorder best met these exacting requirements. Its extreme stability of tape motion, precise timing and consistent accuracy produced laboratory "road-test" results within $1 \%$ of actual conditions.


## LET AMPEX STUDY YOUR REQUIREMENTS

Ampex manufactures the most complete line of magnetic recorders for complex and sensitive automation, communication and data-handling systems. Why not let Ampex application engineers determine what magnetic tape recording can do for you?

For furthep information, send for our 16-page illustrated bulletin, "Data Recording, Machine Control and Process Regulation." Contact your nearest Ampex representative or write to Depf. G-1897.


Wedg－loc．．．The exclusive wedge leads on these DIS－ CAPS lock securely in place on printed circuit assemblies prior to the soldering operation．There is no possibility of the capacitors becoming loose or falling out and the soldered connection is always uniform．

Available in capacities between 2 MMF and $20,000 \mathrm{MMF}$ ， Wedg－loc DISCAPS can be furnished in temperature com－ pensating，by－pass．and stable capacity types．Suggested hole size is a ． 062 square．


Plug－in ．．．RMC plug－in DISCAPS are designed to sim－ plify production line problems on printed circuits．Leads are No． 20 tinned copper（ .032 diameter）and are available up to $11 / 2^{\prime \prime}$ in length．Plug－in DISCAPS are manufactured in temperature compensating，by－pass，and stable capacity types and include the mechanical and electrical features that have made standard DISCAPS the favorite of leading manufacturers．

Write today on your company letterhead for expert engineering help on any capacitor problem．

RADIO MATERIALS CORPORATION<br>GENERAL OFFICE： 3325 N．California Ave．，Chicago 18，III．

[^2]
## Introducing the Eimac 4X250B Radial-beam power tetrode - Higher Power - Easier Cooling <br> 

 - Longer Life4X250B, a new, superior radial.beam power tetrode by Eimac - originators of the famous 4X150A - is now available. Unilaterally interchangeable with the 4X150A in practically all applications, this amazing new bantam for modulator, oscillator and amplifier application from low frequencies into UHF, offers these advantages:

HIGHER POWER-Electrical advances permit an increased plate dissipation rating of 250 watts, plate voltages to 2000 volts and doubled plate power input capabilities of 500 watts.

EASIER COOLING - Development of the Eimac integral-finned anode makes cooling so easy that only one-third the air-pressure and onehalf the cubic feet of air are required. Forced air is unnecessary during standby periods.


For further defails contact our Technical Services Deparfment.

LONGER LIFE—A newly designed, highly efficient oxide cathode and increased temperature tolerances, coupled with Eimac-developed production and testing techniques enable the 4X250B to meet the most critical standards. New techniques in grid production, high vacuum outgassing and product evaluation are among the features that insure uniform incomparable quality and more hours of top periormance.

The small, rugged, versatile $4 \times 2503$ is now available for existing sockets or sockets of yet-to-bedesigned equipment demanding optimum quality and performance.

TYPICAL OPERATION
(per fube, frequeneies to 175 mc ) 4X250B radial-beam power tetrode

| Class C CW FM Phane | Class C AM Phone | Class AB RF Linear |
| :---: | :---: | :---: |
| D-C Plate Voltage 2000v | 1500v | 2000y |
| D-C Screen Voltage 250r | 250v | 350 y |
| D-C Grid Voltage $\quad 90 y$ | -100v | -60y |
| Zero Sig D-C Plate Current - | - | 50ma |
| D-C Plate Current 250 mo | 200 mo | 250ma* |
| D-C Screen Current 12mo | 10 mo | 5ma* |
| O-C Grid Current 22mo | 23 mo | Omo* |
| Peak RF Grid Voltoge 114y | 125v | $60{ }^{*}$ |
| Driving Power $\quad 2.5 \mathrm{w}$ | 2.9 w | - |
| Plate Power Input 500 w | 300w | 500w* |
| Plate Power Output 400 w | 240w | 325w* |
| *Maximum Signal |  |  |

# EITEL-McGULLOUGH, INC. <br>  <br> The World's Largest Manufacturer of Transmitting Tubes 



## New Mallory Cardboard

## Tubular Capacitors

## ... premium performance at no increase in cost

Never before las quality like this been built into cardboard tubular electrolytics. At no increase in price, this new series developed by Mallory offers you a combination of features unique in this type ol capacitor:

> Minimum size, high ripple current ratings, low RF impedance . . ohtained through use of gemaine fabricated plate anoles.
> Long life, high stability temperature rating up to $75^{\circ} \mathrm{C}$...due to fabricated anode and etehed cathode.
> Low leakage current.
> Low-resistance tab-to-lead wire connections . . . welding ends danger of intermitent or high resistance connection.
> Low moisture loss . . cartridge is foil wrapped; wax impregnated cardooard tulse is sealed with wax at both ends.
> High dielectric strength, exceds U.l. requirements, due to improved low-moisture absorbent separators.
> Rugged, flexible leads. . covercd with plastic insulation rated for $105^{\circ} \mathrm{C}$, have L . L. approval.

The new series comes in single, dual, triple and quad sections, with leads all coming from one end or from opposite ends of the cartridge. A complete choice of voltage and capacity ratings is available.
For technical data, write or call Mallory today. A Mallory capacitor engineer will be glad to consult on your circuit reguirements, to suggest possible cost-cutting simplifications based on Mallory's long experience in all types of applications for electronic components.

Parts distributors in all major cities stock
Mallory standard components for your convenience.

Serving Indusiry with These Products:
Electromechanical-Resistors - Switches - Television Tuners - Vibrators
Electrochemical-Capacitors - Rectifiers - Mercury Batteries Metallurgical-Contacts - Special Metals and Ceramics * Welding Materials


Insides of the rase is this foil-wrapped cartridge. Tabs are uedded to the teods, to prevent intermittent romuertions.

Variable width - width of each of 5 pulses can be adjusted independently.


CCJE MODULATED MULTIPLE-PULSE WICRDWAVE SIGNAL GENERATOR Hodel B


Pulse-time modulation-input provided in each of 5 pulse channels for external pulse-time modulation.

## 

Variahle repetition rate-repetition rate of esch group of pulses can be varied.

## CODE MODULATED MULTIPLE-PULSE MICROWAVE SIGNAL GENERATOR

Model B
950-10,750 me


#### Abstract

Generates multi-pulse modulated carrier for beacons, missiles, radar. . . provides 5 independently adjustable pulse channels, 4 interchangeable r-f oscillator heads, precision oscilloscope, self-contained power supplies ...all in one integrated mobile instrument.


The Polarad Model B is an essential instrument for testing beacons, missiles, radar, navigational systems such as DME, Tacan, H. F. Loran, etc., where multi-pulse modulated, microwave frequency energy with accurately controlled pulse width, delay, and repetition rate is required for coding.

## A fully integrated self-contained equipment with these features:

Four Interchangeable Microwave Oscillator Units - all stored in the instrument. . . each with UNI-DIAL control... precision power monitor circuit to maintain 1 mw power output reference level... keying circuit to assure rapid rise time of modulated r-f output... non-contacting chokes.
Five Independently Adjustable Pulse Channels -each channel features variable pulse width and delay; has provisions for external pulsetime modulation.
Precision Oscilloscope with Built-In Wide Band RF Detector for viewing the modulation en-

## SPECIFICATIONS:

Frequency Range:
Band 1...950 to 2400 mc
Band 2... 2150 to 4600 mc
Band $3 \ldots 4450$ to 8000 mc
Band $3 \ldots 4450$ to 8000 mc
Band 4 .... 7850 to $10,750 \mathrm{~m}$
RF Power Output . . . 1 milliwatt maximum ( 0 DBM)
Attenuator:
Output Range . . . 0 to - 127 DBM
Output Accuracy ... $\pm 2 \mathrm{db}$
Output Impedance.. . 50 ohms nominal
RF Pulse Characteristics
a. Rise Time... Better than 0.1 microsecond as measured between 10 and $90 \%$ of maximum amplitude of the initial rise.
b. Decay Time... Less than 0.1 microsecond as measured between 10 and $90 \%$ of maximum amplitude of the final decay.
c. Overshoot. . Less than $10 \%$ of maximum amplitude of the initial rise.
velope and accurately calibrating the r-f pulse width, delay, and group repetition rate. Equipped with built-in calibration markers.
Self-Contained Power Supplies-Model B operates directly from an AC line through an internal voltage regulator. The coded multipulse generator is equipped with an electronically regulated low voltage DC supply. Klystron power unit adjusts to proper voltage automatically for each interchangeable band.

Contact your Polarad representative or write to the factory for detailed information.

Internal Pulse Modulation:
No. of Channels . . . 1 to 5 independently on or off
Repetition Rate... 40 to 4000 pps
Repetition Rate..$\ddot{40}$ to 4000 pps
Pulse Width... 0.2 to 2.0 microseconds Pulse Width . . . 0.2 to 2.0 microseconds
Pulse Defay . . 0 to 30 microseconds Pulse Defay . . . . 0 to 30 microseconds
Accuracy of Pulse Setting . . 0.1 microsecond Accuracy of Pulse Setting . . . 0.1 microsecond
Minimum Pulse Separation . . 0.3 microsecond Minimum Pulse Separation ...0.3 microsecond
Initial Channel Delay ... 2 microseconds from sync. pulse
Internal Square Wave . . 40-4000 pps (separate output)
Pulse Time Madulation:
Frequency . . . 40-400 cps any or all channels
Required Ext. Mod. . . . 1 volt rms min. Maximum deviation ... $\pm 0.5$ microsecond
Power Input (built-in power supply) $105 / 125 \mathrm{v}$. 60 cps 1200 watts.


## OF SYNTHANE HAS DURABILITY, DIMENSIONAL STABILITY,

## DIELECTRIC STRENGTH

Although this sturdy end plate will fit into the palm of your hand, it has in combination all the dielectric strength, the physical properties, and the printalifity the customer requires. It's made of Synthane, a laminated plastic, the same material used in hundreds of other electrical, mechanical, and chemical applications.


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This miniaturized CST-50 variable ceramic capacitor outperforms capacitors several times larger. C.T.C.'s unique design includes a tunable element which virtually eliminates losses due to air dielect ric. This results in wide minimun to maximum capacity range of 1.5 to 12 MMFD.

This tuning sleeve is at ground potential and can be locked firmly to eliminate undesirable capacity change. Each CST-50 is provided with a ring terminal with two soldering spaces.
This is but one of a versatile family of C.T.C. ceramic capacitors of this type, built to C.T.C.'s quality control production standards for guaranteed performance.

All C.T.C. components - standard or custom - are subject to this precision manufacture. Other C.T.C. components include coil forms, coils, terminal boards, terminals, diode clips, insulated terminals and hardware. C.T.C. engineers are glad to consult on your compozent problem. Write now for sample specifications and
prices to Sales Engineering Department, Cambridge Therm:onic Corporation, 456 Concord Ave. Cambridge 38, Mass. On West Coast, contact E. V. Roberts, 5068 West Washington Blyd., Los Angeles 16 or 988 Market St.. San Francisco, Calif.
C.T.C. Capacitor Data: Metallized ceramic farms.

CST-50, in range 1.5 t 13.5 MMFDs.
CST- 6 , in range 0.5 to 4.5 MMFDs .
CS6-6, in range 1 to 8 MMFDs.
CS6-50, in range 3 to 25 MMFDs.
CST-50-D, a differential capacitor with the top hal in range 1.5 th 10 MMFDs and lower half in
range 5 to 10 MMFDs .

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To keep step with the automatic trends of industry, small components, such as the molded fixed resistors and the ceramic capacitors, are now offered by AllenBradley in reel packages wherein the components are attached to a pressuresensitive tape, ideally adapted for automatic assembly or preassembly operations.


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Now comes the commercial.
Take the best available materials (sifted by unrelenting research).
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# Industry's first full year <br> <br> performance warranty on all <br> <br> performance warranty on all transistors announced by <br> <br> General Electric 

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.. outperforms conventional mechanisms of much greater weight in a wide variety of applications . . . yet it's rugged and "tops" in dependability.

Combining improved efficiency and performance with miniaturization, RollerSmith's new Core Magnet Mechanism is an outstanding achievement . . . a precision, self-shielding movement that can be counted upon to increase the prestige of your product through consistently excellent operation.


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## newest line of precision components

 from Texas InstrumentsFor precise resistance values under extreme operating conditions, design with RADELL deposited carbon resistors - now manufactured by Texas Instruments. With resistance tolerance held to $\pm 1 \%$, Texas Instruments RADELL resistors provide exceptional stability plus a wide range of resistance values. Like all TI components, they are manufactured to exacting instriment standards.

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Write for Bulletin No. DL-C 539 giving detailed specifications of all three lines of Texas Instruments RADELL resistors. Your best source for precision components, TI also manufactures a complete line of subminiature transformers as well as custom capacitors, delay lines, special transformers and other reliable electronic components.

Hermetically sealed line - designed for extreme conditions of moisture and temperature. Specially treated ceramic shell effectively seals out moisture and air, resists abusive handling, and assures complete insulation.


MIL-Line - designed for the broad field of military applications. Exclusive multi-layer coating provides envirommental protection substantially equal to hermetic sealing throughout low and middle ranges of resistance. MIL.-Line resistors more than meet MIIL-R-10509A specifications.


# IBM selects DU MONT TYPE 329* as test oscillograph for their new type 702 computer 



When IBM Corporation, world's largest manufacturer of computer equipment, produced their new Model 702, an essential phase of the project involved selection of a cathode-ray oscillograph to go into the field with each computer as standard test equipment. Requirements were strict.
IBM's approach to the problem was to conduct side-by-side evaluation with other competitive instruments. On the basis of actual performance, they selected the Du Mont Type 329 as their test oscillograph.
What are some of the primary reasons why IBM decided on the Du Mont Type 329? Excellent sensitivity-either d.c. or a.c. coupled. precisely calibrated sweeps with movable notch magnification-ideal for making accurate measurements. Brightness-adequate for display of very fast pulses. Synchronization simplicity-
the Type 329 "locks in" on almost any type of signal. Stability-the trace remains steady as a rock despite power line fluctuations, etc. Reliability in service-calibration adjustment requires no extra test gear and is a simple one-step process. And virtually any tube may be replaced without special selection.
Another factor contributing to the selection of the Type 329 was the well known Du Mont Field Service Organization, which assures that regardless of where in the United States the equipment is used, switt, competent service facilities are in the immediate vicinity.
If you have instrumentation requirements. Du Mont facilities are always available for discussion and recommendations. Write us today for complete information on the Type 329, or on any problem you may have relating to cathode-ray instrumentation.

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# Tests show two CATHALOYS most versatile cathode materials 

## New alloys from Superior Tube simplify selection, prolong tube life

Now the engineer's job of selecting the right cathode alloy for practically any electron tube can be a simple choice between two new Cathaloys from Superior Tube.

Cathaloy A-32* is an active alloy characterized by rapid activation, high emission level throughout life, absence of interface impedance, and very low sublimation. These remarkable advantages are the result of using aluminum in place of silicon or magnesium as the reducing agent. The addition of a small percentage of tungsten also makes A-32 approximately $50 \%$ more shock resistant than cathodes without tungsten. Thus A-32 is suitable for virtually any active alloy application, including ruggedized tubes.
Cathaloy P-50 is a passive alloy of carefully controlled analysis that is commercially available in Weldrawn t cathodes as well as Lockseam. $\ddagger$ It can be made in Weldrawn form because of its capacity to take much more severe reductions in cold drawing without rupture than other grades of passive alloys. P-50 is identical in composition with the well-known ASTM Grade 21. The important difference is in the method of melting which improves the uniformity and completeness with which deoxidation is accomplished. All heats are tested in Superior Tube's laboratory before being approved for production.
Ask for complete technical reports on both these new Cathaloys. Write Superior Tube Co., 2506 Germantown Ave., Norristown, Pa.


DN TEST. Laboratory photo of test diodes used in Superior Thbe"s electronic labcratory. Under exhaustive tests, the new Cathaloys display performance characteristics nol present in other alloys.
*Patent applied for



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New, inherently rugged mechanical designs permit safe operation at substantial peak and average powers. Cascade Research has designed these units to optimize iso-lation-to-insertion loss-ratio. Specicl techniques have made possible a reduction in size and weight of the integral permanent magnets. As in all Unilines, no external power source is required.
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Richard P. Gaunt (M'53) has joined the Lockheed Missile Systems I Ivision as a member of the research staff. Mr. Gannt has been associated or many years with the Call'Tech Jet Propulsion Laboratory where he was a Senior Rescarch Engineer in charge of a missile systems gronp.

He has worked in the design and development of the guidance system for

R. P. (BaCNT the Corporal missile concentrating in the design of specialpurpose analog computers and differential analyzers.

Mr. Gaunt was born in Mount Vernon, N. Y. and received the I3.S. in Physics from Brown University in 1948. His hobby is music and at the present time he is studying modern jazz piano. He is a member of the Beta Theta I'i fraternity, the IRF: and Sigma Xi.

Stuart L. Bailey, President of Jansky \& Bailey, Inc., Washington, I). C. recently announced the election of Delmer C. Ports (A'38-SM'45), Chief Engineer, to the position of VicePresident.

Mr. l'orts received the I3.S. degree in Electrical Engincering at the George Washington Injversity and the M.S. degree from Ohio State Eni-

II. C. PORTS versity.

As an engineer for Jansky \& Bailey, which he joined in 1936. Mr Ports was first engaged in antenna design and installation projects and in wave propagation studies. He also worked in the development of high frequency measuring techniques and equipment.

W'ith the advent of World W'ar II, he took part in a mumber of research projects, sponsored by the Office of Scientific Research and Development, related 10 antenna characteristics, propagation phenomena and communication systems. He was in charge of programs to develop terhnigues and measure characteristics of high frequency antennas and supervised research which experimentally isolated someof the factors affecting propagation in the high frequency and very high frepueney regions.

Daring this period Jansky \& Bailey established a laboratory to measure commanication systems in ase by the military: Mr. I'orts assumed the responsibility for this work.

He is a member of the American Insti-

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clare Type hg and hgp Relays offer a combination of high speed, high current- and voltage-handling capacity. and extraordinary uniformity of performance over very long periods.

The relays consist of a magnetic switch, hermetically sealed in a high pressure hydrogen atmosphere in a glass capsule, and a coil, enclosed in a steel vacuum-tube type envelope which has a standard medium-sized octal base. Platinum contact surfaces are continually wetted with mercury by means of a capillary connection to a mercury reservoir below the contacts. Type hap Relays can be factory adjusted to provide either biased or polarized operating characteristics.

For complete iniormation on the neid slare Type hg and hGP Mercury-Wetted Contact Relays, contact your nearest clane representative or address C'. P. Clare \& Co., 3101 Pratt Blud., Chicago 45, Illinois.

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Relay amplifiers
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LONG LIFE: Conservative life expectancy of over a billion operations when operated within ratings.

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NO SONTACT BOUNCE

## MECHANICAL FEATURES

- Small chassis space required
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REL-14-2 and REL-14-3 have slightly larger dimensions.

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Youcan Rely on...

F. Stanley Atchison (A'53-SM'53) has been appointed Technical Director of the U. S. Naval Ordnance L.aboratory, Corona, California. He succeeds Ralph A. Lamm who is joining the stalf of Bendix Aviation Corporation.

Dr. Atchison was born in Missouri and received the I'h.I). in physis from the State I'niversity of Iowa in 1942. Juining the

F. S. Archison staff of the National
lharean of Standards in 1942 ds a physicist in the Ordnance Development Division, he worked on proximity fuzes and went to the Mariana Islands as technical advisor to the Army Air Force during their first use of proximity fuzed bombs at the invasion of Iwo Jima. Dr. Atchison was subsernently put in charge of the Missile Intelligence Section, where he was concerned with the design of electronic systems for guided missile control. Ite noved to California in 1951 when the National Bureau of Standards established its Corona Laboratories, and directed a special research project on developments for electronic computers. Last year he was appointed head of the Phisical Science Department, which is engaged in research in physics, chemistry, and electronics.

Dr. Atchison is president of the Sigma Xi Club) at the University of California at Riverside and a member of the American Physical Society.

The TelAutograph Corporation has recently appointed R. G. Leitner (SM'53) to the post of Chief Engineer. In his new position, Leitner will direct the development of Tel.Autograph communication systems and also the company's expanded program into automation, nucleonics, and electronic instrumentation.

Mr. Leitner formerly served as Chief Research Engineer for Fackard Bell Company, and prior to that time in a similar


Douglas chose the new Kollsman KS-54 Cabin Pressure Control System for their new DC-7C's because of the many decided advantages it offers over the other existing systems.

LIVING-ROOM COMFORT IN THE CABIN . . . There is no annoying ear-popping because cabin pressure is held practically constant under cruising conditions. Even when cabin pressure is changing, the rate of change is so smoothly controlled that the actual change of pressure is unnoticable.

PEACE OF MIND IN THE COCKPIT . . . When the controls are set, the system is fully automatic and thoroughly reliable - especially so because of the simplicity of the Kollsman design.

NO WORRY IN THE MAINTENANCE SHOP . . . The components are simple and rugged, proven dependable and require a minimum of maintenance. There are no sensor contacts or filters to clean, no complex tubing to worry about.

WRITE for special folder giving full technical details on the new Kollsman KS-54 System, or ask to have a sales engineer visit you.


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## AnII E Peofle

(cominted from poge 52.4)
pensition at Lear, Incorporated. He was Chairman of the 1952 Western Electronic Show and (onvention and has been a Director of the West Cowst Blectronic Mannfacturers dosociation ats well as lice-Chairman of the las Angeles Comncil.

About 200 of his asembates juined in a luncheon this spring at the Naval Research Laloratory to monor Samuel D. Summers (SM'50), consultant to the Electrenics Disision, who is retiring from Civil Service.

A native of Hornbeaik, Temm., Mr. Summers attended Memphis State College; he received the Bachelor of Science in Electrical Enginerering from the Tri-State College in Angola, Indiana. He did graduate work at the Massachusetts Institute of Technology and the Linversity of Michigan and received the Master's degree in e.e. from the C"niversity of Maryland.

He was an instructor in electrical engineering at Tri-State College in 192.3 and later became a professor at the school. For three years, he was head of the Electrical Enginering I epartment. From 1925 to 1931, Mr. Summers was an electrical engineer with the Commonwealth \& Southern Corporation at Jackson, Michigan, where he designed electrical power stat tions and transmission lines.

During World War 11, Mr. Summers served as an electronics officer in the Navy, and was assigned to the Naval Research Latoratory as a member of the Avation Electronics Service I'nit, He joined NRI, asanadvisor on problems relating to naval electrical and efectronics applications: in 1945.

Arr. Summers is a member of the Ameriean Institute of Electional Engineers.


The following transfers and admissions were approved and are now effective:

Transfer to Senior Member
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This is a development which calls for immediate changes in purchasing specifications for Tape Wound Cores, because introduction of the Aluminum Core Bex means designing your torgids around four important new advantages:

1. Use of an aluminum core box means the new Magneties, Inc. tape wound cores will withstand temperatures of at least $450^{\circ} \mathrm{F}$.
2. Because of the unustal seal provided by forming the aluminum over the silicone glass seal, true vacum impregnation of your coils is now possible. Varnish cannot penetrate the core box and affect magnetic properties of the tape.
3. The strong aluminum construction absolutely prevents deflection of the core box when coils are wound-a distortion-free construction which means no change of magnetic propertics.
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Because of the many artvantages of these new Magnetics, Inc. Tape Wound Cores, it will pay you many times over to specify "Aluminum Core Boxes" on your next orter.
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Immediately available in 109 standard sizes, using all commercially available magnetic materials.

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SECON'S highly engineered fine Wire is being used to meet critical SPECIFICATIONS BY MANUFACTURERS OF IMPORTANT DEFENSE AND MILITARY END-USE ITEMS.

## Precision Wire-Wound Potentiometers

In supplying jrecious metal alloy wire for these, Secon not ouly conforms to the physical and electrical characteristics on the manufacturer's specifications, but also winds a prototype potentiometer from each melt, which is tested for life, noise, and other characteristics which cannot be specified on the wire. Roundness of so small a magnitude that it cannot be measured is a carefully controlled characteristic which receives Secon's contimuous attention.

## Direct-Heoted Cathodes in

Electronic Vacuum Tubes
Wire and ribbon for use here are individually prepared for each manufacturer to insure satisfactory operation. Secon sets aside the melts until the manufacturer has ascertained the emission and life characteristics of the melt. Approved Secon melts are then used exclusively to supply the manufacturer who made the tests,

## Electro-Ploted Grid Wire for <br> Electronic Vacuum Tubes

Precious metals used for these are carefully selectel for purity. Only high purity gold, rhodium, silver and others are employed.

## Strain Gauge Wires

These are most carefully selected, in both precious and base metals. Samples of Secon melts are tested by the manufacturer of the strain gauge for temperature coefficient of resistance, gauge factor, and other important characteristics. To insure uniformity, Secon sets aside approved melts for the exclusive use of the manufacturer who made the tests.

## New Wire Products for Semi-

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Developed through special rescarch for application in these fields, the new products include:

Gold: fine gold in purities up to $99.99 \%$; and doped gold alloys.
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II'hisker Wires: in base metals as well as hard platinum alloys, with close tolerances on straiglitness and hardness in all types.
Lead-in W'ires of a great variety such as timned copper wire or ribbon.

Secon specializes in the development, research and production of special alloys and pure metals, processed to very small diameter wire-in all shapesround, oval, flat, ribbon, grooved-for highly engineered applications in electronics, instrumentation, ordnance, aviation, nuclear physics, atomic energy, guided missiles, attomotive industry, and other fields.

Close tolerances and controlled specifications can be held on many important characteristics such as: resistance, tensile strength, elongation, surface appearance, special spooling, purity, torque, linearity, composition, cross section, weight per unit length, uniform plating, dependable insulation, temperature coefficient of expansion and resistance, and strain sensitivity.

Secon end-products include:

| - Fine W'ire drawn to 0.0003" diameter |  |
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| - | Electro-Plated Wire and Hibbon |
|  | pecial Solder |
| - | Enamieled and Inaulated Wire |
|  | irani Gauke Wire |
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| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  | 3-5 | 5 | 10 | 20 | 50 | 100 | 150 | 200 |
| 30 | $200 \mu \mathrm{~A}$ @ - 20 V |  |  |  |  |  | H1) $215 \%$ |  |  |
| 40 | $10_{\mu} \mathrm{A}$ ( $)^{\text {a }}$ - 10 V | 1N148* |  |  |  |  |  |  |  |
| 60 |  |  | $\begin{aligned} & 1 \times 116 \\ & 1.90 \\ & 1 \times 126^{*} \end{aligned}$ | $\begin{aligned} & 1 \times 117 \\ & 1 \times 05 \end{aligned}$ | $\begin{aligned} & 1 \times 118 \\ & 1 \times 96 \end{aligned}$ | $\begin{aligned} & \text { H1) } 2167 \\ & \text { HI) } 2166 \\ & \text { HI } 2155 \end{aligned}$ | (11) 2173 (11) 2174 H1) $216 z$ |  | $\begin{aligned} & \text { 111) } 2160 \\ & \text { 111) } 2171 \\ & 111) 217 z \end{aligned}$ |
| 80 |  | $\begin{aligned} & 1 \times 67 \mathrm{~A} \\ & 1.189 \end{aligned}$ | 1.191** <br> 1N102** <br> 1N198* | $\begin{aligned} & 1 \times 9 \\ & 1 \times 97 \end{aligned}$ | $\begin{aligned} & 1 \times 100 \\ & 1 \times 98 \end{aligned}$ | $\begin{aligned} & \text { H1) } 2151 \\ & \text { H1) } 2168 \end{aligned}$ $\text { HH) } \geq 169$ | $11!\geq 1.50$ $\text { (11) } 2163$ $\text { HI } 217.5$ |  | $\begin{aligned} & \text { III } 21.58 \\ & \text { III) } 21.57 \end{aligned}$ $\text { H1) } 2154$ |
| 100 | $\begin{gathered} 180 \mu .1 @-100 \mathrm{~V} \\ 500 \mu .1 @-100 \mathrm{~V} \\ 625 \mu .1 \text { @ } 100 \mathrm{~V} \\ 300 \mu .1 @-505 \\ 50 \mu .1 @-50 \mathrm{~V} \end{gathered}$ | 1N68A <br> 1. 127* <br> (II) 20.51 |  |  |  | 111) 2170 | H1) 2165 | (1I) 2154 | (i) 2161 |
| 150 | 500u. 1 () -150 V |  | 1.V5.513 |  |  |  |  |  |  |
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## Performance Characteristics

## Repetition Rate

## Fall Time

 Rise Time Input Impedance (pulse)Output Impedance
(pulse)
Trigger Pulse
Output Voltage

Load Current

Power
Requirements
Operating
Temperature

40 kc max.
$8 \mu \mathrm{sec}$ to resistive locad 2 usec

3500 ohms
3500 ohms
12 volts of $.5{ }_{10} \mathrm{sec}$ duration min. (a) $40 \mathrm{kc} / \mathrm{sec}$. 20 volts peak to peak (unit clamped as -20 and 0 volts)
2 ma of current may be drawn without destroying clamped levels
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They exceed gll MIL-R-10509A specifications as you can see from the comparison table below. Note, for example, that they take iull power at ambient temperatures up to $120^{\circ} \mathrm{C}$ instead of only $40^{\circ} \mathrm{C}$. Thus, they are ideal for use in aircraft and guided missiles. The same fact, of course, will result in much longer life when they are operated at lower temperatures.

POLYOHMS are well suited to replace bulky, expensive and highly inductive wire-wound resistors.
The resistor will remain well with. in its $1 \%$ tolerance even under the stringent moisture test which allows a $5 \%$ change. Its temperature coefficient is always lowet than both the $R$ and $X$ characteristics.
POLYOHMS are manufactured in $1 / 2,1$, and 2 watt sizes with facilities controlled by the Signal Corps. They are presently avail. able only for government end use. Please request samples on com. pany letterhead.

TABLE OF TEST RESULTS

| TEST | MIL-R-10509A <br> Allowable change | POLYOHM Test Results <br> (Median Value) |
| :---: | :---: | :---: |
| Temperature cycling | $1 \%$ | $.03 \%$ |
| Low temperature exposure | $3 \%$ | $.08 \%$ |
| Short time overload | $.5 \%$ | $.03 \%$ |
| Load life @ $40^{\circ} \mathrm{C}-1000 \mathrm{hrs}$. | $1 \%$ | $.2 \%$ |
| @ $120^{\circ} \mathrm{C}-1000 \mathrm{hrs}$. | - | $.5 \%$ |
| Temp. coef. ppm $/{ }^{\circ} \mathrm{C}$ (spec.X) |  |  |
|  | $\pm 500$ | -150 |
| (spec.R) | $\pm 300$ | -150 |
| Moisture resistance test | $5 \%$ | $.3 \%$ |



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Ideal for production-line testing and laboratory work, this new VTVM provides direct readings without interpolation. Features illuminated digital scale with decimal point and polarity sign; 12 ranges ( $A C, D C$, ohms); frequency response to 250 mc with auxiliary probe; accuracy: $1 \%$ on DC and ohms, $2 \%$ on $A C$. Cuts multiple scale confusion and learning curve error.


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*RETMA Report
months of 1955.

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Director, 1955


Howard Vollum was bor: on May 31, 1913, in P'ortland, Oregon. He attended Columbia University in Portland from 1931 to 1933, transferring to Reed College in 1934. In 1936 he received the $13 . A$. degree in physics from the latter school.

Upon graduation from college, he spent the next few years servicing and installing home, auto, and aircraft radios and constructing electronic devices. From 1940 to 1941 he was Supervisor of Radio Project, NYA, in Portland. Mr. Vollum served as an officer in the U. S. Army Signal Corps from 1942 to 1946. Ilis first two years in service he spent at ADRIIE, in Malvern and Christ Church, England, working on coast artillery fire control radar. Ile was awarded the Legion of Merit for this work. For the next two years he was stationed
at the Evans Signal Laboratory in Belmar, New Jersey, in charge of a subsection concerned with the use of radar by ground forces. As a result of this contribution, the Oak Leaf Cluster was added to his award.

In 1946 Mr . Vollum helped to found Tektronix, Inc., of which he is now President. Mr. Vollum is known for his work on the development of the cathode-ray oscilloscope. In recognition of his achievements, Portland University awarded him an honorary Sc.D. degree in 1953.

Mr. Vollum became a Senior Member oi the IRE: in 1950, and received the Fellow Award in 1955, "for his contribution to the development and manufacture of electronic laboratory instruments." In 1954 he was Chairman of the Portland Section.

# Changes in the IRE Dues Structure 

J. D. Ryder, President

Because of the breadth of the electronics field, including as it does most aspects of electrical engineering and many areas in physics, the IRE has become one of the large and well-respected professional engineering societies of the world. As such, W. R. Hewlett, our president in 1954, believed it undesirable that our organization should have so many qualified professional members in the Associate grade of membership, and therefore unable to vote and participate fully in IRE activities. At the March, 1955, meeting of the Board of Directors, certain changes were made in the membership and dues structure of the IRE, which it is hoped will channel new members more directly into membership grades properly representative of their professional qualifications.

For many years the IRE dues pattern has provided for annual dues of $\$ 10$ for the first five years of Associate membership, all other dues being $\$ 15$ per year. This arrangement has furnished in the past, a financial incentive to initial application at the Associate, or non-professional level, and normal human inertia, or reluctance to fill out and file transfer blanks, has kept many professionallyqualified members there. It would seem more desirable, however, that a potential member's professional qualifications, and not a favorable dues condition, should determine his initial grade of IRE membership, and this the new dues structure aims to provide.

While other engineering societies are finding the going difficult and are raising dues, the IRE remains in strong financial condition, and it seemed inadvisable and unnecessary to adopt the obvious possibility of eliminating the $\$ 10$ initial Associate dues level. Accordingly, the Board of Directors has adopted the reverse policy and has ruled that for all new members elected after July 1, 1955, regardless of grade of membership, the annual dues shall be $\$ 10$ for the first three years of membership, after which the dues rise to the present maximum level of $\$ 15$ annually. This will represent a dues reduction for new elections as Members or Senior Members. It is hoped thereby that new applicants will enter in the grades for which they are qualified, and thus will strengthen the IRE
through increase of qualified professional membership.

Several other changes have been made in the bylaws governing membership. Age limits have been eliminated in all grades, the Board feeling that anyone who has achieved the required professional standing should be permitted the privileges of that standing, irrespective of age. It has also been pointed out that the Associate grade requires merely "an interest in radio," whereas the graduate of an accredited four-year school who has specialized in radio, electrical engineering, or allied studies certainly has progressed beyond mere interest in radio. This fact has been recognized through bylaw changes allowing such graduates to enter as, or transfer to Member level, the lower professional grade, by granting of three years of professional credit for the four school years instead of the previous two year credit. IRE Student Members graduating from such curricula will henceforth transfer to Member instead of Associate grade, upon expiration of Student Member status on graduation.

It should be noted that the changes in dues structure are to affect only new members elected after July 1, 1955, and do not apply to members admitted before that time. Thus Associates elected prior to that date will continue their five years at the rate of $\$ 10$, whereas Associates, Members, or Senior Members elected after July 1, 1955, will have three years at the new rate. To newly elected Members and Senior Members the revised plan represents a slight reduction in dues.

It seems desirable to reemphasize the feeling of the Board that the strength of the IRE depends on sufficient numbers of professionally qualified members. and that it is of great importance for every nember to be in the highest membership grade for which he is professionally qualified. To this end. Section Membership Chairmen stand ready to aid. It is also suggested that every Associate and Member study page 5 of the 1954 Directory to determine if he has the qualifications necessary to a higher membership grade. Such transfer can be of much value to each individual and to the IRE in furthering its professional responsibilities.

# Frequency Aging of High-Frequency Plated Crystal Units* 

A. W. WARNER $\dagger$, MEMBER, IRE

> The following paper was procured and recommended for publication in the Proceedings by the Professional Group on Ultrasonics Engineering.-The Editor


#### Abstract

Summary-The frequency aging of high-frequency crystal units is explained in terms of residual contamination, which may be a partial molecular layer. Experimental data are given on methods designed to reduce frequency aging to a minimum.


THE ALIOWABLE change in frequency of a crystal unit with time is of ten specified in parts in $10^{6}$ per month, and for the more precise frequency standards is specified in parts in $10^{8}$ per week. These tolerances are several orders of magnitude beyond those required of other standards used in communications measurements. To maintain and improve such tolerances requires the application of a high degree of skill and engineering ingenuity. It is the purpose of this article to show the relationship betweeen various fabrication methods and frequency aging, and to present typical frequency aging data on plated AT crystal units in metal and glass enclosures.


Fig. $1-$ Aging record at 70 degrees C. of a CR32 type crystal unit' 44 megacycles, fifth overtone, enclosed in an HC6 metal holder.

The resonant frequency of a crystal unit is determined by its mechanically vibrating section, consisting of the quartz plate, electrodes deposited on its two major faces, and to some extent the mounting or support wires. ${ }^{1}$ At sufficiently high frequencies, or with specially contoured quartz plates, ${ }^{2}$ the mounting can be effectively divorced from the vibrating system, leaving a part of the crystal plate and its associated metal electrodes as the fre-

[^4]quency determining elements. Any change in the clastic constants or the mass of these elements will cause a change in frequency.

The aging of typical AT crystal units can be best explained by the transfer to or from the quartz plate of material other than quartz or electrode metal. This is not to say that aging cannot le caused by loss of quartz from the surface and by migration within the metal electrode, but with present-day methods of etching quartz and the use of noble metals for electrodes, aging from this source is much less prevalent than from contaminants.

Fig. 1 shows the aging record at 70 degrees C. of a 44 -megacycle, fifth-overtone, CR32-type crystal unit enclosed in the commonly used metal IIC6 holder. The aging is two parts per million per month for the first month of operation. This aging rate is about normal for


Fig. 2-Aging record at 70 degree C . of a contaminated metal enclosed crystal unit, $15-\mathrm{mc}$, third-overtone.
the CR32 type of crystal unit. The downward direction of the aging curve is attributed to the slow transfer of contamination to the crystal plate from the inner walls of the metal enclosure. This assumption seems reasonable in view of the fact that the quartz crystal plate and its gold electrode can be cleaned more effectively than can the metal enclosure.

Fig. 2 shows the aging record of a similar crystal, except that additional contamination. probably soldering flux, is present. This is surmised from the fact that the unit failed to pass a test for low vapor content. The aging rate is four times as bad and, since the direction of frequency change is reversed upon exposure to room
temperature, it is evident that equilibrium is a function of temperature. Although this crystal unit might eventually age at an acceptably low rate, any change in temperature, such as an oven shut down, would start a new aging cycle.


Fig. 3-Aging record of metal enclosed, $89-\mathrm{mc}$, seventh-overtone crystal units at 70 degrees C .

Fig. 3 shows the aging record of a group of three 89mc , seventh-overtone cry'stal units in metal holders at 70 degrees $C$. At this frequency the quartz plate is less than half as thick as that of Fig. 2, yet no aging is apparent within the error of measurement, $\pm 1$ part per million. These units were baked for 24 hours prior to assembly. Fig. 4 shows the aging record of two thirdovertone crystal units made in the laboratory, with great care taken in the cleaning and sealing methods. The aging is about 1 ppm for the extended period of four months, five to ten times better than the crystal in Fig. 1.


Fig. 4-Aging record of two third-overtone laboratory-made crystal units, metal enclosed.

In view of the foregoing experiments, it was reasoned (1) that a glass enclosure, with a smooth surface which can be easily degassed and cleaned, should be superior to a metal enclosure, and (2) that by the use of glass, aging rates low enough for primary frequency standard
use might be achieved. 'Io demonstrate this a group of crystal plates having a large frequency determining dimension, 10 times that of the crystal unit in Fig. 1 were tried, first in evacuated HC6 metal enclosures and then in glass bulbs. The frequency measuring equipment was improved to measure parts in $10^{8}$ rather than $10^{6}$. Using a 25 X expanded frequency scale, it can be seen in Figs. 5 and 6 that in the metal enclosure the frequency aging is slightly erratic and downward, and in the glass enclosure it is uniform and upward. The conclusions reached from this experiment were (1) that the


Fig. 5—Aging record of 5-mc, fifth-overtone crystal units in HC6 metal enclosures, at 60 degrees C .


Fig. 6-Aging record of $5-\mathrm{mc}$, fifth-overtone crystal units in T5-1/2 glass bulbs, at 60 degrees C .
glass enclosure was contaminated less than the crystal plate surface, which is apparently losing mass to its surroundings, and (2) that for further improvement of frequency aging the crystal plate surface should now be improved.

The next experiments were performed on polished quartz plates. It was reasoned that a smooth surface would not only have less surface area, but also could be more easily cleaned. The contaminants are not imbedded in tiny crevices and can be removed by rela-
tively short exposures to cleaning agents. Fig. 7 shows the aging record of the polished crystal plates in glass enclosures. These are twice as good as the unpolished plates and are uniform, but the indications are still that the crystal plate is not entirely free of contamination.


Fig. 7-Aging record of polished 5 -mc, fifth-overtone crysta units in glass bulbs, at 60 degrees $C$.

This was investigated further in an aging test designed to show the relative amounts of contamination removed by various processes. Referring to Fig. 8, crystal units number 5 and 6 were made as in the previous experiment. Crystal unit number 4 is a similar


Fig. 8-Aging record of various polished $5-\mathrm{mc}$, fifth-overtone crystal units in glass bulbs at 60 degrees C. with a 24 -hour 100 degrees C . bake at intervals.
unit that had been through the previous aging test. Crystal units 1, 2, and 3, were made in the same manner as units 5 and 6 , except that they were vacuum baked in a special apparatus just prior to final sealing. During the aging test the crystals were removed periodically and exposed to a temperature of 100 degrees C. for 24
hours. It will be seen that the vacuum baking technique had the effect of removing about 80 per cent of the remaining contamination, making the crystal unit much less susceptible to frequency changes due to temperature interruptions.

Fig. 9 shows the aging record of a group of five crystal units made in accordance with the above principles; that is, smooth surfaces and maximum removal of contamination. The average aging of the five units is nearly zero, the maximum excursion 10 parts in a billion in 30 days.


Fig. 9—Aging record at 75 degrees C. of crystal units baked 20 hours at 140 degrees C .

From these data one can conclude empirically that the mechanism of frequency aging of high-frequency crystal units is one of transfer of mass to and from the crystal plate. The rate of transfer and the degree of permanence after transfer will be a function of the vapor pressure of the contaminant and the degree of adherence, which may be molecular, chemical, or mechanical, between the contaminant and the electrode material. For this reason, it is not likely that any degree of permanence would be achieved by stabilizing a contaminated crystal after sealing, as by heat cycling and the like. Also, it is probably not possible to predict an aging curve, since the equilibrium reached at any given operating condition is not likely to repeat itself.

The relationship between contaminant and frequency change may be clearer if one calculates the mass involved. A change in frequency of 1 part in $10^{9}$ in the crystal units of Fig. 9 requires a change in mass at the surface of $2 \times 10^{-4}$ micrograms per square centimeter. If this mass were, for example, an oil film having a density of 1 , it would be only two hundredths of an angstrom thick. Since a molecular layer is at least $4 \AA$ thick, this represents only a partial molecular layer.

It should be recognized, however, that a high degree of frequency stability is obtainable from crystal units mounted in the widely used HC6 metal enclosure, i.e., 1 part in $10^{6}$ for the first month of operation with continuing improvement as long as operating conditions are unchanged. Where frequency stability greater than this is required, as in primary frequency standards, the cleaner glass-enclosed crystal units are preferable, despite their larger size.

# Some Gyrator and Impedance Inverter Circuits* 

B. P. BOGERT $\dagger$


#### Abstract

Summary-The use of feedback amplifiers to provide impedance inversion is considered. If the proper types of feedback connections are used, and the input and output impedances of the amplifier and its gain are chosen suitably, the input impedance of the circuit will be approximately proportional to the reciprocal of the load impedance. The approximation may be improved by the insertion of negative impedance elements to compensate for the residual impedances in the inverter. When this is done, the circuit exhibits gyrator properties. Experimental verification of one of the circuits, using high quality stable amplifiers, gave a range of impedance inversion of two decades.


II'T IS THE purpose of this paper to discuss some 'feedlback amplifier circuits which act as impedance inverters and as gyrators. The gyrator ${ }^{1}$ may be considered to be a two terminal-pair circuit which has the property that the phase shift for transmission in one direction differs by 180 degrees from that for transmission in the other direction, over a broadband of frequency. ${ }^{2}$ This is a consequence of the fundamental property of a gyrator, which is that the transfer impedance $z_{12}$ be $R, z_{21}$ be $-R$, and that $z_{11}$ and $z_{22}$ be zero. Another property which follows from the above is that of imperlance inversion. If, as a starting point, we consider circuits capable of impeclance inversion, it is possible to arrive at some active circuits for gyrators.

We consider the relation between the input impedance and the load impedance of amplifiers with external feedback connections of specified types. By using the proper feedback connections, and by suitable choice of amplifier input and output impedances and gain, the circuit ap)proximates an impedance inverter. The residual impedances nay be corrected by addition of negative imperdance elements, and the resulting circuits possess the properties of a gyrator. Since the gyrator circuits so obtained employ active elements, their over-all stability is an important consideration.

Consider first the fourpole shown in Fig. 1(a). The input impedance $Z_{i n}$ when an impedance $Z_{L}$, terminates the output is given by the expression

$$
\begin{equation*}
Z_{i n}=z_{11}-\frac{z_{12} z_{21}}{z_{22}+Z_{L}} . \tag{1}
\end{equation*}
$$

In order to make an impedance inverter, we must have $z_{11}=z_{22}=0$, and $z_{12} z_{21}$ negative real. If we cannot strictly realize the first condition, we should have, in the impedance range of interest,

$$
\begin{equation*}
z_{22} \ll Z_{L} \tag{2}
\end{equation*}
$$

[^5]and
\[

$$
\begin{equation*}
z_{11} \ll \frac{z_{12} z_{21}}{Z_{L}} \tag{3}
\end{equation*}
$$

\]

Consider now the circuit of Fig. 1(b). It consists of an amplifier having an input impedance $Z_{1}$, a passive output impedance $Z_{2}$ and open circuit gain $k$. The output is fed back into the input using a series connection at the output end and a parallel connection in the input circuit.


Fig. 1-(a) Basic fourpole with current and voltage conventions. (b) Impedance inverter circuit. (c) Alternative form of circuit of Fig. 1(b). (d) Grounded emitter transistor stage.

If the output circuit is terminated in an impedance $Z_{L}$, the input imperlance $Z_{\text {in }}$ is:

$$
\begin{equation*}
Z_{i n}=\frac{Z_{1}\left(Z_{2}+Z_{L}\right)}{Z_{1}(k+1)+Z_{2}+Z_{L}} . \tag{4}
\end{equation*}
$$

If we make

$$
\begin{equation*}
Z_{1}(k+1)+Z_{2}=0, \tag{5}
\end{equation*}
$$

and

$$
\begin{equation*}
Z_{2} \gg Z_{L} \tag{6}
\end{equation*}
$$

then

$$
\begin{equation*}
Z_{i n} \cong \frac{Z_{1} Z_{2}}{Z_{L}}=\frac{Z^{2}}{Z_{L}}, \tag{7}
\end{equation*}
$$

which has the desired property. This circuit is easily recognized as a form of feedback amplifier, ${ }^{3}$ but its action as an impedance inverter does not appear to be particularly well known.

An alternative circuit is obtained by interchanging the feedback leads and replacing $k$ by $-k$. This circuit is shown in Fig. 1(c), which makes it evident that we are dealing with the ordinary reactance tube circuit, except that for operation as an impedance inverter, the phase shift in the amplifier must be zero when $Z_{1}$ and $Z_{2}$ are resistances. Thus we are led to consider the use of a
${ }^{3}$ R. B. Blacknan. "Effect of feedback on impedance," Bell Sys. Tech. Jour., vol. 22, p. 276 (F ig. 4); October, 1943.
transistor, using the circuit shown in Fig. 1(d). This circuit is essentially the grounded emitter transistor circuit, and the input impedance of such a stage is given in terms of the transistor constants as: ${ }^{4}$

$$
\begin{equation*}
R_{i}=r_{e}+r_{b}+\frac{r_{e}\left(r_{m}-r_{e}\right)}{r_{\bullet}+r_{c}-r_{m}+R_{L}} . \tag{8}
\end{equation*}
$$

and, looking from the output end:

$$
\begin{equation*}
R_{0}=r_{e}+r_{c}-r_{m}+\frac{r_{e}\left(r_{m}-r_{e}\right)}{r_{e}+r_{b}+R_{e}} . \tag{9}
\end{equation*}
$$

If we employed a transistor having the idealized properties $r_{b} \rightarrow 0 ; r_{c}-r_{m} \rightarrow 0 ; r_{e} \rightarrow 0$ and $r_{m} \rightarrow \infty$ such that $r_{\mathrm{g}} r_{m}=R^{2}$, then a good impedance inverter could be made by using a grounded emitter transistor stage.

It is convenient to regard these circuits as feedback a mplifiers, and to characterize them in terms of the feedback connections at the input and output. The circuit we have been discussing is parallel connected at the input and series connected at the output. It will be denoted by PS. Let us see how other circuits behave as regards impedance transformations. We consider, in order, the series-series (SS), series-parallel (SP), parallelseries (PS), and parallel-parallel (PP). Using the sym-


Fig. 2-General forms of feedback amplifier circuits discussed. $S S=$ series-series; $S P=$ series-parallel ; $P S=$ parallel series; $P P$ $=$ parallel-parallel. First letter refers to input feedback connection, second letter to output feedback connection. Meaning of amplifier symbol shown.
bols and conventions shown in Fig. 2 we have for the impedance transformations:

$$
\begin{align*}
& \mathrm{SS}: Z_{\text {in }}=(1-k) Z_{1}+Z_{2}+Z_{L}=\alpha+Z_{L}  \tag{10}\\
& \mathrm{SP}: Z_{\text {in }}=\frac{Z_{1} Z_{2}+Z_{L}\left[Z_{1}(1-k)+Z_{2}\right]}{Z_{2}+Z_{L}}=\frac{Z_{1} Z_{2}+\alpha Z_{L}}{Z_{2}+Z_{L}} \\
& \mathrm{PS}: Z_{\text {in }}=\frac{Z_{1}\left(Z_{2}+Z_{L}\right)}{Z_{1}(1-k)+Z_{2}+Z_{L}}=\frac{Z_{1}\left(Z_{2}+Z_{L}\right)}{\alpha+Z_{L}}  \tag{12}\\
& \mathrm{PP}: Z_{\text {in }}=\frac{Z_{1} Z_{2} Z_{L}}{Z_{1} Z_{2}+Z_{L}\left[Z_{1}(1-k)+Z_{2}\right]}=\frac{Z_{1} Z_{2} Z_{L}}{Z_{1} Z_{2}+\alpha Z_{L}} \tag{13}
\end{align*}
$$

where $\alpha=Z_{1}(1-k)+Z_{2}$, and is assumed real.

[^6]Inspection of (10) and (13) show that it is impossible to obtain impedance inversion with the SS and PP feedback connections. We have already discussed the PS connection so it remains to consider the SP circuit. For good inversion we need

$$
\begin{equation*}
Z_{2} \ll Z_{L}, \quad \alpha=0 \tag{14}
\end{equation*}
$$

in which case

$$
\begin{equation*}
Z_{i n}=\frac{Z_{1} Z_{2}}{Z_{L}} \tag{15}
\end{equation*}
$$



Fig. 3- (a) Ideal cathode follower type impedance inverter. (b) Gyrator formed from SP inverter circuit with addition of series resistance $-R$ in output lead. (c) Gyrator formed from SP inverter circuit with addition of shunt resistance $-R$ across input. (d) Gyrator formed from PS inverter circuit with addition of shunt resistance $-R$ across output. (e) Gyrator formed from PS inverter circuit with addition of series resistance $-R$ in input lead.

A possible circuit is shown in Fig. 3(a), which involves the use of an ideal cathode follower. ${ }^{5}$ The impedance to be inverted is placed from cathode-to-ground and the inverted impedance appears from cathode-to-grid. ${ }^{6}$
For the SP and PS cases we have for the fourpole impedance matrix:

$$
\begin{align*}
& \mathrm{SP}:\|z\|=\left\|\begin{array}{cc}
\alpha & -Z_{2} \\
Z_{1}-\alpha & Z_{2}
\end{array}\right\|  \tag{16}\\
& \mathrm{PS}:\|z\|=\left\|\begin{array}{cc}
Z_{1} & -Z_{1} \\
Z_{2}-\alpha & \alpha
\end{array}\right\| . \tag{17}
\end{align*}
$$

The admittance matrix in each case is

$$
S P:\|y\|=\left\|\begin{array}{cc}
Y_{1} & Y_{1}  \tag{18}\\
\alpha^{\prime}-Y_{2} & \alpha^{\prime}
\end{array}\right\|
$$

${ }^{5}$ J. Shekel, "The gyrator as a 3-terminal element," Proc. IRF, vol. 41, pp. 1014-1016; August, 1953.

- If the circuit of Fig. 3(a) is redrawn to bring into evidence the combination of the internal feedback of the catiode follower stage and the external SP feedback under discussion, a differently appearing circuit results, which consists of a grounded cathode stage with an impedance $Z_{1}$ connected between the grid and plate. Although the circuit now appears to belong to the PP family, this is not strictly the case, since the external feedback path has other than zero series impedance. Zero series impedance in feedback paths is a somewhat arbitrary limitation for the types of feedback circuits under discussion.

$$
\operatorname{PS}:\left\|y^{\|}\right\|=\left\|\begin{array}{cc}
\alpha^{\prime} & Y_{2}  \tag{19}\\
\alpha^{\prime}-Y_{1} & Y_{2}
\end{array}\right\|
$$

where

$$
\begin{equation*}
Y_{1}=1 / Z_{2}, Y_{2}=1 / Z_{2}, \quad \text { and } \quad \alpha^{\prime}=Y_{2}(1-k)+Y_{1} \tag{20}
\end{equation*}
$$

Examination of (10) through (13) shows that the expression $\alpha=Z_{1}(1-k)+Z_{2}$ enters into each equation in a rather fundamental way. It is not difficult in a practical experimental circuit to satisfy the condition $\alpha=0$ (which implies $\alpha^{\prime}=0$ ). It must be noted if $\alpha \leqq 0$ the circuit is potentially unstable, ${ }^{7}$ so that generally, $\alpha$ must be small but greater than 0 . If we neglect $\alpha$ we have, from (11) and (12)

$$
\begin{equation*}
\mathrm{SP}: Z_{\text {in }}=\frac{Z_{1} Z_{2}}{Z_{2}+Z_{L}} \tag{21}
\end{equation*}
$$

and

$$
\begin{equation*}
\mathrm{PS}: Z_{i n}=\frac{Z_{1}\left(Z_{2}+Z_{L}\right)}{Z_{L}} \tag{22}
\end{equation*}
$$

while the matrices (16) and (18) become

$$
\mathrm{SP}:\|z\|=\left\|\begin{array}{rr}
0 & -Z_{2}  \tag{23}\\
Z_{1} & Z_{2}
\end{array}\right\| \text { and }\|y\|=\left\|\begin{array}{rl}
Y_{1} & Y_{1} \\
-Y_{2} & 0
\end{array}\right\|
$$

and (17) and (19) become

$$
\text { PS: }\|z\|=\left\|\begin{array}{cc}
Z_{1} & -Z_{1}  \tag{24}\\
Z_{2} & 0
\end{array}\right\| \text { and }\|y\|=\left\|\begin{array}{cc}
0 & Y_{2} \\
-Y_{1} & Y_{2}
\end{array}\right\|
$$

If we consider the $\mathrm{SP} z$-matrix, we see that if an impedance $-Z_{2}$ is placed in series in the output circuit, we have a perfect impedance inverter, and if $Z_{1}=Z_{2}=R$, and $k=2$ (so that $\alpha=0$ ), we have

$$
\|z\|=\left\|\begin{array}{cc}
0 & -R  \tag{25}\\
R & 0
\end{array}\right\|
$$

which defines a gyrator. In addition, we can make a gyrator if an admittance $-Y_{1}$ is shunted across the input terminals, with $Y_{1}=Y_{2}=1 / R$, so that

$$
\|y\|=\left\|\begin{array}{cc}
0 & 1 / R  \tag{26}\\
-1 / R & 0
\end{array}\right\|
$$

which is the admittance matrix corresponding to (25). These circuits are shown in Figs. 3(b) and 3(c). In the same way, a gyrator can be made using the PS circuit by either shunting $-Y_{2}$ across the output terminals or by inserting $-Z_{1}$ in series with an input terminal, as shown in Figs. 3(d) and 3(e).

It is interesting to note that in both SP and I'S circuits, to correct for the residual $z_{11}$ or $z_{22}$ (which ever occurs), the series negative impedance is placed on the " P " or parallel feedback side, and the shunt negative admittance is placed on the " S " or series feedback side.

[^7]Thus, when $Z_{1}=Z_{2}=R$, and $k=2(\alpha=0)$, the SP and PS circuits can be regarded as identical except for a reversal in amplifier direction.

Since these circuits involve active elements, problems of stability are important. If we suppose that the terminations are passive, but otherwise arbitrary, the circuits must remain stable. In the first place, the condition that $\alpha=0$ is the boundary between stable and potentially unstable operation, as has been mentioned. In addition the negative resistances used to make the gyrators must be of the proper stability type. The series negative resistances must be open-circuit stable and the shunt conductances must be short-circuit stable. ${ }^{8}$ The range of impedance seen by the series negative resistance runs from $R$ to $\infty$, while the range of admittance seen by the shunt negative conductance runs from 0 to $1 / R$.

Experimental confirmation of one of these circuits was made, using the circuit of Fig. 3(e). The amplifiers used were of a design having accurately controlled input and output impedances ( 600 ohms) and were very stable. To obtain the negative resistance of -600 ohms, an SS circuit was used with $Z_{1}=Z_{2}=600 \mathrm{ohms}, Z_{L}=0$, and $k=3$ [see (10)]. The theoretical gain setting corresponding to $k=3$ is 9.54 db and the actual one used was 9.8 db .

A measure of the stability of the amplifiers used is given by the fact that the negative resistance obtained by the SS circuit referred to varied by less than 5 ohms from day to day. This resistance variation corresponds to a gain variation of approximately 0.04 decibel.


Fig. 4-Measurements of $\left|Z_{\text {in }}\right|$ vs $R_{L}$ of an impedance inverter circuit of PS type with series resistance $-R$ in input lead [Fig. 3(c)] $R=600$ ohms.

Measurements of $Z_{i n}$ versus $R_{L}$ at 400 cps were made using a Technology Instrument Company Type 310A $Z$-Angle Meter, the results of which are shown in Fig. 4. The inversion was very good over a two-decade range $40 \leqq R_{L} \leqq 4,000$ ohms. The angle of $Z_{\text {in }}$ was less than 10 degrees for $80 \leqq R_{L} \leqq 3,000$ ohms.

[^8]When $R_{L}$ was placed at the input, and $Z_{i n}$ measured at the output, the impedance was inverted, but the inversion range was reduced.

A 0.572 mfd condenser was used as $Z_{L}$ and the impedance $Z_{\text {in }}$ at 400 cps was measured as 535 / 90 degrees, whereas the computed impedance $\omega \mathrm{C}(600)^{2}=517$ ohms.

Computations were made of the fourpole impedances using the $Z_{i n}$ versus $R_{L}$ data discussed above. These impedances are

$$
\begin{aligned}
& z_{11}=-2.8+j 24.8 \mathrm{ohms} \\
& z_{12}=+598 \\
& z_{21}=-598
\end{aligned}
$$

$$
z_{22}=-4.1+j 13.8
$$

which indicate that the experimental circuit approaches an ideal gyrator quite closely.

The applications of these circuits to practical use is limited by the stability and accuracy requirements of the amplifiers employed. The closer the circuit approaches an ideal impedance inverter, the closer it approaches potential instability. For this reason, these circuits might prove useful in those cases in which the residual impedances of the inverter current would not require compensation by additional negative impedance circuits, or the requirements on $\alpha$ would not be too stringent.

# A Bridge for Measuring Audio-Frequency Transistor Parameters* 

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#### Abstract

Summary-A bridge is described which measures the smallsignal parameters of point contact and junction transistors at a frequency of 1 kc . The impedance parameters of point-contact transistors are measured for the grounded-base connection, while a set of parameters representing a compromise between the impedance parameters and the $h$ parameters is measured for junction transistors operating in either the grounded-base or grounded-emitter connections. This set includes the short-circuit input impedance $h_{11}$, the short circuit current amplification factor $h_{21}$, a paralleled resistance $r_{22}$ and capacitance $C_{22}$ representing the open-circuit output impedance, and two feedback resistances $r_{12}$ and $r_{12}{ }^{\prime}$ which in the particular case of the grounded-base connection represent the "lowfrequency" base resistance $r_{b}$ and the "high-frequency" base resistance $r_{b}{ }^{\prime}$ respectively. It is also shown that the $\alpha$ cut-off frequency of junction transistors can be calculated with good accuracy from the bridge measurements.


## Introduction

EQUIPMENT for measuring small-signal parameters is an important adjunct to development work on transistors and transistor circuits, and such equipment has been described by a number of atuthors. ${ }^{1-6}$ These authors have generally favored the use

[^9]of impedance parameters when dealing with point-contact transistors and either the admittance or the $h$ palrameters when dealing with junction transistors. The bridge which is described in the present paper has been designed, in accordance with the usual practice, to measure the impedance parameters of point-contact transistors. However, the set of parameters measured for junction transistors represents a compromise between the impedance and the $h$ parameters, and is discussed more fully in the next section. The bridge which operates at a frequency of 1 kc measures all resistive and reactive parameter components which play a significant part in the audio-frequency operation of transistors. Parameters are measured for the groundedbase connection of point-contact transistors and for the grounded-base or grounded-emitter connection of junction transistors. The other possible connections have not been considered because point-contact transistors are normally used only in the grounded-base connection for small-signal work, and grounded-collector measurements on a junction transistor do not appear to yield any more useful information than can be gotten from the other two connections. It has also been found possible to deduce a fairly accurate figure for the $\alpha$ cut-off frequency of alloyed-junction transistors from the measured phase lag of the parameter $\alpha_{c b}$ at 1 kc .

## Parameter Systems

Transistor small-signal parameters can be specified in a number of ways, depending on which of the four variables, input voltage ( $v_{1}$ ), input current ( $i_{1}$ ), output voltage ( $v_{2}$ ), and output current $\left(i_{2}\right)$, are chosen as independent variables. The combinations of $i_{1}$ with $i_{2}, i_{1}$ with


Fig. 1-Grounded-base equivalent $T$-circuits for a point-contact transistor.
$v_{2}$ and $v_{1}$ with $i_{2}$ as independent varimbles have all received extensive treatment in transistor literature, but in this paper attention will be restricted to the first two combinations which lead, respectively, to the impedance parameters and the $h$ parameters.

## Impedance Parameters

The well-known relationships which yield the impedance parameters are ${ }^{7}$

$$
\begin{align*}
& v_{1}=z_{11} i_{1}+z_{12} i_{2} .  \tag{1}\\
& v_{2}=z_{21} i_{1}+z_{22} i_{2} . \tag{2}
\end{align*}
$$

Eqs. (1) and (2) lead to the definitions

$$
\begin{align*}
& z_{11}=v_{1} / i_{1} \mid i_{2}=0  \tag{3}\\
& z_{12}=v_{1} / i_{2} \mid i_{1}=0  \tag{4}\\
& z_{21}=v_{2} / i_{1} \mid i_{2}=0  \tag{5}\\
& z_{22}=v_{2} / i_{2} \mid i_{1}=0 . \tag{6}
\end{align*}
$$

It is also convenient to introduce the current gain parameter $\alpha_{21}$ at this stage, although strictly speaking it belongs among the $h$ parameters. Here we write

$$
\begin{equation*}
\alpha_{21}=-i_{2} / i_{1} \mid v_{2}=0, \tag{7}
\end{equation*}
$$

and it may also be seen that

$$
\begin{equation*}
\alpha_{21}=z_{21} / z_{22} . \tag{8}
\end{equation*}
$$

In order to give the impedance parameters a physical basis it is customary to express them in terms of the parameters of the transistor equivalent $T$-circuit which for a point contact transistor can take either of the forms shown in Fig. 1(a) and (b). For a junction transistor operating in the grounded-base connection, either of the forms shown in Fig. 2(a) and (b) are applicable; while Fig. 3(a) and (b) apply to the grounded-emitter connection. In the audio-frequency range the parameters of point contact transistors are essentially resistive, and collector capacitance can be ignored. However, col-

[^10]

Fig. 2-Grounded-base equivalent $T$-circuits for a junction transistor.


Fig. 3-Grounded-emitter equivalent $T$-circuits for a junction transistor.
lector capacitance is included in the junction transistor circuits because its effect can become significant at frequencies of a few hundred cycles per second. A further important feature of these circuits is the division of the base resistance into two components $r_{b}{ }^{\prime}$ and $r_{b}{ }^{\prime \prime}$, where the component $r_{b}{ }^{\prime}$ results from ohmic resistance in the base region, and $r_{b}{ }^{\prime \prime}$ is an equivalent feedback resistance which results from "space-charge layer widening" as described by Early. ${ }^{8} r_{b}{ }^{\prime}$ is usually termed the "high-fre-
${ }^{8}$ J. M. Early, "Effects of space-charge layer widening in junction transistors," Proc. IRE, vol. 40, pp. 1401-1406; November, 1952.
quency base resistance" since it dominates the feedback at high frequencies when most of the collector signal current flows through $C_{c}$. In the grounded emitter equivalent circuit it is common practice to omit component of feedback resistance shown as $r_{b}{ }^{\prime \prime}(1-\alpha)$ in Figs. 3(a) and 3(b). However, it is included here because its presence can be detected with sensitive measuring equipment of the type described in this paper. Owing to the relatively rapid frequency variation of the factor ( $1-\alpha$ ) the parameters $\alpha /(1-\alpha), r_{c}(1-\alpha), C_{c} /(1-\alpha)$ and $r_{b}{ }^{\prime \prime}(1-\alpha)$ become frequency-dependent at a much lower frequency than other parameters not involving the factor $(1-\alpha)$. However, with average transistors the values measured at 1 kc are usually applicable up to high audio frequencies.

For a point-contact transistor or for a grounded-base junction transistor operating at frequencies so low that $C_{c}$ can be ignored, the impedance parameters are essentially resistive and are given by the simple relationships

$$
\begin{align*}
& z_{11}=r_{11}=r_{c}+r_{b}  \tag{9}\\
& z_{12}=r_{12}=r_{b}  \tag{10}\\
& z_{21}=r_{21}=r_{m}+r_{b} \approx r_{m}  \tag{11}\\
& z_{22}=r_{22}=r_{c}+r_{b} \approx r_{c} . \tag{12}
\end{align*}
$$

It may also be seen that

$$
\begin{equation*}
a_{21}=\alpha_{c 4}=r_{21} / r_{22} \approx \alpha \tag{13}
\end{equation*}
$$

On the other hand, when collector capacitance is taken into account the impedance parameters become complex quantities whose resistive and reactive components are all frequency depenclent. This makes the impedance parameters cumbersome to apply and it is desirable to seek some form of simplification. For the purposes of this paper, simplification has been achieved by using the equivalent circuit of Fig. 4 which employs a notation suitable for either the grounded-base or grounded-emitter connection.


Fig. 4-Four-pole equivalent circuit using modified impedance parameters.

This circuit is derived by considering first the operation at very low frequencies where only the parameters $r_{11}, r_{12}, r_{21}$, and $r_{22}$ play a part. It is then observed that so long as $1 / \omega C_{c} \gg r_{b}^{\prime}$ or $(1-\alpha) / \omega C_{c} \gg r_{c}{ }^{\prime}$ the collector capacitance can be taken into account by adding a capacitance $C_{22}$ across $r_{22}$ and by introducing an additional component of feedback voltage $r_{12}{ }^{\prime} i_{2}^{\prime \prime}$ resulting from the
flow of quadrative current through a resistance $r_{12}{ }^{\prime}$. Here $r_{12}{ }^{\prime}=r_{b}{ }^{\prime}$ for the grounded-base connection and $r_{12}{ }^{\prime}=r_{\mathrm{a}}{ }^{\prime}$ for the grounded emitter connection.

Summarizing the parameter values applicable to Fig. 4, we may write for the grounded base connection (adding the appropriate subscript $b$ )

$$
\begin{align*}
r_{11 b} & =r_{\epsilon}+r_{b}  \tag{14}\\
r_{12 b} & =r_{b}  \tag{15}\\
r_{12 b}^{\prime} & =r_{b}^{\prime}  \tag{16}\\
r_{21 b} & =r_{m}+r_{b} \approx r_{m}  \tag{17}\\
r_{22 b} & =r_{c}+r_{b} \approx r_{c}  \tag{18}\\
C_{22 b} & =C_{c} \tag{19}
\end{align*}
$$

For the grounded emitter connection

$$
\begin{align*}
r_{11 \epsilon} & =r_{\epsilon}+r_{b}  \tag{20}\\
r_{12 \mathrm{e}} & =r_{\epsilon}  \tag{21}\\
r_{12 \mathrm{e}}^{\prime} & =r_{\epsilon}^{\prime}  \tag{22}\\
r_{21 \mathrm{e}} & =r_{m}+r_{\epsilon} \approx r_{m}  \tag{23}\\
r_{22 \epsilon} & =r_{c}(1-\alpha)+r_{\mathrm{t}} \approx r_{c}(1-\alpha)  \tag{24}\\
C_{22 b} & =C_{c} /(1-\alpha) . \tag{25}
\end{align*}
$$

This simplification of junction-transistor impedance parameters is achieved at the expense of dealing separately with the in-phase and quadrature components of the electrode currents and voltages, but such procedures are standard in bridge measurements.

## $h$ Parameters

The use of $i_{1}$ and $v_{2}$ as independent variables leads to the relationships

$$
\begin{align*}
& v_{1}=h_{11} i_{1}+h_{12} v_{2}  \tag{26}\\
& i_{2}=h_{21} i_{1}+h_{22} v_{2} . \tag{27}
\end{align*}
$$



Fig. 5-Four-pole equivalent circuit using $h$ parameters.
These relationships correspond to the equivalent circuit of Fig. 5, and it may also be seen that

$$
\begin{align*}
& h_{11}=v_{1} / i_{1} \mid v_{2}=0  \tag{28}\\
& h_{12}=v_{1} / v_{2} \mid i_{1}=0  \tag{29}\\
& h_{21}=i_{2} / i_{1} \mid v_{2}=0  \tag{30}\\
& h_{22}=i_{2} / v_{2} \mid i_{1}=0 . \tag{31}
\end{align*}
$$

The $h$ parameters have the advantage of including the important current gain factor $h_{21}\left(=-\alpha_{11}\right)$ explicitly,
and when dealing with junction transistors they have the further advantage as compared with the impedance parameters of not requiring the measuring apparatus to provide an open-carcuit collector termination $i_{2 m 0}$. This latter requirement is a difficult although not insurmountable one, in view of the high collector resistance of junction transistors.

The parameters $h_{11}$ and $h_{21}$ are virtually unaffected by the presence of collector capacitance since they are measured with the collector shorted. For the groundedbase connection it may be shown that

$$
\begin{align*}
& h_{11 b}=r_{e}+r_{b}(1-\alpha)  \tag{32}\\
& h_{e 1 b}=-a_{c e} \approx-\alpha . \tag{33}
\end{align*}
$$

For the grounded-emitter connection

$$
\begin{align*}
& h_{11 e}=r_{b}+r_{c} /(1-\alpha)  \tag{34}\\
& h_{21 e}=-\alpha_{c b} \approx \alpha / 1-\alpha . \tag{35}
\end{align*}
$$

Thus these parameters are virtually pure resistances except in so far as they are affected by the frequency variation of $\alpha$.
The parameters $h_{12}$ and $h_{22}$ can be conveniently expressed in terms of the parameters $r_{12}, r_{12}{ }^{\prime}, r_{22}$, and $C_{22}$, viz.

$$
\begin{align*}
& h_{12}=r_{12} / r_{22}+j \omega C_{22} r_{12}{ }^{\prime}  \tag{36}\\
& h_{22}=1 / r_{22}+j \omega C_{22 .} . \tag{37}
\end{align*}
$$

In view of the general suitability of the $h$ parameters for application to junction transistors, the utility of the special treatment of impedance parameters shown in Fig. 4 may be questioned. However, the author is of the opinion that for general laboratory use it is more convenient to deal with the resistance $r_{22}$ than the conductance $1 / r_{22}$. Similarly the parameters $r_{12}$ and $r_{12}{ }^{\prime}$ are more readily visualized than the complex voltage ratio $h_{12}$. Consequently a compromise proposal has been adopted here for junction transistors in which the parameters $h_{11}$ and $h_{21}$ are directly measured, but instead of $h_{12}$ and $h_{22}$ the quantities $r_{12}, F_{12}{ }^{\prime}, r_{22}$, and $C_{22}$ are measured. Where the complete set of $h$ parameters is needed, as for instance in network calculations, the values of $h_{12}$ and $h_{22}$ can be easily calculated from (36) and (37).

The bridge if necessary could be modified to read $h_{22}$ and $h_{12}$ directly, but so far the need has not seemed great enough to warrant the additional circuit complications. For point-contact transistors the resistance parameters and the $\alpha$ (i.e., $h_{21}$ ) parameter are normally measured although $h_{11}$ can be measured, if necessary, so long as its value is positive.

## Estimation of $\alpha$ Cut-off Frequency

It may be shown that the frequency variation of $\alpha$ (i.e., $\alpha_{c e}$ ) can be described by the relationship

$$
\begin{equation*}
\alpha=\frac{\alpha_{0}}{1+j \kappa \omega / \omega_{\alpha}}, \quad \omega \ll \omega_{\alpha}, \tag{38}
\end{equation*}
$$

where $\alpha_{0}=$ low frequency value of $\alpha, \omega_{\alpha}=$ cut-off angular frequency, i.e., the frequency at which $|\alpha| \alpha_{0} \mid=0.707$, and $\kappa$ is a constant whose value is determined by the nature of the physical process giving rise to the $\alpha$ cutoff. Usually the cutoff is determined by the dispersion in transit time of minority carriers crossing the base, in which case $\kappa$ has a theoretical value of 1.21 .

At low frequencies, (38) may be rewritten with good accuracy as

$$
\begin{equation*}
\alpha \approx \alpha_{0}\left(1-j \kappa \omega / \omega_{\alpha}\right), \quad \omega \ll \omega_{\alpha} . \tag{39}
\end{equation*}
$$

This relationship in principle offers the possibility of calculating $\omega_{\alpha}$ by measuring the phase lag $\tan ^{-1} \kappa \omega / \omega_{\alpha}$ at a known low frequency $\omega$. In the case of the present bridge operating at 1 kc , the phase angle is too small to measure with any degree of accuracy. However, the phase angle of $\alpha_{c b}$ is much larger than that of $\alpha_{c e}$, as may be seen from the following analysis:

$$
\begin{align*}
\boldsymbol{\alpha}_{c b} & =\frac{\alpha}{1-\alpha} \cdot \\
& =\frac{a_{0} /\left(1+j \kappa \omega / \omega_{\alpha}\right)}{1-\alpha_{0} /\left(1+j \kappa \omega / \omega_{\alpha}\right)}, \\
& =\frac{\alpha_{0}}{1-\alpha_{0}+j \kappa \omega / \omega_{\alpha}}  \tag{40}\\
& \approx \frac{\alpha_{0}}{1-\alpha_{0}}\left\{1-j \kappa \omega / \omega_{\alpha}\left(1-\alpha_{0}\right)\right\}, \omega \ll \omega_{\alpha}\left(1-\alpha_{0}\right) . \tag{41}
\end{align*}
$$

This shows that $\alpha_{c b}$ has a phase lag of $\tan ^{-1} \kappa \omega$ $/ \omega_{\alpha}\left(1-\alpha_{0}\right)$, which is very much greater than that of $\alpha_{c e}$, owing to the presence of the factor $1 / 1-\alpha_{0}$.

It has been found possible to measure this phase lag with good accuracy at 1 kc , and the values of $\omega_{\alpha}$ calculated from the phase lag with $\kappa=1.21$ are found to agree fairly closely with the directly-measured values of $\omega_{\alpha}$.


Fig. 6-Equivalent circuit for $h_{111}$.
Mueller and Pankove ${ }^{9}$ using equipment developed by Giacoletto ${ }^{8}$ have described an alternative method of calculating $\omega_{\alpha}$ from bridge measurements on junction transistors. They show that the parameter $h_{11}$ can be represented by the network of Fig. 6, where the resistance $r_{b b}$ originates in the base layer ohmic resistance, and the capacitance $C_{b^{\prime},}$, and the resistance $r_{b^{\prime},}$ result from diffusion of minority carriers in the base. The value

[^11]of $C_{b^{\prime}}$, is given as $q I \epsilon / k T \omega_{\alpha}$, where $q / k T=39$ volt $^{-1}$ at room temperature. Hence when $C_{b^{\prime}, ~}$ is known $\omega_{a}$ can be calculated.

The bridge described here is capable of measuring an equivalent capacitative component of $h_{11}$ referred to the base terminal, and the value of $C_{b^{\prime}}$, could then be calculated if $r_{b b^{\prime}}$ were known. $r_{b b^{\prime}}$ tends to have a value close to that of $r_{b}^{\prime}$, but it would appear that there can be appreciable differences between these two parameters in the case of alloyed-junction transistors. As shown by Giacoletto, the values of $r_{b b^{\prime}}, C_{b^{\prime}, \text { and }} r_{b^{\prime}, ~ c a n ~ o n l y ~ b e ~}^{\text {co }}$ determined accurately by a multifrequency test set, and it must therefore be concluded that with the present bridge only a rough value of $\omega_{\alpha}$ can be deduced from $h_{11}$ e measurements. This has been verified by actual measurements.

## Principle of Measurement

The bridge circuit is derived from the one used in the General Radio vacuum tube bridge and described by Tuttle. ${ }^{10}$ The adopted principle makes the measurement of $h_{21}$ and $h_{11}$ interdependent. Similarly the measurement of $r_{22}$ and $C_{22}$ is linked with the measurement of $r_{12}$ and $r_{12}{ }^{\prime}$, while the measurement of $r_{11}$ is linked with the measurement of $r_{21}$.

## Measurement of $h_{21}$ and $h_{11}$

Fig. 7 shows the basic circuit for measuring these parameters. In this and the following circuits the voltages $e_{1}, e_{2}$, and $e_{3}$ are small $1,000 \mathrm{cps}$ voltages derived from low-impedance attenuators connected to a number of secondary windings on an input transformer, the in-
$e_{3}$. For reasons which will become apparent later, it is convenient to make $k_{2}$ and $k_{3}$ continuously adjustable while $k_{1}$ is adjustable by factors of ten, i.e., it takes the values $1.0,0.1,0.01$ etc. It will be noticed that the voltage $e_{2}$ appears twice in Fig. 7. This is arranged by the use of a 1:1 transformer.

The resistance $R_{1}$ is made at least a hundred times greater than the input impedance so that the input current can be taken as $i_{1}=e_{1} / R_{1}$ within 1 per cent accuracy. Emitter or base bias current is obtained by connecting a suitably by-passed de supply to the "low" end of $e_{1}$. Collector bias is applied through the primary of the detector transformer $T_{1}$. The latter is followed by a highgain detector amplifier.

The polarity of the voltages shown in Fig. 7 is suitable for positive values of $h_{21}$, i.e., for the grounded emitter connection. For the grounded base connection, $e_{2}$ must be reversed.

For a detector null to be obtained $\left(i_{D}=0\right)$ it is seen that

$$
\begin{equation*}
i_{2}=i_{R}+i_{Q} . \tag{45}
\end{equation*}
$$

At the same time $v_{2}=0$ so that

$$
\begin{equation*}
i_{2}=h_{21} i_{1}=h_{21} e_{1} / R_{1} . \tag{46}
\end{equation*}
$$

Also $i_{R}=e_{2} / R_{2}$ and $i_{Q}=-j \omega C e_{2}$, whence

$$
\begin{align*}
h_{21} & =\frac{R_{1} e_{2}}{R_{2} e_{1}}\left(1-j \omega C R_{2}\right)  \tag{47}\\
& =\frac{R_{1} k_{2}}{R_{2} k_{1}}\left(1-j \omega C R_{2}\right) . \tag{48}
\end{align*}
$$



Fig. 7-Masic circuit for measuring $h_{21}$ and $h_{11}$.
put voltages to the attenuators being carefully equalized. The voltages $e_{1}$ and $e_{2}$ are exactly in phase with each other but the voltage $e_{3}$ has an adjustable phase angle. Accordingly we may write

$$
\begin{align*}
& e_{1}=k_{1} e_{0}  \tag{42}\\
& e_{2}=k_{2} e_{0}  \tag{4.3}\\
& e_{3}=k_{3}(1+j \beta) e_{0}, \tag{44}
\end{align*}
$$

where $k_{1}, k_{2}$, and $k_{3}$ are attenuation factors, $e_{0}$ is the attenuator input voltage, and $\tan ^{-1} \beta$ is the phase angle of

[^12]Thus the real part of $h_{21}$ can be read directly from the setting of the $k_{2}$ attenuator multiplied by the factor $R_{1} / R_{2} k_{1}$ which can be arranged to be a power of ten. For the grounded base connection the real part of $h_{21}$, i.e., $\alpha_{c e}$, is the only significant term.

In the grounded emitter case it was shown in (41) that $h_{21}$, i.e., $\alpha_{c b}$, has a significant phase lag equal to $\tan ^{-1} \kappa \omega / \omega_{\alpha}\left(1-\alpha_{0}\right)$. Hence equating this to the phase lag of (48) yields

$$
\begin{equation*}
\kappa \omega / \omega_{\alpha}\left(1-\alpha_{0}\right)=\omega C R_{2} \tag{49}
\end{equation*}
$$

whence

$$
\begin{equation*}
\omega_{\alpha}=\kappa /\left(1-\alpha_{0}\right) C R_{2} . \tag{50}
\end{equation*}
$$

Once a null balance has been established for the $h_{21}$ measurement the transistor input impedance is by definition $h_{11}$. A voltage $h_{11} i_{1}$ is then developed at the transistor input terminal and this can be measured by closing $S_{1}$ and adjusting $e_{3}$ to restore the null balance. When this is done we have

$$
\begin{equation*}
e_{3}=h_{11} i_{1} \tag{51}
\end{equation*}
$$

or

$$
\begin{align*}
h_{11} & =R_{1} e_{3} / e_{1}  \tag{52}\\
& =R_{1} k_{3}(1+j \beta) / k_{1} . \tag{53}
\end{align*}
$$

Thus the real part of $h_{11}$ is given directly by the attenuating factor $k_{3}$ multiplied by the factor $R_{1} / k_{1}$ which can be arranged to be a power of ten. For the grounded base connection, $h_{16}$ is essentially resistive and $\beta$ is quite small. ${ }^{11}$ As mentioned earlier, $h_{11 \epsilon}$ has a substantial capacitative component so that negative values of $\beta$ must be provided to obtain a balance.

## Measurement of $r_{22}, C_{22}, r_{12}$ and $r_{12}{ }^{\prime}$

The method of measurement is illustrated in Fig. 8. It will be appreciated that for point-contact transistors a simpler analysis applies in which $C_{22}$ and $r_{12}{ }^{\prime}$ are ignored.


Fig. 8-Basic circuit for measuring $r_{22}, C_{22}, r_{12}$, and $r_{12}{ }^{\prime}$.
Here the resistance $R_{1}$, through which the input bias current is fed, is made large enough to represent an effective ac open circuit at the input terminal. When a null is achieved with this circuit the collector voltage must be equal to $e_{2}$ and $i_{D}=0$ so that

$$
\begin{equation*}
i_{R}+i_{Q}=i_{2}^{\prime}+i_{2}^{\prime \prime}, \tag{54}
\end{equation*}
$$

whence

$$
\begin{equation*}
e_{1} / R_{2}+j \omega C e_{2}=e_{2} / r_{22}+j \omega C_{22} e_{2} . \tag{55}
\end{equation*}
$$

Equating in-phase and quadrature components in this equation yields

$$
\begin{equation*}
r_{22}=R_{2} e_{2} / e_{1}=R_{2} k_{2} / k_{1} \tag{56}
\end{equation*}
$$

and

$$
\begin{equation*}
C_{22}=C . \tag{57}
\end{equation*}
$$

[^13]This shows that $r_{22}$ can be read from the $k_{2}$ attenuat or multiplied by the factor $R_{2} / k_{1}$. n addition, $C_{22}$ is obtained by calibrating the variable capacitor $C$.

The foregoing analysis has neglected the effect of stray capacitance. It is found in practice that only the stray capacitance $C_{\text {s }}$ shown dotted in Fig. 8 has any significant effect. This capacitance appears at a high impedance point and, in effect, augments the collector capacitance. Stray capacitances at other points are shunted across low impedances and can be ignored. The magnitude of $C_{s}$ can be determined by placing a resistance of known self-capacitance (determined, say, by a $Q$-meter) between the collector terminal and ground. The setting of the quadrature balancing capacitor $C$ when the bridge is in balance then gives a measure of $C_{s}$ which can be deducted from subsequent measurements with a transistor in circuit. It is to be noted that $C_{s}$ does not influence the $\alpha$ measuring circuit since the collector ac voltage is then zero.

In measuring $r_{12}$ and $r_{12}^{\prime}$ a mull is first established in respect of $r_{22}$ and $C_{22} . S_{1}$ is then closed and $e_{3}$ is adjusted in magnitude and phase to restore the null. In this case

$$
\begin{align*}
e_{3} & =r_{12} i_{2}^{\prime}+r_{12}^{\prime} i_{2}^{\prime \prime}  \tag{58}\\
& =r_{12} e_{1} / R_{2}+j \omega C_{22} e_{2} r_{12}^{\prime}, \tag{59}
\end{align*}
$$

whence

$$
\begin{equation*}
k_{3}(1+j \beta)=r_{12} k_{1} / R_{2}+j \omega C_{22} k_{2} r_{12}{ }^{\prime} . \tag{60}
\end{equation*}
$$

Equating in-phase and quadrature components yields

$$
\begin{align*}
r_{12} & =R_{2} k_{3} / k_{1}  \tag{61}\\
r_{12}^{\prime} & =\beta k_{3} / \omega C_{22} k_{2}, \tag{62}
\end{align*}
$$

or, alternatively,

$$
\begin{equation*}
r_{12}^{\prime} / r_{12}=\beta / \omega C_{22} r_{22} . \tag{63}
\end{equation*}
$$

In this case $r_{12}$ can be read directly from the $k_{3}$ attenuator multiplied by the factor $R_{2} / k_{1}$. On the other hand, the value of $r_{12}{ }^{\prime}$ is not directly indicated but must be calculated from (62) or (63).


Fig. 9-Basic circuit for measuring $r_{21}$ and $r_{11}$.

## Measurement of $r_{21}$ and $r_{11}$

The basic circuit for measuring these parameters is shown in Fig. 9 and is illustrated for the sake of generality with regard to a junction transistor, although such measurements are usually only performed on point-
contact transistors. Once again, $R_{1}$ is made very much greater than $r_{11}$ so that $i_{1}=e_{1} / R_{1}$. When a null is established on the collector side, the in-phase component of collector current is reduced to zero; $i_{2}{ }^{\prime}=0$. Hence

$$
\begin{equation*}
r_{21} i_{1}=e_{2} \tag{64}
\end{equation*}
$$

and

$$
\begin{equation*}
i_{2}^{\prime \prime}=j \omega C_{22} e_{2}=j \omega C e_{2}, \tag{65}
\end{equation*}
$$

whence

$$
\begin{equation*}
r_{21}=e_{2} R_{1} / e_{1}=k_{2} R_{1} / k_{1} \tag{66}
\end{equation*}
$$

and

$$
\begin{equation*}
C_{22}=C . \tag{67}
\end{equation*}
$$

Thus $r_{21}$ is read from the bridge in very much the same way as $r_{22}$, and the capacitance balancing conditions, including the effects of $C_{s}$, are identical with those discussed in the previous section.

In measuring $r_{11}$, the null is first established on the collector side, $S_{1}$ is closed and $e_{3}$ is adjusted to restore the null balance. Since $i_{2}{ }^{\prime}=0$ in this case,

$$
\begin{equation*}
e_{3}=r_{11} i_{1}+r_{12}{ }^{\prime} i_{2}^{\prime \prime} \tag{68}
\end{equation*}
$$

or

$$
\begin{equation*}
k_{3}(1+j \beta) e_{0}=r_{11} k_{1} e_{0} / R_{1}+j \omega C_{22 K_{2} e_{0} r_{12} 2^{\prime}} \tag{69}
\end{equation*}
$$

Whence, equating the in-phase terms,

$$
\begin{equation*}
r_{11}=\kappa_{3} R_{1} / \kappa_{1} . \tag{70}
\end{equation*}
$$

Thus $r_{11}$ is read with the same scale factors as $r_{12}$.
It may also be seen that the quadrature terms of (69) offer the possibility of measuring $r_{12}{ }^{\prime}$, but since this has already been done it need not be considered any further.

In all three basic bridge circuits there is an element of indirectness in the measurements carried out on the input side of the transistor, since they depend on a prior balance being performed on the output side. However this method has the advantage that in Fig. 9 the residual in-phase component of the collector current can be made as low as $10^{-10} \mathrm{amps}$, i.e., comparable with the transistor noise current. This is equivalent to maintaining an effective ac impedance of thousands of megohms at the collector, a value which would be impossible to achieve with more conventional methods. Furthermore, the method of Fig. 8 allows accurately known currents as low as $10^{-8} \mathrm{amps}$ to be passed into the collector terminal for the purpose of measuring $r_{12}$ and $r_{12}{ }^{\prime}$. Other more direct methods of carrying out measurements on the input side of the transistor would require the detector to be transferred to that side. The present arrangement which keeps the detector on the output side results in a minimum of switching and allows the power gain of the transistor to be used to improve the balance sensitivity.

## Complete Circuit

The complete circuit embodying the basic circuits of Figs. 7, 8 and 9 is shown in Fig. 10 (opposite). Alternat-
ing current at 1 kc is applied to the bridge attenuators through a transformer $T_{1}$ having a $20: 1$ step-down ratio between the primary and each secondary. The primary is driven by a cathode-follower in order to keep the secondary output impedance down to a few ohms. The secondary voltage is adjustable up to a maximum value of about 1 volt. This transformer should be designed to have equal secondary leakage reactances so that the secondary voltages are exactly in phase. In the experimental transformer used here the middle secondary was found to have a slightly higher leakage reactance than the outer two, and the loading arrangement shown in Fig. 10 which places approximately 1,100 ohms on the middle secondary and approximately 100 ohms on the outer secondaries was found necessary in order to bring the secondary currents exactly into phase with each other. Finally the small trimming resistors $R_{5}$ and $R_{10}$ were added to make the full-scale attenuator voltages equal within $\pm 0.1$ per cent. These adjustments were carried out by performing an $\alpha$ measurement with the emitter and collector terminals shorted together. In this case if $R_{1}$ and $R_{2}$ are carefully equalized and the quadrature balancing capacitor is disconnected, $e_{1}$ and $e_{2}$ must be exactly equal in amplitude and phase in order to obtain a null. The bridge should then, of course, register an $\alpha$ of unity. A check on $e_{3}$ can then be made by connecting it temporarily in place of $e_{1}$.

The attenuator controlling $e_{1}$ takes the form of a simple voltage divider which yields values of $k_{1}$ equal to $1.0,0.1,0.01$, or 0.001 . Its output impedance is 100 ohms or less depending on the setting used. The attenuator controlling $e_{2}$ has a coarse and fine control which give a setting accuracy of $\pm 0.1$ per cent of the full-scale value. Nine 10 -ohm resistors $R_{3, A} \ldots I$ and the 10.5 ohm potentiometer $R_{4}$ are used here, the value of 10.5 ohms being adopted to give a small overlap between the coarse steps. For the sake of smoothness $R_{4}$ has been constructed in the form of a single-turn slide wire potentiometer. A switch $S 7$ is provided to reverse the polarity of $e_{2}$ wherever required.

The method of collector capacitance balancing shown in Fig. 10 represents an improvement over the simple scheme described earlier and follows the technique adopted in the General Radio bridge. It will be seen that two voltages $\pm e_{2}$ are generated with the aid of the center-tapped autotransformer $T_{2}{ }^{12}$ and are applied to the stators of the differential variable capacitor $C_{1}$ whose rotor is connected to the collector terminal. With this arrangement the effective balancing capacitance can take positive or negative values and can be varied smoothly through zero effective value. This avoids the limitations which would be imposed by the minimum capacitance of a single-ended variable capacitor. Often it is found that the built-in capacitor is not large enough to secure a balance, and in this case an external

[^14]

Fig. 10-Complete circuit of transistor bridge.
variable capacitor is connected to the terminals marked "Ext. Cap. $C_{1}$." Alternatively the external capacitor can be used to obtain greater precision than the internal capacitor is capable of providing.

The attenuator controlling $e_{3}$ includes the coarse and fine potentiometer system $R_{15 A} \cdots R_{15 I}$ and $R_{16}$
which are identical with $R_{3 A} \cdots R_{3 I}$ and $R_{4}$. These are preceded by a network including $R_{12}, R_{13}, R_{14}$ and $S_{9}$ which is arranged to provide two ranges of $\kappa_{3}$, namely 0 to 0.01 and 0 to 0.001 . When the phase angle of $e_{3}$ requires adjusting, an external decade capacitor is introduced in the position marked $C_{2}$ or $C_{3} . C_{2}$ is used when
the phase angle of $e_{3}$ must be advanced, as when measuring $r_{12}{ }^{\prime}$, and $C_{3}$ is used when measuring $h_{11 \epsilon}$. Analysis shows that the magnitude of the quantity $\beta$ appearing in (44) is given to good approximation by relationship

$$
\begin{equation*}
\beta=\omega C_{2} R_{12} \quad(|\beta|<1) \tag{71}
\end{equation*}
$$

In this case (63) can be rearranged to the form,

$$
\begin{equation*}
r_{12}^{\prime} / r_{12}=R_{12} C_{2} / r_{22} C_{22} \tag{72}
\end{equation*}
$$

Further analysis shows that for $|\beta|>1$ (which occasionally occurs) the value of $r_{12}$, as read from the bridge, should be increased by a factor $\left(1+.01 \beta^{2}\right)$. This is because the presence of $C_{2}$ causes the in-phase component of $e_{3}$ to increase slightly. Notice that this arrangement for shifting the phase angle of $e_{3}$ is an improvisation which might well be replaced by some other method giving a more direct indication of the value of $r_{12}{ }^{\prime}$.

The voltage $e_{3}$ is coupled to the transistor input terminal through the push button switch $S_{1}$ and a $1 \mu F$ capacitor $C_{6}$ which is shunted by a 1 -megohm resistor in order to prevent the accidental accumulation of charge on $C_{6}$.

Three values are provided for the resistance $R_{1}$, namely 1 megohm, 100 kilohms and 10 kilohms. The first value is used when working with small input bias currents, when at the same time the transistor input impedance tends to be high. The second value is used for moderate values of bias current, i.e., up to 3 ma with a 300 -volt de supply. The third value is used for occasional measurements at higher bias currents. With this value of $R_{1}$ it may be necessary to make corrections to the measured parameter value because the transistor input impedance may be greater than $R_{1} / 100$.

For the resistance $R_{2}$ alternative values of $100 \mathrm{k} \Omega$ or $10 \kappa \Omega$ are provided, the former value belng most frequently used. The 10,000 -ohm value sometimes permits a higher signal current level to be achieved with a consequent sharpening of the null point.

The detector transformer is specially wound with a double-shielded primary, the inner shields being connected to the "low" end of the primary to secure a guard-ring effect. The primary is parallel resonated by the capacitor $C_{5}$ at 1 kc to improve the detection sensitivity and at the same time to provide some discrimination against hum pickup and transistor noise. By airgapping the transformer core, the effects of dc saturation are minimized and the $Q$ at 1 kc is improved. Induced voltages resulting from stray fields are cut down below the limits of detection by proper orientation of the transformer and by enclosing it in a mumetal shield.

Fig. 11 is a photograph showing the layout of controls on the panel of the bridge.

## Performance Data

With the component values shown in Fig. 10 values of $\alpha$ up to 1,000 and values of $r_{22}$ or $r_{21}$ up to 100 meg ohms can be measured. The range of values of $h_{11}, r_{11}$, and $r_{12}$ which can be measured depends on the type of
transistor and the connection being used but is adequate for all normal purposes. Experience has shown, however, that it would be advantageous to have an additional attenuating position on $S_{9}$ yielding a $0-0.0001$ range of values of $\kappa_{3}$. This would improve the accuracy when measuring small values of $r_{12}$ in association with large values of $r_{22}$ (i.e., $h_{12}$ very small).


Fig. 11-Photograph of transistor bridge.
In checking the accuracy of the bridge, certain tests can be carried out such as an $\alpha$ measurement with the emitter and collector terminals shorted which, as mentioned earlier, should give an indicated $\alpha$ of unity. A comprehensive check on $r_{22}$ and $C_{22}$ measurements can be carried out with the aid of precision resistors and capacitors, and a similar check can be devised for $r_{11}$ and $h_{11}$ measurements by connecting a precision resistor between the emitter and base terminals and using a vacuum tube to supply an amplifying link between the emitter and collector terminals. These tests show that where the transistor noise level is low enough to yield a sharp null point the resistance measurements should be accurate within 1 or 2 per cent except where the parameters take extremely high or low values. $C_{\varepsilon 2}$ measurements are most accurate when associated with high values of $r_{22}$, in which case an accuracy of the order of $\pm 2$ per cent $\pm 1 p F$ is obtainable. High values of $r_{22}$ need to be corrected for the inherent leakage resistance of the bridge which can be measured with the collector terminal open-circuited.

The bridge is capable of measuring parameters over a wide range of biasing conditions extending down to as little as 0.1 volt collector bias and $1 \mu a$ emitter current.

Under these conditions special precautions must be taken to ensure that signal voltages and currents are small in comparison with bias values without at the same time allowing the signal to be obscured by noise. The best compronise is achieved by increasing the signal level to the point at which the null setting just begins to shift owing to curvature of the transistor characteristics.

The correlation between the $\alpha$ cut-off frequency as calculated from the bridge measurements and the directly measured value has been found to be best with alloyed-junction transistors. In this regard a test was carried out with 24 alloyed-junction transistors, mostly type OC70, OC71, and some CK722's. The value of $\kappa$ calculated from (50) as an empirical constant to harmonize the measured $\alpha$ cut-off frequency with the $1-\mathrm{kc}$ phase measurements worked out as $\kappa=1.17$ with a standard deviation of $\pm 0.05$. The difference between the measured value of 1.17 and the theoretical value of 1.21 is not considered significant in view of the spread of measurements and the possibility of small systematic errors.

In the case of grown-junction transistors, of which relatively few have been available to the author, apparrent values of $\kappa$ ranging from 0.7 to 1.4 have been deter-
mined; indicating that the factors determining the cut-off frequency of these transistors are more complex than those applying to alloyed-junction transistors. On the whole it is felt that, in the absence of an " $\alpha$ sweeper" or other convenient $\alpha$ cut-off measuring equipment, the feature of being able to calculate from the bridge measurements a fairly accurate $\alpha$ cut-off frequency for alloyed-junction transistors should prove quite useful.

Values of $r_{b}^{\prime}$ measured on the bridge are found to fall within $\pm 10$ per cent of the values measured directly at high frequencies and in this regard it may be of interest to note that with grown-junction transistors, values of $r_{b}{ }^{\prime}$ greater than $r_{b}$ are often measured. Here the value of $r_{b}{ }^{\prime \prime}$ is negative, an effect which Early ${ }^{13}$ has shown to be due to "base-resistance modulation." In alloyed-junction transistors $r_{b}$ is always several times greater than $r_{b}{ }^{\prime}$ for moderate values of emitter current. A corollary of these remarks in regard to $r_{b}$ and $r_{b}{ }^{\prime}$ is that $r_{t}{ }^{\prime}$ is greater than $r_{e}$ except when $r_{b}{ }^{\prime \prime}$ is negative.

## Acknowledgment

Thanks are due to Mr. C. D. Howarth, who attended to the constructional details of the bridge.
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# Skin Resistance of a Transmission-Line Conductor of Polygon Cross Section* 

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#### Abstract

Summary-If a conductor cross section is any straight-sided polygon that can be circumscribed on a circle, it is found to have the same skin resistance as a conductor whose cross section is this circle. For example, a square wire has the same resistance as a round wire of the same radius, though the square perimeter is $4 / \pi$ times as great. This "polygon rule" is derived from the "incrementalinductance rule" of the skin effect, published in 1942. It applies equally to inner or outer conductors, though the current distribution is very different. It applies to some unusual shapes that are difficult to compute by any other method. $\ddagger$


T- HE SKIN EFFEC'1 is the well-known phenomenon of high-frequency currents concentrating just under the exposed surfaces of a conductor. The effective resistance of a cylindrical conductor may be computed by integrating the power losses of such cur-

[^15]rents associated with its surfaces. This has been done for cylindrical conductors of some simple cross sections, and more laboriously for various polygon cross sections.

The theorem to be presented gives a very simple evaluation of this resistance for a certain category of polygon cross sections. It applies to all polygons that can be circumscribed on a circle. Two conductors of the same material have the same skin resistance if their cross sections are respectively a circle and a polygon of such size and shape that it can be circumscribed on this circle. This rule applies to inner or outer conductors bounding a high-frequency magnetic field. A special case of the former is an antenna wire. The conditions of validity will be specified in more detail.

It is essential to distinguish between the skin resistance and the inductive reactance of a conductor, because they are obtained by different methods of computation. Conformal mapping can be utilized to obtain directly, for different shapes, the ratio of sizes that will have the same reactance. To this operation must be added the integration of power dissipation over the
surface, to obtain the ratio of sizes that will have the same resistance. The present theorem is an implicit solution for the latter in a variety of cases.

For example, a square wire has the same skin resistance as a circular wire of the same diameter but only $\pi / 4$ as great a perimeter. This is illustrated in Fig. 1, as a basis for stating the theorem and its proof.

> Skin Depth: $d=\frac{1}{\sqrt{\pi f \mu_{0} \sigma}}$
> Surface Rcsistivity: $R_{s}=\rho / d=\sqrt{\pi f \mu_{0} \rho}$
> Skin Kesistancc: $\quad R=R_{N} l / 2 \pi r$

Fig. 1-Cross-sectional circle and circumscribed polygon (square) having equal resistance.

The underlying principle is the "incremental-inductance rule" previously stated by the author. ${ }^{2,3}$ Its basis will be restated briefly.

In the idealized skin effect, the current appears to be distributecl uniformly in the skin depth (d) just under the surface of the conductor. The average depth of the current is then half the skin depth ( $d / 2$ ).

In a perfect condluctor, current would travel on the surface, since there would be zero penetration. This is the usual assumption in inductance formulas for high frequencies; coaxial line is the most common example.

The depth of penetration effects a proportional increment of inductance, just as if the conductor surface receded by an amount equal to the half-depth ( $d / 2$ ). It happens that the reactance of this increment of inductance is equal to the skin resistance, so the latter may be evaluated by computing the former. This is the "incremental-inductance rule."

Referring to Fig. 1, there are shown two cross sections of the same radius, a circle and a circumscribed square. These are given as examples of all the polygons that can be circumscribed on a circle. We wish to compare these two shapes with respect to the increment of inductance caused by reducing the radius ( $r$ ) by a relatively small amount ( $\Delta r \ll r$ ). In each case, this change of size is accompanied by no change of shape.

Here we rely on the well known principle that the change of inductance depends only on the ratio of change of size, if the shape remains the same. The simplest example of this principle is the coaxial line, whose inductance depends only on the shape (determined by the ratio of radii) and not on the size. In Fig. 1, either cross section may represent the inner conductor of a coaxial line whose outer conductor is represented by a concentric circle of much greater radius. If then the radii of both inner and outer conductors are reduced
in the same relative amount $(\Delta r / r)$, the inductance remains the same. This proves that the change of the inductance of either inner conductor is the same as that of the outer conductor; hence the circle and the circumscribed polygon have the same change of inductance.

From the preceding relationship and the incrementalinductance rule, it is deduced that the circle and the circumscribed polygon have the same skin resistance. This theorem is designated, the "polygon rule."

The actual skin resistance may be evaluated by the well-known formulas in Fig. 1; the symbols are defined as follows (MKS rationalized units):

```
    \(r=\) radius of circular cross section of conductor
\(\Delta r=\) effective reduction of radius by penetration
    \(\rho=\) resistivity of conductor
    \(\sigma=\) conductivity of conductor
\(\mu_{0}=\) magnetivity of free space (also in the conductor,
        assumed nonmagnetic)
    \(f=\) frequency
    \(l=\) length of conductor
    \(d=\) skin depth
\(R_{s}=\) surface resistance of a square (equal length and
        width)
    \(R=\) skin resistance of a conductor (of specified radius
        and length).
```

As an example of the polygon rule applied to inner and outer conductors, Fig. 2 shows a coaxial line of square cross section. The squares are circumscribed on the circles shown in dotted lines. By the same reasoning, either square has the same skin resistance as its inscribed circle. (With respect to reactance, the pair of squares gives less than the pair of circles, and occupies more space.) Either square may be used with the other circle, since the polygon rule applies independently to each of the two conductors.


Fig. 2-Concentric circles, and circumscribed squares having equal resistance (but less reactance in greater space).

In this and further cases, the equality of resistance for the circle and the circumscribed polygon is realized to the extent that the following conditions are approximated. (The first of these is the usual assumption for the simple formulas of the skin effect.)

## Conditions

1. The skin depth is much less than $\frac{2}{4}$ the thickness of any substantial part of the conductor. (At the angles of a polygon, this condition is met if the skin depth is a very small fraction of the radius of the inscribed circle, as is true in many applications.)


Fig. 3-Regular polygons circumscribed on equal circles, examples of inner or outer conductors.
2. Outside of an "inner" condluctor (or isolated wire) any other conductors affecting the field pattern are spaced at a distance sufficient to provide that they have a negligible effect on the current distribution.
3. Inside of an "outer" condluctor, the inner conductor is near the center of the inscribed circle and is small enough to provide that the current distribution on the outer conductor is substantially the same as would be obtained with a fine wire in the center.

Commenting on the second and third conditions applied to a coaxial line, the polygon rule requires that the ratio of radii be great enough to prevent either conductor causing appreciable distortion of the field at the other conductor. The ratio of radii may be closer to unity as either or both of the conductors assumes a higher order of symmetry or regularity.

Figs. 3 and 4 show examples of regular polygons and elongated polygons circumscribed on equal circles and therefore having the same values of skin resistance. Every one is closed and hence may serve as inner or outer conductor; if the latter, the associated inner conductor is centered in the circle.

Fig. 5 shows examples of right polygons that have some sides open. Every open conductor is assumed to extend outward as far as the field is appreciable. Therefore these are suitable only for outer conductors.
(•)

(b)

(c)

(d)

(e)


Fig. 4-Elongated polygons circumscribed on equal circles, examples of imer or outer conductors.


Fig. 5-Right polygons with some sides open, examples of outer conductors.

The basic requirement of this category of polygons is simple. They comprise every cross-sectional contour that retains the same shape if the entire surface recedes by a specified small amount. It appears that this category includes all polygons formed of straight lines tangent to a circle, and excludes all other shapes. The circle is one limiting case. Every one of these contours presents the same skin resistance if circumscribed on the same size of circle.

With respect to all imner conductors, the contours with more acute outer angles make less effective use of their perimeters, because there is more extreme concentration of current near the outer angles. By analogy with electric potential gradient, the current density is theoretically infinite at any angle whose outer side is exposed to the magnetic fied. In spite of this fact, it is noted that the square suffers very little, effectively: utilizing $\pi / 4$ of its perimeter. An opposite extreme is the elongated rhombus, Fig. 4 (b).
Fig. 4(e) as an inner conductor might be expected to approximate the behavior of a rectangle. On the contrary, it has much less effective utilization of its contour because of the acute angles.

Referring to the more extreme shapes in Figs. 3 and 4, their equality of resistance as imner and outer conductors is remarkable and unexpected, because the current distribution is radically different in these alternative functions.

Every example of the polygon rule is an evaluation of a certain definite integral. Some of the more unusual cases may he integrals that camot be evaluated by any procedure known to mathematicians. A long table could be prepared on the basis of this one rule.
The polygon rule offers a fascinating variety of examples based on a single theorem. It has some pratical
utility in computing or estimating the skin resistance of inner and outer conductors of various polygon cross sections. Its greatest value lies in the ideas on le perceived in its examples, particularly the effect of extreme current concentration on acute angles exposed to the fiekl. It is another interesting application of the basic "incremental-inductance rule."

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# Active-Error Feedback and Its Application to a Specific Driver Circuit* 

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#### Abstract

Summary-A short discussion of the advantages and disadvantages of active-error feedback in amplifier design is given. Such feedback can yield all the advantages of ordinary negative feedback without gain reduction and is particularly suitable for use in reducing the distortion of individual amplifier stages.

Active-error feedback is applied to a cathode follower by amplifying the difference between its input and output voltages, then adding the amplified error to the output. The resulting driver has very low output impedance and low distortion and is especially useful for driving the grid of an output tube far positive. A direct-coupled version of the circuit using ordinary miniature tubes had an output impedance of 5.6 ohms and could supply several hundred milliamperes of positive current. The theory of the circuit agrees with experiment, and the distortion of the driver when driving an output tube grid to the diode line is found to be far less than that of an ordinary cathode-follower driver.


[^16]
## I. Active-Error lieedback

 ( $\int \begin{aligned} & \text { NIIIKE ORIDIN IRY negative feedback, where } \\ & \text { a portion of the available gain of an amplifier is }\end{aligned}$ expended in obtaining the benefits of feedback, active-error feedback ( $\triangle \mathrm{FF}$ ) is a type of feedback with which no such direct gain reduction occurs. With . IEF, a portion of the output signal from an amplifier or singleamplifier stage is subtracted from the input signal, then the resulting difference amplified in an external circuit whose gain is equivalent to the extra gain necessary with ordinary feedback. If the portion of the output subtracted is nominally equal to the input, the difference is proportional to the error or distortion in the output. This error is then injected back into the original circuit with the proper polarity to reduce the output error. Although the principle of $A E F$ has been used in servomechanisms, it does not seem to have been as fully ex-ploited in amplifier design as it deserves to be. ${ }^{1}$ Therefore, it is worthwhile to discuss its advantages and disadvantages in this application in some detail and to present a specific example of this type of feedback.


Fig. 1-(a) Block diagram showing connection of active error feedback around the amplifier of gain $A_{0}$; (b) block diagram showing a method of combining active error feedback and ordinary negative feedback.

Fig. 1(a) shows a block diagram of a general AEF circuit. The circuit, the gain of which is to be stabilized, the distortion and output impedance of which are to be reduced, and the frequency response of which is to be improved, has a mid-frequency numerical gain of $A_{0}$. We have drawn this block diagram in terms of the positive mid-frequency numerical gains $A_{0}$ and $G_{0}$ rather than the complex phasor gains $A(f)$ and $G(f)$ in order to show explicitly the possible signs which may occur in the midfrequency region. The symbol

indicates addition and

[^17]
subtraction of the two input voltages. A variable voltage entering the junction at a plus sign goes through unchanged in sign, but a voltage entering at a minus sign has its polarity inverted. The plus-or-minus signs within circles in Fig. 1(a) go together as do those without circles, but the signs of the two sets may be specified independently.

The block diagram shows that the output voltage is multiplied by a factor $\alpha$, the result subtracted from the input voltage, and the resulting error voltage amplified by the factor $G_{0}$. Since only error voltage is amplified in this branch of the circuit, the amplifier of mid-band gain $G_{0}$ need handle only fairly small signals and need not itself be distortionless. Finally the amplified error voltage is added to the input in such a phase sense that it reduces the difference between the input and $\alpha$ times the output. It is usually most convenient to make $\alpha$ the pure numeric $A_{0}{ }^{-1}$. Then the AEF tends to make the output follow the input with no gain reduction.

Analysis of the block diagram yields the following result for the over-all gain $e_{\text {out }} / e_{\text {in }}$,

$$
\begin{equation*}
e_{\text {out }} / \varepsilon_{\text {in }}=A_{0}\left[1+G_{0}\right] /\left[1+\alpha A_{0} G_{0}\right]=A_{0} . \tag{1}
\end{equation*}
$$

The second equation follows on taking $\alpha=A_{0}{ }^{-1}$. If we continue to take $\alpha=A_{0}{ }^{-1}$ and generalize (9) for complex phasor gains, we obtain

$$
\begin{equation*}
\frac{e_{\text {out }}}{\mathbf{e}_{\text {in }}}=\frac{A(f)\lfloor 1+G(f)\rceil}{1+G(f) A(f) / A_{0}} . \tag{2}
\end{equation*}
$$

So long as $G(f) A(f)$ is considerably greater than $A_{0}$, (2) reduces closely to

$$
\frac{e_{\mathrm{out}}}{\mathrm{e}_{\mathrm{in}}} \cong A_{0},
$$

the midband gain. We thus see that AEF can considerably extend the flat response of the $A$-circuit provided that the frequency response of the $G$-circuit is initially the wider and that $G_{0}$ is considerably larger than unity. A straightforward calculation also shows that harmonic components and the output impedance are each reduced by the factor $\left|1+G(f) A(f) / A_{0}\right|$, which will be considerably greater than unity over the frequency range of interest. Finally, (2') shows that the fundamental-signal gain of the circuit is stabilized by the AEF circuit to the mid-frequency value when $\alpha=A_{0}{ }^{-1}$. Thus, the circuit yields the usual advantages of negative voltage feedback without the usual decrease of gain. The additional gain required is of course supplied by the active $G$-circuit. As in any feedback circuit, it is necessary, in order to avoid regeneration, that $G(f) A(f) / A_{0}$ become less than unity before the phase shift of the combination reaches 180 degrees. The usual Nyquist criterion for stability is applicable here with $\beta$ given by $-G(f) / A_{0}$.

A combination of AEF and negative feedback can be
applied to an amplifier as shown in Fig. 1(b). However, analysis of this circuit shows that the combination acts as though the extra gain of the AEF were directly in the normal negative feedback loop. Thus, although the effective negative feedback is increased, the AEF has not appreciably simplified the problem of equalizing the a mplifier and feedback paths to avoid regeneration and to achieve unconditional stability. This latter statement needs qualification in one way. Since the gain $G$ is essentially outside the main amplification path, its phase and amplitude may be conveniently controlled without the necessity (which might arise with the same total gain used only with negative feedback) of having to equalize the gain $A$ and possibly thereby reduce the effective feedback at high or low frequencies. In addition, if the entire circuit is to be direct coupled, the splitting of the effective feedback into two paths in the fashion of Fig. 1(b) will usually require a smaller dc supply voltage than would be needed had all the available gain been distributed serially in the direct amplification path. When a large amount of effective negative feedback is required, its realization in a direct-coupled amplifier with reduced supply voltages may be an important economic advantage.


Fig. 2-Block diagram showing alternative connection of active error feedback.

The AEF circuit of Fig. 1(a) may be rearranged to inject the amplified error voltage into the output rather than the input of the amplifier of gain $A(f)$. The resulting circuit, with some of the possible signs indicated, is shown for midband gains in Fig. 2. If $\alpha$ is taken as $A_{0}{ }^{-1}$ as usual, the complex gain of the circuit is found to be

$$
\begin{equation*}
\frac{e_{\text {out }}}{e_{\text {in }}}=A_{0}\left[\frac{A(f)+G(f)}{A_{0}+G(f)}\right] . \tag{3}
\end{equation*}
$$

Thus, the gain will be stabilized to the value $A_{0}$ over a wide frequency range as long as $G(f)$ is appreciably larger in magnitude than $A(f)$. Here it is necessary for stability that $G(f) / A_{0}$ become less than unity before the phase shift of $G(f)$ reaches 180 degrees.

The output impedances $Z_{A}$ of amplifier $A$ and $Z_{G}$ of amplifier $G$ will be connected together across the load in the circuit of Fig. 2. The effective output impedance of the combination (the internal impedance of the composite unit) is readily found to be

$$
\begin{equation*}
Z_{\text {ieff }}=\frac{Z_{A} Z_{G}}{Z_{G}+Z_{A}\left[1+G(f) / A_{0}\right]} \cong \frac{Z_{G}}{1+G(f) / A_{0}}, \tag{4}
\end{equation*}
$$

where the second equation follows when $|G(f)| / A_{0} \gg 1$ and when $Z_{G}$ and $Z_{A}$ are comparable. These conditions also lead to the gain given by (3).

When $A_{0}$ is large, it will usually be inconvenient to make $G_{0} / A_{0} \gg 1$. In this case, the AEF circuit of Fig. 1 (a) will be more suitable than that of Fig. 2. However, when AEF is applied around an individual stage of relatively low gain, the circuit of Fig. 2 may become preferable. This may be particularly the case when added power or current handling capacity is required, since the $A$ and $G$ amplifier outputs are effectively in parallel and thus need each supply only part of the total required output power or current. Examples are a driver which must supply appreciable undistorted current, or a power output stage. The former will be discussed in more detail in the next section.

The distinction between amplified (or active) negative feedback and AEF should be emphasized. Amplified negative feedback would be obtained if the amplifier $G$ amplified a portion $\alpha$ of the output only. It is only when the error between a portion of the output and the input is amplified that AEF is obtained. It may be noted that amplified negative feedback produces the same reduction in output impedance that AEF does, but that while AEF stabilizes but does not reduce the midband gain, amplified negative feedback reduces it by about the same factor that the output impedance is reduced. It is obvious that while the present discussion has dealt only with AEF involving the output voltage, an AEF circuit could be applied which would make the output current, rather than the output voltage, follow the inpl:t voltage (or current).

## II. The Augmented Cathode Follower

For many applications, a circuit having wide dynamic range and low output impedance is desirable. For example, the direct-coupled driver of an output tube which is to be driven into the positive-grid region must have such characteristics. The input resistance of such a tube may be as low as 100 ohms when its grid is driven far positive. Further, this resistance is a strongly nonlinear function of grid voltage. To a void a ppreciable distortion, the driver of such a tube must itself have an output impedance considerably below 100 ohms and must, at the same time, be capable of supplying large positive peak grid currents.
An arbitrarily low output impedance can be obtained from an ordinary plate-loaded amplifier by applying sufficient negative voltage feedback around it. However,
the load current must flow through the output plate resistor, which is often undesirable, and the change of dc voltage level between the grid and plate of the output tube may complicate the use of such a circuit in a direct-coupled amplifier. Even if the driver tube itself is a cathode follower whose output impedance is reduced by ordinary inverse feedback around previous amplifier stages, these stages will be in the direct amplification path, again complicating its use in a direct-coupled circuit. In the present section, we show how these difficulties may be avoided by applying AEF to a cathodefollower driver. The resulting direct-coupled circuit has both very low output impedance and no appreciable change in de voltage level between input and output.

Fig. 3 indicates one way of adapting the AEF circuit of Fig. 2 to a cathode follower. We shall call the resulting circuit a parallel augmented cathode-follower driver (PACFD). The type of AEF shown in Fig. 2 is particularly applicable to the cathode follower because the latter's gain is near unity and thus the external gain $G$ need only be greater than unity to be effective in reducing output impedance and distortion. Further, $\alpha$ can be conveniently taker equal to unity.


Fig. 3-The parallel augnented cathode-follower driver.

As shown in Fig. 3, the difference between the input $e_{1}$ and the output $e_{k}$ of tube $V_{1}$ is amplified by the differential amplifier ${ }^{2}$ consisting of $V_{3}$ and $V_{4}$, then applied to the grid of the parallel cathode follower $V_{2}$ to reduce the error between $e_{1}$ and $e_{k}$. In this direct-coupled circuit, it is desirable that $V_{3}$ be of the same tube type as $V_{1}$ and $V_{2}$, in order that operating biases be correct. In an ac coupled version of the circuit both $V_{3}$ and $V_{4}$ could be, for example, the halves of a single 12AX7. It is worth mentioning that a cathode follower can be augmented in another way by using the tube half $V_{2}$ as a cathode follower in series with $V_{1}$ so that the cathode of

[^18]$V_{2}$ is connected to the plate of $V_{1}$. Then the grid of $V_{2}$ could be direct-coupled to the plate of $V_{4}$ without the voltage divider necessary in Fig. 3. We shall designate such a unit a series augmented cathode-follower driver (SACFD). The SACFD is superior to an ordinary cathode-follower driver (CFD) but inferior to a PACFD, as we shall see below. In addition, its dynamic range is limited, for a given supply voltage value, by the necessary voltage division across $V_{1}$ and $V_{2}$ in series, which does not occur with the PACFD.

A straightforward analysis of the midband equivalent circuits of the SACFI) and PACFD yield the following results for their gains and internal impedances:

$$
\begin{align*}
G_{S} & =\left[\mu\left(1+g_{1}\right)+\mu^{2}\right] /\left[\mu g_{2}+(1+\mu)^{2}+(\mu+2) r_{p} / R_{k}\right],  \tag{5}\\
r_{i S} & =r_{p} /\left[\mu g_{2} /(2+\mu)+(1+\mu)^{2} /(2+\mu)+r_{p} / R_{k}\right],  \tag{6}\\
G_{P} & =\mu\left(1+g_{1}\right) /\left[\mu g_{2}+2(1+\mu)+r_{p} / R_{k}\right],  \tag{7}\\
r_{i P} & =r_{p} /\left[\mu g_{2}+2(1+\mu)+r_{p} / R_{k}\right] . \tag{8}
\end{align*}
$$

$G_{S}$ and $r_{i S}$ refer to the SACF1), $G_{P}$ and $r_{i P}$ to the P.ICFD. In the above equations, the arithmetical gains $g_{1}$ and $g_{2}$ of the differential amplifier are those indicated on Fig. 1; they are slightly unequal, with $g_{2}$ the larger. Note that the algebraic gain corresponding to $g_{2}$ is negative. ${ }^{2}$ It is also assumed that the tube halves $V_{1}$ and $V_{2}$ have the same characteristics. For most purposes, we shall ignore the small difference between $g_{1}$ and $g_{2}$ and designate them both by g . The above equations show that if $\mu \mathrm{g}$ is sufficiently large and $r_{p} / R_{k}$ small, both $G_{S}$ and $G_{P}$ will approach unity closely. Further $r_{i s}$ will approach $r_{p} /(\mu+g)$ and $r_{i p}$ will be approximately $r_{p} / \mu g$. Note that were amplified negative feedback used in the PACFD (e.g., by grounding the grid of $V_{3}$ for input signals) instead of AEF, $g_{1}$ would then be zero, and $G_{P}$ would be reduced to about $g_{2}{ }^{-1}$ while $r_{i p}$ would remain unchanged.

For comparison with the above results, the equations pertaining to an ordinary cathode follower are

$$
\begin{align*}
G & =\mu /\left[1+\mu+r_{p} / R_{k}\right],  \tag{9}\\
r_{i} & =r_{p} /\left[1+\mu+r_{p} / R_{k}\right] . \tag{10}
\end{align*}
$$

When $r_{p} / R_{k}$ is small and $\mu$ appreciably larger than unity, we see from these results that to good approximation the output impedance of the SACFD is reduced over that of an ordinary cathode follower of the same characteristics as $V_{1}$ by the factor $(\mu+g) / \mu$ and that of the PACFD is reduced by the factor $g$. The principal reason for the difference is that the error voltage at the plate of $V_{4}$ is degenerated in the SACFD by a factor of about $\mu$ when applied to the plate of $V_{1}$ and so is less effective in reducing the output error than is that of the PACFD. Such degeneration is instrumental in reducing the dynamic range of the SACFD even further. Since the PACFD makes superior use of the same tubes required in the SACFI), we shall concentrate on the former in the rest of this work.

It may be noted that the double cathode follower ${ }^{3}$ achieves, with two tubes in series, about the same smallsignal gain and output impedance as the PACFD. The top input tube is plate loaded and its cathode connected to the plate of the bottom tube. The bottom tube is itself driven from the plate of the top tube. Neither the SACFI nor the double cathode follower are comparable to the PACFI) as drivers, however. In the SACFI), the driving current must pass through both the upper series tube and the lower cathode-follower tube. In the double cathode follower, it must pass through both the load resistor $R_{L}$, which should be appreciably greater than $r_{p}$, and through the upper tube. In the PACFI), the driving current is supplied by both the cathode-followers $V_{1}$ and $V_{2}$ of Fig. 3, essentially in parallel. The dymamic range and current handling capacity of the PiClil) are thus much superior to those of the other two circuits.

## ill. Comparison of Theory and Experiment

The circuit of Fig. 3 was constructed with the parameter values and tubes slown. It was found that its noload gain was 0.986 . Next, the output voltage was measured as a function of total load resistance $R_{L}$ (the parallel combination of $R_{k}$ and any added load) for a fixed input voltage. The measurements were carried out at $10^{4} \mathrm{cps}$ using a $30 \mu f$ oil capacitor in series with a variable load resistance; only at the lowest load resistances was the capacitative reactance of importance.


Fig. 4-l)ependence on load resistance of the normalized output voltage of the l'ACFl).

Fig. 4 shows the load dependence of the output voltage $e_{k}$ normalized with respect to that without load $\left(e_{k}\right)_{\infty}$. The theoretical line of this figure was calculated using (7) with $R_{k}$ replaced by $R_{L}$. The values $g=70, \mu=16$ and $r_{p}=6.45$ kilohms were employed; these values are in reasonable agreement with published curves. Fig. 4 shows that these values are indeed a good choice, and that theory and experiment are in agreement. In addli-

[^19]tion, the internal impedance, defined as the added load necessary to make $e_{k} /\left(e_{k}\right)_{\infty}=0.5$, is shown to be 5.6 ohms . For comparison, the internal impedances of the SACFD and CFD using the same tubes were found to be of the order of 70 and 370 ohms, respectively. The above definition of $r_{i P}$ leads to the same result for this quantity as that given in (8), which was calculated on the basis of a grounded input and a measuring signal applied to the output. Alternatively, if $r_{i p}$ is again determined by loading the output but defined as the added load required to make $e_{k} / e_{1}=0.5$, the expression for $r_{i P}$ becomes
\[

$$
\begin{equation*}
r_{i P}=r_{P} /\left[\mu\left(2 g_{1}-g_{2}\right)-2-r_{p} / R_{k}\right] . \tag{11}
\end{equation*}
$$

\]

For large $\mu \mathrm{g}$, it does not differ appreciably from (8).
Next, the amplified error voltage $e_{2}$ (see lig. 3) was measured under the same conditions as above for a fixed input voltage $e_{1}$. The normalized quantity $e_{2} / e_{1}$ is plotted in Fig. 5 vs $R_{L}$. The small-signal equivalent circuit yields a value for this ratio of

$$
\begin{align*}
e_{2} / e_{1} & =\left[\mu\left(2 g_{1}-g_{2}\right)\right. \\
& \left.+g_{1}\left(2+r_{p} / R_{k}\right)\right] /\left[\mu g_{2}+2(1+\mu)+r_{p} / R_{k}\right\rfloor . \tag{12}
\end{align*}
$$

This quantity is slightly greater than unity even for $R_{k}$ infinite. The solid line of 1 Fig. 5 was calculated from (12), replacing $R_{k}$ by $R_{L}$ and using the same values for the tube parameters as those used for Fig. 4. Again, agreement between theory and experiment is exceptionally goorl. It is of interest to note that at very large loads $e_{2} / e_{1}$ may be much greater than unity; its maximum value will be approximately $g$ if this value can be achieved without overdriving the tube $V_{4}$.

Finally, it should be pointed out that the data of Figs. 4 and 5 were measured with values of $e_{1}$ of the order of 0.1 volt or less. The equivalent circuit and the resulting formulas only hold as long as operation is in a linear region. When negative peaks are to be produced across a load sufficiently large that the peak current required exceeds the quiescent current in $R_{k}$, the tubes $V_{1}$ and $V_{2}$ will be cut off and negative peak limiting will occur. Only by employing voltages sufficiertly small that such limiting did not occur could an accurate undistorted value of $e_{k}$ be obtained when very low load resistances were used. This negative peak limiting is the reason why a single P\CFI) or a pair in push-pull cannot be conveniently used to drive a load like a loudspeaker directly even though the small-signal impeclances may be matched.

## IV. Comparison of Grid-Driver Circuits

The PACFD is ideally suited for a grid driver. Because it uses two cathode-followers essentially in parallel ( $V_{1}$ and $V_{2}$ ), it can supply twice the peak positive grid current of a single unit. In addition, as the current increases, the $g_{m}$ and $\mu$ of both tubes increase and the $r_{p}$ 's fall. For example, at 50 ma per tube-half, the $\mu$ and $g_{m}$ of a 5687 are approximately 19 and $12,000 \mu \mathrm{mhos}$, re-


Fig. 5-Dependence on load resistance of the normalized error voltage of the PACFD.
spectively. Using $g=70$, (8) or (11) predict an internal impedance of the l'A( $F 1$ ) of only about 1.15 ohms instead of the value of 5.6 ohms found for small signals with the circuit of Fig. 3 .


Fig. 6-Comparison between the intermodulation distortion of three drivers when direct-coupled to an output tube grid.

In Fig. 6, we give a comparison between the distortion generated by a ( CHD , a S.ACFl), and a PA(CI) when clirect coupled to a power tube grid load. The lowest line, marked "no load," shows the distortion in the unloaded PACFI) output. 'This distortion arises almost entirely from the preceding amplifier stage. The output tube
was an 807 , triode connected, with 400 volts on the plate. It had an unbypassed 25 -ohm cathode resistor and formed half of a push-pull output circuit with output transformer and resistive load. The other half of the push-pull output circuit was, in each case, driven by a driver identical to that measured. ${ }^{4}$ The intermodulation distortion was measured at the output of the driver and employed 60 and $5,600 \mathrm{cps}$ signals, mixed $4: 1$. The de bias of the output tube was adjusted to -42.5 volts so that the grid was driven positive when the rms driver voltage exceeded 30 volts. It is this positive grid region which is presented in Fig. 6.

The dotted line is the approximate peak grid current supplied by the driver. When the rms driver voltage is 70 volts, the grid is driven positive by 56.5 volts peak, and we see that it draws a peak current of about 200 ma. It is obvious from lig. 6 that the SACFI) is a considerable improvement on the ( Cl I , and the PACFI) an improvement on the SACFI) over most of the range considered. For applied voltages greater than 60 to 65 volts rms, the grid of the output tube loses control of the output current on positive peaks; the point at which control is lost defines the diode line of the output tube. It is seen from the figure that the distortion of all the drivers increases rapidly for larger voltages. Oscillographic observations showed, however, that the PACFD was capable of driving the grid of the output tube considerably beyond the point where the output voltage of the output tube began to show peak clipping arising from diode-line limiting. Even in this region, however, appreciable distortion of the grid signal could not be observed on the CRO.
${ }^{4}$ The push-pull driver circuit used in these measurements incorporated a special feedback loop which reduced even-order harmonic distortion greatly at the driver outputs. Therefore, the internodulation distortion results obtained at one of the push-pull driver outputs may be appreciably smaller, especially for the case of the CFD, than would be attained in practice without such a feedback loop. Nevertheless, the distortion curves still afford a valid comparison between the relative distortion of the three types of drivers.

# A Semiconductor Diode Multivibrator* 

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#### Abstract

Summary-The operation and design of a novel semiconductor multivibrator circuit is described. Consisting of one double-base diode, one diode, three resistors, and one capacitor, the diode multivibrator affords a circuit economy of two-to-one over the more conventional Eccles-Jordan configuration. The active element is the double-base diode, which is a three-terminal single-junction nega-tive-resistance semiconductor device. Applications of the diode multivibrator to square-wave oscillators and delayed-pulse generators are illustrated. Control of the period of oscillation or variation of the pulse-delay time is accomplished by the variation of a single capacitor. Hence, this new circuit simplifies the design of astable, monostable or bistable multivibrators.


## Introduction

TIII: INCREASING use of quantized signals as "information carriers" in modern electronic systems has caused a growing concern over the power requirements, complexity, and expense of pulse-type circuits. Thus the advent of the transistor as a lowpower, long-lifetime, and sub-miniature active element was welcomed most enthusiastically by pulse-circuit engineers as a possible solution to the dilemma of increasing system complexity and size.


Fig. 1-Vacuum-tube multivibrator.

One of the oldest, and perhaps most fundamental, of the digital-type circuits is the multivibrator. ${ }^{1}$ As a twostate circuit, analogous to the mechanical relay, the multivibrator extended the use of relaying to highspeed operations. And since the "on-off" device is virtually the "nerve-cell" of all counting and logic circuits, the multivibrator ultimately became a fundamental building block of modern electronic computers.

The semiconductor multivibrator to be described in this paper may be considered as a diode flip-flop. It has been made possible by the invention of a new semicon-

[^20]ductor device - the double-base diode. ${ }^{2-1}$ In the cliode multivibrator one diode "drives" another in a reciprocal "on-off" relationship. A significant departure from the well-known Eccles-Jordan circuit, which is illustrated in Fig. 1, using vacuum-tubes, and in Fig. 2, using transistors, is that only one active circuit element is required for the regenerative action.


Fig. 2-Junction-transistor multivibrator.

The astable and monostable diode multivibrators will be treated here. The use of the former as a square-wave oscillator and the latter as a delayed-pulse generator will be clescribed. Because of the almost two-to-one reduction in circuit components required by the diode multivibrator, when compared to the Eccles-Jordan transistor configuration, a higher degree of circuit simplicity, miniaturization, and economy may be achieved.


Fig. 3-Diode multivibrator.

## Principles of Operation

The basic circuit configuration of the diode multivibrator is illustrated in Fig. 3. Its relative simplicity is immediately apparent when compared to the transistor

[^21]multivibrator shown in Fig. 2. Operation of the astable diode multivibrator may be described as follows. Capacitor $C$ is charged from the battery supply through the resistance $R_{2}$ and the diode, $D$. During the charging cycle of the capacitor, $D$ is conducting but the doublebase diode is in the cut-off state. When the potential across the capacitor becomes equal to or greater than the peak-point potential of the double-base diocle, the latter becomes unstable and switches into the conducting state. The junction potential ( $b$ in the circuit of Fig. 3) is then clamped to a low value-almost to the ground potential, thus causing the diode $D$ to become cut off. When the diode is in its nonconducting state, point $a$ is virtually isolated from point $b$. The capacitor will then discharge through the resistance $R_{1}$ until the potential at $a$ is approximately equal to the junction potential of the double-base diode. At this instant, the diode becomes conducting again. When the diode reverts to its conduction state, the current through the junction of the double-base diode decreases and the latter is driven into its cut-off state. Capacitor $C$ will then recharge and the cycle will be repetitive.

The waveforms generated by the astable circuit of Fig. 3 are illustrated in Fig. 4. Since capacitor $C$ alternately charges and discharges through a fixed resistance ( $R_{2}$ and $R_{1}$, respectively), the waveform at $a$ consists of a periodic exponential rise and decay. During the time


Fig. 4-Diode multivibrator waveforms.
$D$ is conducting, the waveform at $b$ will be almost identical to that at $a$. However, when the double-base diode becomes conducting, the potential at $b$ is clamped to a near-ground value until the capacitor has completed its discharge cycle. When the double-base diode is in its cut-off state, the current through $R_{3}$ will be comparatively low. However, when the double-base diode switches "on," its bar resistance drops by an order of magnitude and the current through $R_{3}$ increases. Thus, the current through $R_{3}$ will be either high or low, depending upon the operating state of the double-base diode. Consequently, the waveform across $R_{3}$ (at point C) will be a square-wave as illustrated in Fig. 4. It is apparent that the frequency and symmetry of this square wave is dependent upon the time constants associated with $R_{1}, R_{2}$ and $C$.

## Graphical Analysis

A better understanding of the operation of the diode multivibrator may be obtained by considering the equivalent circuits illustrated in Fig. 5. Figure 5(a) illustrates the dc equivalent circuit of the diode multivibrator when $D$ is conducting. The capacitor is omitted from the dc circuit and the diode is assumed to have negligible forward resistance. Writing the loop equations for the circuit of Fig. $5(\mathrm{a})$, the following set is obtained:

$$
\begin{align*}
& E=\left(R_{1}+R_{2}\right) I_{1}-R_{1} I_{d}  \tag{1a}\\
& O=-R_{1} I_{1}+R_{1} I_{d}+V_{d} . \tag{1b}
\end{align*}
$$



Fig. 5-Dc equivalent circuits; (a) diode conducting, ( $b$ ) diode nonconducting.

In (1b), $V_{d}=f\left(I_{d} ; E, R_{3}\right)$, which represents the input characteristics of the double-base diode for a given battery supply, $E$, and load resistance, $R_{3}$. Eqs. (1a) and (1b) are easily solved for $V_{d}$ as a function of $I_{d}$. Thus

$$
\begin{equation*}
V_{d}=\frac{R_{1}}{R_{1}+R_{2}} E-\frac{R_{1} R_{2}}{R_{1}+R_{2}} I_{d} . \tag{2}
\end{equation*}
$$

When the diode $D$ in the circuit of Fig. 3 is nonconducting, the steady-state equivalent circuit of Fig. (5b) may be obtained. It is assumed that the resistance $R_{1}$ is effectively isolated from the double-base diode by the very high-back resistance of the diode. For the circuit of Fig. (5b),

$$
\begin{equation*}
V_{d}=E-R_{2} I_{d} . \tag{3}
\end{equation*}
$$

The application of (2) and (3) to the operating characteristics of the double-base diode permits a graphical load-line analysis similar to that used in vacuum-tube and transistor circuits design. This is illustrated in Fig. 6 (next page). For the condition that the diode conducts, the steady-state input load line is determined by (2) and is represented by the dashed line. The intersection of load line with ordinate axis is at a point

$$
V_{d}=E \frac{R_{1}}{R_{1}+R_{2}},
$$

and the slope of the load line is the parallel combination of $R_{1}$ and $R_{2}$. For the condition that the diode is nonconducting, the load-line characteristic is determined by (3) and is represented by the solid line (slope $=R_{2}$ ) in Fig. 6.

In order for the diode multivibrator to be astable, or free running, the following conditions must be satisfied.

1. When the diode is conducting, the input load line cannot intersect the double-base diode characteristic in the cut-off region.
2. When the diode is nonconducting, the input load line must intersect the double-base diode operating characteristic in the transition, or negative-resistance, region. Consequently, the circuit conditions for the freerumning multivibrator are as follows:

$$
\begin{align*}
\frac{R_{1} E}{R_{1}+R_{2}} & >V_{p}  \tag{4a}\\
\frac{E}{R_{2}} & \leqq I_{v} . \tag{4b}
\end{align*}
$$

In (4a) and $4(\mathrm{~b}), V_{p}$ is the peak-point potential of the double-base diode and $I_{v}$ is the input current corresponding to its valley point. The operating path of the multivibrator, in relation to the input characteristics of the double-base diode, may be approximately determined from the graphical analysis, as indicated in Fig. 6.


Fig. 6--Operating characteristic of the diode multivitrator.

## Design Analysis

From the steady-state circuit analysis it has been possible to determine the conditions required for sustained oscillation of the multivibrator. The frequency and symmetry of the generated waveforms may be determined from a consideration of the charging and discharging cycles of the capacitor, $C$, in Fig. 3 . It may be noted that when $D$ is conducting, the capacitor is charging. This also corresponds to the conclition that the double-base diode is in its cut-off state and that the potential at $c$ is comparatively high. Hence, the diode conducts during the positive cycle of the square-wave and is nonconducting during the negative cycle of the square wave. The charging and discharging equivalent circuits, corresponding to the positive and negative portions of the output waveform, are illustrated in lig. 7.

In terms of the Laplace operator, $s$, the voltage transform for the circuit of $\mathrm{Fig} .7(\mathrm{a})$ is

$$
\begin{equation*}
\frac{V_{a}}{E}(s)=\left[\frac{R_{1}}{R_{1}+R_{2}}\right]\left[\frac{1}{1+\frac{R_{1} R_{2} C s}{R_{1}+R_{2}}}\right] \tag{5}
\end{equation*}
$$

Closing the switch, $S$, is assumed to be equivalent to I) suddenly becoming conducting. Eq. (5) is solved in
the time domain for the following initial conditions:

$$
\begin{equation*}
\text { at } t=0: \quad V_{i}=\frac{E}{s} ; \quad V_{a}=V_{v}^{\prime} \tag{6}
\end{equation*}
$$

Eq. (6) indicates that a step voltage of magnitude $E$ is applied to the circuit of Fig. 7 (a) at a time when $C$ is charged to a potential of $V_{v}^{\prime}$. Using the standard Laplace transformation techniques (5) and (6) may readily be solved to give the following result in the time domain:

$$
\begin{align*}
V_{a}(t)= & \frac{R_{1} E}{R_{1}+R_{2}}\left[1-\exp \left(-\frac{\left(R_{1}+R_{2}\right) t}{R_{1} R_{2} C}\right)\right] \\
& +V_{v}^{\prime} \exp \left(-\frac{\left(R_{1}+R_{2}\right) \ell}{R_{1} R_{2} C}\right) \tag{7}
\end{align*}
$$

The time that it takes the capacitor to charge to the peak-point voltage, $V_{p}$, may be calculated from (7) by substituting $V_{a}(t)=V_{p}$ on the left-hand side of the equation. Thus

$$
\begin{equation*}
t_{p}=-\frac{R_{1} R_{2} C}{R_{1}+R_{2}} \ln \left[\frac{\frac{R_{1} E}{R_{1}+R_{2}}-V_{p}}{\frac{R_{1} E}{R_{1}+R_{2}}-V_{v}^{\prime}}\right] \tag{8}
\end{equation*}
$$

where $t_{p}$ denotes the duration of the positive portion of the square wave.

Similarly, the duration of the negative portion of the square wave can be obtained from the equivalent circuit of Fig. $7(1)$ ). For the initial condition that $V_{a}(0)$ $=V_{p}$ and calculating the time it takes the capacitor to discharge to a value $V_{v}^{\prime}$, it is found that

$$
\begin{equation*}
t_{n}=-R_{1} C \ln \frac{V_{v}^{\prime}}{V_{p}^{r}} \tag{9}
\end{equation*}
$$

In (9), $t_{n}$ is the time duration of the negative portion of the output wave.

(b)

Fig. 7-Low-frequency ac equivalent circuits; (a) positive cycle, (b) negative cyole.

The total period, $t_{T}$, of the square wave is the sum, $t_{p}+t_{n}$. Hence

$$
\begin{align*}
t_{T}= & -R_{1} C\left\{\frac{R_{2}}{R_{1}+R_{2}} \ln \left[\frac{1-\left(\frac{V_{p}}{E}\right)\left(\frac{R_{1}+R_{2}}{R_{1}}\right)}{1-\left(\frac{V_{v}^{\prime}}{E}\right)\left(\frac{R_{1}+R_{2}}{R_{1}}\right)}\right]\right. \\
& \left.+\ln \frac{V_{v}^{\prime}}{V_{p}}\right\} \tag{10}
\end{align*}
$$

For convenience, the design parameters are defined:
$P=\frac{V_{p}}{E}, \quad L^{\cdot}=\frac{V_{v}^{\prime}}{E}, \quad a=\frac{R_{1}}{R_{2}}, \quad X=\frac{t_{p}}{t_{n}}$.
If a symmetrical square wave is desired, $t_{r}$ must equal $t_{n}$. From (10) therefore, using constants defined by (11). the following condition for symmetry is established:

$$
\begin{equation*}
\frac{1-P\left(1+\frac{1}{a}\right)}{1-L^{\prime}\left(1+\frac{1}{a}\right)}=\left[\frac{U}{P}\right]^{1+a} \tag{12}
\end{equation*}
$$

If the ratio of the positive portion of the square wave to total period is defined as figure of symmetry, $S$, then

$$
\begin{equation*}
S=\frac{Y}{1+X} \quad \text { where } \quad X=\frac{t_{p}}{t_{n}} \tag{13}
\end{equation*}
$$

It can then be shown that, for any desired figure of symmetry, (10) must satisfy the condition

$$
\begin{equation*}
\frac{1-P\left(1+\frac{1}{a}\right)}{1-U\left(1+\frac{1}{a}\right)}=\left[\frac{U}{P}\right]^{(S / 1-S)(1+a)} \tag{14}
\end{equation*}
$$

For a symmetrical waveform, $S=0.5$. Eq. (12), therefore, is a special case of (14).

The output waveform across $R_{3}$ depends upon the change in the bar resistance of the double-base diode as the latter oscillates between the cut-off and conducting states. Denoting the equivalent lar resistances by $R_{B C}$ and $R_{B S}$, where the subscripts $C$ and $S$ refer to the cutoff and conducting states of the (louble-base cliode, respectively, the equivalent circuit of Fig. 8 may be used


Fig. 8-I.ow-frecuency equivalent-output circuit.
to represent the low-frequency output circuit of the multivibrator. The peak-to-peak amplitude of oscillation across $R_{3}$ is then given by

$$
\begin{equation*}
\left|E_{0}\right|=\left[\frac{1}{1+\frac{R_{3}}{R_{B C}}}-\frac{1}{1+\frac{R_{3}}{R_{B S}}}\right][F] \tag{15}
\end{equation*}
$$

If it is clesired to find the value of $R_{3}$ which maximizes the output voltage amplitude, (15) may be clifferentiated and solved for a maximum in the usual manner. It is found that the value of $R_{3}$ required to obtain a maxi-
mum $E_{0}$ is equal to the geometric mean of the two bar resistances, $R_{B C}$ and $R_{B S}$. Thus

$$
\begin{equation*}
R_{3 m}=\sqrt{R_{B C} R_{B S}} \tag{16}
\end{equation*}
$$

The maximum output voltage, found by substituting (16) into (15), is

$$
\begin{equation*}
\left|E_{0}\right|_{\max }=\left(\frac{1-\sqrt{R_{B S} / R_{B C}}}{1+\sqrt{R_{B S} / R_{B C}}}\right)(I) \tag{17}
\end{equation*}
$$

The peak-point voltage of the double-base diode is linearly related to the interbase potential. ${ }^{4,5}$ Denoting the constant of proportionality by $I I_{12}$, one can write

$$
\begin{equation*}
V_{p}=I_{12} I_{B}^{\prime} \tag{18}
\end{equation*}
$$

In (18), the factor $H_{12}$ is always less than unity and is a constant of the device. When the double-base diode is in its cut-off state, the interbase potential is

$$
\begin{equation*}
V_{B}^{\prime}=\frac{R_{B C}}{R_{B C}+R_{3}} E \tag{19}
\end{equation*}
$$

Substituting (18) into (19) gives

$$
\begin{equation*}
P=\frac{\mathrm{I}_{P}^{\prime}}{E}=\frac{I_{12}}{1+R_{3} / R_{B C}} \tag{20}
\end{equation*}
$$

For maximum output voltage, given by (16), (20) is


Fig. 9.
Fig. 9 illustrates a typical set of input characteristics for a clouble-base diode with a load resistance in the base-two lead. The interbase voltage, $V_{b}$, is constant ( 5.8 volts) only cluring the cut-off state. It is important to note that the valley point of the curves in Fig. 9 moves to the left as the load resistance is increased. However, the valley point approaches a limiting value as the load resistance becomes very large. ${ }^{5}$ Consequently, $V_{v}{ }^{\prime}$ in (6) to (11) depends upon both the input and output load resistances and therefore $U$ in (11) may be treated as an arbitrary circuit design parameter.
${ }^{5}$ J. J. Suran, "Low-frequency circuit theory of the double-base diode," Trans. IRE, vol. ED 2; April, 1955.

## Design Procedure

Based upon the foregoing analysis, one can outline an approximate design procedure for the free-running diode multivibrator. Procedure given below is based on the maximum voltage criterion as defined by (16) and (17).

E(q. (4a) and (4b) specify the two necessary conditions required for oscillation. From (4b),
(a)

$$
\frac{E-V_{v}}{R_{2}}<I_{v}
$$

where $V_{v}$ and $I_{v}$ are the voltage and current, respectively, which correspond to the valley point of the double-base diode input characteristic. Since $V_{v}$ and $I_{v}$ approach constant limits as $R_{3}$ is increased, the limiting values for the valley point may be used in condition 1. Thus, $R_{2}$ can be selected. U'sing this value of $R_{2}, I_{v}{ }^{\prime}$ may be estimated (see Fig. 6). From bar-characteristic curves of the double-base diode, such as illustrated in Fig. $10, R_{B S}$ can be approximated. $R_{B C}$ may be determined fairly accurately from the characteristic curves of Fig. 10. Thus, using (16),

$$
\begin{equation*}
R_{3}=\sqrt{R_{B S} R_{B C}} \tag{b}
\end{equation*}
$$



Fig. 10.
The design parameter $P$ can now be calculated from (20) or (21).
(c)

$$
P=\frac{I_{12}}{1+R_{3} / R_{B C}}
$$

In (c), $I_{12}$ is the voltage ratio defined by (18). Knowing $I_{v}{ }^{\prime}, V_{v}{ }^{\prime}$ may be estimated and

$$
\begin{equation*}
U=\frac{V_{v}^{\prime}}{E} \tag{d}
\end{equation*}
$$

With $P$ and $U$ determined, the ratio $a=R_{1} / R_{2}$ may be
found from (14) for any desired figure of symmetry. A family of curves for $S=0.5$, based upon (12), is illustrated in Fig. 11. The k-values in Fig. 11 correspond to both $U$ and $P$. Hence, the intersections of these curves determine the $a$-values for respective $U$ and $P$ parameters when a symmetrical waveform is desired. $R_{1}$ may be determined from the design parameter $a$. Thus
(e)

$$
R_{1}=a R_{2}
$$



Fig. 11-Symmetry curves for $S=0.5$.
It this point, $a$ must satisfy the condition [from (4a)]

$$
\begin{equation*}
\frac{a}{1+a}>P \tag{f}
\end{equation*}
$$

If (f) is not satisfied, (a) must be re-closen and the design procedure repeated. However, if (f) is satisfied, the design of the multivibrator circuit is complete except for the specification of capacitor $C . C$ is calculated on the desired frequency, from (10). Clearly, for $S=0.5$,

$$
\begin{equation*}
C=\frac{1}{-2 R_{1} f \ln U} \tag{g}
\end{equation*}
$$

## Example:

Suppose that it is clesired to design a $10-\mathrm{kc}$ diode multivibrator, having a symmetrical waveform $(S=0.5)$, and operating from a voltage source of 12 volts. Assume that the double-base diode characteristics are given by Figs. 9 and 10. The first step in the design procedure is to select a value of $R_{2}$ which will satisfy (41)). From Fig. 9, it is apparent that if the input load line intersects the abscissa axis at $I_{d 0}=1.0 \mathrm{ma}$, the astable requirement will almost certainly be satisfied. Hence, from Fig. 6,

$$
R_{2}=\frac{E}{I_{c 0}}=\frac{12}{0.001}=12,0000 \mathrm{hms}
$$

The current flowing into the junction of the double-base diode, when the latter is in its conducting state, will be approximately 1 ma . This is readily established by superimposing value of $R_{2}$ for 12 -volt source, on input characteristics of Fig. 9. From Fig. 10, for $I_{d}=1 \mathrm{ma}, R_{B S}$
may be estimated. Thus, $R_{B S} \approx 1.8 \mathrm{k}$, and $R_{B C} \approx 6.5 \mathrm{k}$. Using maximum voltage criterion defined by (16),

$$
R_{3}=\sqrt{(6.5 \mathrm{k})(1.8 \mathrm{k})}=3.5 \mathrm{k} \text { ohms }
$$

Having determined $R_{3}$ and $R_{B C}$, interbase potential of double-base diode, during cut-off state, may be calculated.

$$
V_{B C}=\left(\frac{6.5 \mathrm{k}}{6.5 \mathrm{k}+3.5 \mathrm{k}}\right)(12)=7.8 \text { volts. }
$$

For $V_{B C}=7.8$ volts, $V_{v}^{\prime}$ (see Fig. 6) may be estimated from the input characteristics. Hence $V_{v}^{\prime} \approx 2.0$ volts. From Fig. 9,

$$
H_{12}=\frac{5.2 \text { volts }}{5.8 \text { volts }}=0.9
$$

The design parameters, $U$ and $P$, defined by (11) and (20), may now be calculated.

$$
\begin{aligned}
& P=\frac{H_{12}}{1+R_{3} / R_{B C}}=\frac{0.9}{1+\frac{3.5 \mathrm{k}}{6.5 \mathrm{k}}}=0.58 \\
& U=\frac{V_{v}^{\prime}}{E}=\frac{2 \text { volts }}{12 \text { volts }}=0.167
\end{aligned}
$$

For a symmetry condition of $S=0.5$, the ratio $a=R_{1} / R_{2}$. may be determined from the design curves of Fig. 11. The intersection of the two curves corresponding to $\mathrm{k}=0.58$ and 0.167 , respectively, is the desired value of $a$. from the input characteristics. Hence $N_{v}{ }^{\prime} \approx 2.0$ volts. from Fig. 11, $a=1.55$, and

$$
R_{1}=a R_{2}=(1.55)(12 \mathrm{k})=18.6 \mathrm{k} \text { ohms }
$$

$R_{1}, R_{2}$ and $R_{3}$ have now been determined to a first approximation. A necessary requirement for oscillation, however, as derived from (4a), is that

$$
\frac{a}{1+a}>P, \quad \text { or } \quad \frac{1.55}{2.55}>0.58
$$

Since the latter inequality is satisfied, the three resistance values will be adequate to sustain oscillations. Capacitor $C$ may now be determined from the frequency requirement. For the symmetrical condition ( $S=0.5$ ),

$$
\begin{aligned}
C=\frac{1}{-2 R_{1} f \ln U} & =\frac{-10^{-6}}{(2)(18.6)(10) \ln 0.167} \\
& =.0015 \mu \mathrm{f}
\end{aligned}
$$

The output amplitude of the waveform is given by (17):

$$
\left|E_{0}\right|_{\max }=\frac{1-\sqrt{\frac{1.8 \mathrm{k}}{5.6 \mathrm{k}}}}{1+\sqrt{\frac{1.8 \mathrm{k}}{6.5 \mathrm{k}}}}[12]=3.7 \text { volts }
$$

It should be noted that (17) does not inchude the variation of the double-base diode parameters with fre-
quency. Experiment shows that the outlined analysis is accurate to within $\pm 10$ per cent for frequencies up to approximately 20 kc . Waveforms obtained from experimental circuits are shown in Figs. 12 and 13.


Fig. 12-Experimental waveforms: (a) germanium dbd-1 kc; (b) silicon $\mathrm{dbd}-10 \mathrm{kc}$.


Fig. 13-Silicon double-base diode waveforms: (a) top- 55 kc , bot-tom- 7 kc ; (b) top-rise time, bottom-fall time, full sweep $=6 \mu \mathrm{sec}$.

## Pulse-Delay Generator

The diode multivibrator may be made monostable if

$$
\begin{equation*}
\frac{R_{1} E}{R_{1}+R_{2}}<V_{p} \tag{22a}
\end{equation*}
$$

and

$$
\begin{equation*}
E / R_{2}<I_{v} \tag{22b}
\end{equation*}
$$

Eq. (22a) fixes the stable operating point of the double-base diode in the cut-off region and (22b) insures
that this is the only stable operating point. If (22a) and (22b) are satisfied, a positive pulse will trigger the double-base diode from the "off" state to the "on" state. The double-base diode will then remain conductive until the capacitor discharges through $R_{1}$. When the diode reverses at the end of the capacitor discharge cycle, the double-base diode becomes nonconductive. But since it is stable in cut-off state, the multivibrator circuit remains stable until the next positive trigger pulse is applied. Thus, the regenerated output waveform duration is

$$
\begin{equation*}
t_{D}=-R_{1} C \text { in } \frac{V_{v}^{\prime}}{V_{p}} \tag{23}
\end{equation*}
$$

[see (9) ]. Fig. 14 shows the waveform generated by monostable multivibrator. Minimum spacing of the trigger pulses is obviously limited by the circuit's time constants.


Fig. 14-Monostable multivibrator waveform.
On the other hand, if

$$
\begin{align*}
\frac{R_{1} E}{R_{1}+R_{2}} & >V_{p}  \tag{24a}\\
\frac{E}{R_{2}} & >I_{n}, \tag{2+b}
\end{align*}
$$

a monostable circuit, having a stable operating point associated with the conductive state of the double-base diode, is obtained. Negative pulses may then be used to trigger the circuit into its regenerative cycle and the output waveform duration will be given [see (8)] by;


Fig. 15-Delayed-pulse generator.
Use of the monostable multivibrator as a clelayedpulse generator is illustrated in Fig. 15. In this circuit, $R_{D}$ and $C_{D}$ are used as a simple differentiating network and diode $D_{D}$ filters out the pulses of unwanted polarity. The output waveform consists of pulses which are gen-
erated by differentiating the trailing edge of the multivibrator output. Diode $D_{0}$ filters out the pulses which are generated by the leading edge of the multivibrator waveform. Hence, the output of the delayed-pulse generator consists of a train of pulses which have the same polarity and repetition rate of the input pulses but which are delayed in time by an interval, $t_{D}$, determined by the time constants of the monostable circuit.

Fig. 16 illustrates the relationship between the pulse delay $t_{d}$ and the magnitude of the multivibrator capacitor $C$, which was obtained for an experimental de-layed-pulse generator circuit. In the circuit of Fig. 16, the conditions defined by (22a) and (22b) are required for operation. Time delays from $50 \mu \mathrm{sec}$ to 2 msec have been obtained for pulse repetition rates from $0-5 \mathrm{kc}$. As is apparent from Fig. 16, the time delay is related to the magnitude of the capacitor $C$ in a linear manner. This relationship is convenient in design and facilitates constructing very simple variable-delay pulse generators.


Fig. 16-Experimental delayed-pulse generator characteristics.

## Conclusion

The diode multivibrator consists of three resistors, one capacitor, one diode and one double-base diode. When compared to the corresponding Eccles-Jordan transistor circuit, the diode configuration affords a two-to-one economy in circuit components. Furthermore, since the diode multivibrator consists of one diode and one double-base diode, as compared to the two transistors required for the Eccles-Jordan circuit, an additional cost economy may be achieved.

Such factors as circuit simplicity, easy designability and component and device economy make the diode multivibrator a considerable competitor to the corresponding transistor circuit. The advantages of the diode circuit should be particularly significant in such complex systems as digital computers and counters where component cost and network complexity can be restrictive.

## Acknowiedgment

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# The Effect of the Source Distribution on Antenna Patterns* 

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#### Abstract

Summary-The response pattern of an antenna is modified when the source subtends an appreciable angle. Under these conditions the observed or resultant pattern is a function of both the true antenna pattern and the source distribution. This problem is important in radio astronomy and in radar. The general problem is formulated and solved for the particular case of a radio telescope antenna consisting of a 40 wavelength broadside array. Graphs are presented showing the effect of sources of various widths on the observed patterns. Using such graphs it is possible to deduce the approximate source extent from an observed pattern. The direct analytical solution for the source distribution from the observed pattern is also considered and the limitations of the various methods are discussed.


## Introduction

TYIIE TRUE RESPONSE pattern of a receiving antenna is obtained when the radiator is a point source situated at a sufficient distance from the antenna. The distance is sufficient if an increase in the distance produces no detectable change in the pattern. Let the true pattern of a receiving antenna be as shown in Fig. 1(a). If the point source is replaced by an extended source at the same distance, the observed pattern is modified as suggested in Fig. 1(b).


Fig. 1-(a) Intema pattern for a point source, (b) for an extended source.

In radio astronomy many of the celestial sources are of sufficient angular extent to modify the observed pattern and it is desirable to be able to deduce the source distribution (extent and shape) from the observed pattern. A similar situation exists in radar when the object under observation is of sufficient angular extent.

Referring to Fig. 2, the general problem of the effect of the source distribution on the observed antenna pattern may be stated as follows:

$$
\begin{equation*}
G\left(\phi_{0}\right)=\frac{1}{A} \int F\left(\phi+\phi_{0}\right) f(\phi) d \phi ; \tag{1}
\end{equation*}
$$

[^22]where
$G\left(\phi_{0}\right)=$ observed or resultant pattern,
$F\left(\phi+\phi_{0}\right)=$ true antenna pattern (as measured with a point source),
$f(\phi)=$ source distribution,
$A=\int f(\phi) d \phi=$ effective angle subtended by source.

All patterns in (1) are proportional to power.


Fig. 2-Antenna pattern, source pattern, and resultant or observed pattern.

In (1) the general problem has been simplified to the one-dimensional case where the patterns are functions only of one co-ordinate. For the purpose of the following discussion this has the advantage that the fundamental problem is retained intact but the analysis is greatly simplified. The simplified situation stated in (1) often occurs in practice as, for example, when the antenna pattern in the direction $(\theta)$, normal to $\phi$, is sufficiently broad compared to the source extent in this direction. In (1) the antenna pattern, $F\left(\phi+\phi_{0}\right)$, and observed pattern, $G\left(\phi_{0}\right)$, are usually known while the source distribution, $f(\phi)$, is unknown. To determine the source distribution it is necessary to solve the integral equation (1). This can be done by assuming various source distributions and calculating the corresponding distributions, $G\left(\phi_{0}\right)$. If a calculated $G\left(\phi_{0}\right)$ distribution can be obtained that agrees with the actual observed distribution $G\left(\phi_{0}\right)$, one can conclude that the assumed source distribution $f(\phi)$ used for this case represents true source distribution or its equivalent. Or one can solve the integral equation in a straightforward manner. This is usually the more difficult procedure. The indirect and direct methods of solution will be discussed in that order.

It is assumed that the cosmic signals under consideration are incoherent so that the power received is proportional to the incident power flux from the source integrated over the solid acceptance angle of the antenna.

Eq. (1) applies not only to antenna problems but to many other situations where instead of an antenna pattern there is a modifying function $F\left(\phi+\phi_{0}\right)$ which disturbs the actual distribution $f(\phi)$ so as to yield the observed function $G\left(\phi_{0}\right)$.

## Cases of Point Source and I.arge Extended Source

Referring to Fig. 2 the value of the observed distribution $G\left(\phi_{0}\right)$, when the main lobe of the antenna is displaced by an angle $\phi_{0}$ from the center line of the source, is given by

$$
\begin{equation*}
G\left(\phi_{0}\right)=\frac{1}{A} \int_{-\alpha}^{+\infty} F\left(\phi+\phi_{0}\right) f(\phi) d \phi \tag{2}
\end{equation*}
$$

where

$$
A=\int_{-\alpha}^{+\alpha} f(\phi) d \phi
$$

The over-all source extent is $2 \alpha$. The distributions are all power patterns so that $A$, the area uncler the source pattern, is in fact the total power flux of the source.

For a point source the source pattern in Fig. 2 collapses to a single vertical spike at $\phi=0(\alpha=0)$. The observed pattern is then given by (2) with $f(\phi)=0$ except at $\phi=0$. For this case (2) reduces to $G\left(\phi_{0}\right)=r^{\prime}\left(\phi_{0}\right)$ or, in general,

$$
\begin{equation*}
G(\phi)=F(\phi) \tag{3}
\end{equation*}
$$

Thus, for a point source the observed pattern is identical with the true antenna pattern.

At the other extreme let us consider the case of an extended source that is much wider thatn the antenna pattern as suggested in Fig. 3. L.et the source be a step


Fig. 3-Case of source pattern that is much wider than antenna beam width.
function equal to unity between $+\alpha$ and $-\alpha$ and zero elsewhere. The resulting observed pattern $G\left(\phi_{0}\right)$ from (2) is as shown in the figure. It is to be noted that in the range of $\phi_{0}$ between $\alpha-\beta$ and $-(\alpha-\beta)$ the observed distribution is a constant like the source although reduced by a factor $B / A$, where $B$ is the area under the antenna pattern and $A$ is the area under the source pat-
tern $(=2 \alpha)$. $\operatorname{tern}(=2 \alpha)$.

## O.S.U. Radio Telescope Antenna

The preceding cases are idealized. Turning now to the case of an actual antenna, such as the Ohio State University radio telescope antenna, let us consider the ef-
fect of the source distribution for three conditions of antenna operation: (1) All elements in phase (singlelobe pattern), (2) two halves of antenna in phase opposition (split-lobe pattern), and (3) comparison arrangement [resultant pattern equal to the difference of (1) and (2)].

The O.S.U. radio telescope antenna, shown by the photograph in Fig. 4, consists of an array of 96 helical beam antennas mounted on a steel ground plane 160 feet long (east-west) by 22 feet wide. The entire antenna pivots like a meridian transit instrument. At 250 mc the antenna is approximately 40 wavelengths long by 5.6 wavelengths wide. With all helices in phase the beamwidths at 250 mo are approximately 1 degree in right ascension (east-west) by 8 degrees in declination.


Fig. 4--Photograph of the Ohio State University
radio telescope antenna.
Although the antenna can be operated at frequencies between 200 and 300 nce the patterns at orly the center frequency of 250 me are discussed in this article. All helices are right-handed so that the antenna is responsive to the right circularly polarized component of the incident radiation, which is usually of an incoherent nature. In operation the antenna is set at a fixed declination and as the earth rotates a trace or profile is obtained on the recorder as a celestial raclio source crosses the meridian. This recorder profile is the observed pattern $G(\phi)$. The helices are arranged in 4 rows with 24 helices in each row. The long (24-helix) rows are east-west and determine the pattern of the antenna in right ascension while the short (4-helix) rows are at right angles and determine the pattern in declination. Since the beamwidth in declination is sufficiently wide ( 8 degrees) it will be convenient to reduce the problem to the one-dimensional case and consider only the right ascension pattern and the effect of the source shape on it.

The east-west rows of 24 helices are as shown in Fig. 5(a) with a uniform spacing $d$ between helices. The total field pattern of this array is given by the product of the individual helix pattern and the pattern of an array of 24 isotropic point sources with a spacing $d$ (array pattern). However, the array pattern is so much sharper than the helix pattern that for angles near the meridian (broadside to the antenna array) total antenna

[^23]pattern is substantially the same as array pattern. The array has a uniform amplitude distribution.

## Case with All Helices in Phase (Single-Lobe Pattern)

Considering first the case where all helices are in phase the normalized field pattern is substantially that of a linear array of 24 isotropic point sources of equal amplitude and spacing as given by ${ }^{2}$

$$
\begin{equation*}
E_{A}=\frac{1}{n} \frac{\sin \frac{n \psi}{2}}{\sin \frac{\psi}{2}} \tag{4}
\end{equation*}
$$

where

$$
\begin{aligned}
\psi & =(2 \pi d / \lambda) \sin \phi_{\psi,} \\
d & =\text { spacing between helices ( }=1.69 \text { wavelengths at } \\
& 250 \mathrm{mc}), \\
n & =24 \text { (number of helices in a row in } \phi \text { direction), } \\
\phi_{0} & =\text { angle from meridian to center of source, } \\
\lambda & =\text { wavelength. }
\end{aligned}
$$

The normalized power pattern of the antenna is equal to $E_{A}{ }^{2}$. This is the true antenna pattern $F(\phi)$ as shown in Fig. 5(b). Let the source distribution be a step function equal to unity between $+\alpha / 2$ and $-\alpha / 2$ and zero elsewhere. For this case (2) then becomes

$$
\begin{equation*}
G\left(\phi_{0}\right)=\frac{1}{\alpha} \int_{-\alpha / 2}^{+\alpha / 2} \frac{\sin ^{2}\left[\frac{n d_{r}}{2} \sin \left(\phi+\phi_{0}\right)\right]}{n^{2} \sin ^{2}\left[\left(d_{r} / 2\right) \sin \left(\phi+\phi_{0}\right)\right]} d \phi \tag{5}
\end{equation*}
$$

where $d_{r}=2 \pi d / \lambda=$ spacing of helices in radians.


Fig. 5-(a) 24-helix array, (b) antenna pattern with assumed source pattern.

Eq. (5) for the observed distribution $G\left(\phi_{n}\right)$ can be obtained from (5) by graphical integration [by measuring shaded area in Fig. 5(b)], or (5) can be evaluated analytically. Proceeding with the latter solution one can put $\sin \left(\phi+\phi_{0}\right)=\phi+\phi_{0}$ provided both $\phi_{0}$ and $\alpha$ are small. Introducing this approximation into (5) and integrating yields

[^24]\[

$$
\begin{align*}
G\left(\phi_{0}\right)= & \frac{1}{n}+\frac{2}{n^{2}} \sum_{m=1}^{n-1} \frac{m \sin \left[(n-m) d_{r} \frac{\alpha}{2}\right]}{(n-m) d_{r} \frac{\alpha}{2}} \\
& \cdot \cos \left[(n-m) d_{r} \phi_{0}\right] . \tag{6}
\end{align*}
$$
\]

By evaluating (6) for various values of $\alpha$ the curves of Fig. 6 were obtained. These show the effect of the source extent $\alpha$ on the observed pattern $G\left(\phi_{0}\right)$ of the O.S.U. radio telescope antenna. These curves have been normalized (maximum value set equal to unity). For small source extent ( $\alpha$ small) the observed pattern is nearly the same as the true antenna pattern $(\alpha=0)$, while for larger source extent the observed pattern tends to conform more to the source shape. If the source distribution is a step function but $\alpha$ is not known, one can deduce its value from Fig. 6 provided $\alpha>\frac{1}{2}$ degree.


Fig. 6-Single-fobe patterns that woukl be observed with 40 -waveleugth broadside array for assumed uniform source distributions of various angular extent ( $\alpha$ ). The sharpest pattern is for the case of a point source $(\alpha=0)$.

If $\alpha<\frac{1}{2}$ degree the observed pattern differs so litule from the antenna pattern that it is impractical to deduce its value unless the source is sufficiently strong for the small amount of broadening to be accurately measured. It is also to be noted in Fig. 6 that with increase in $\alpha$ the minor lobe amplitude of the observed distribution tends to decrease. Since the array is a long one the curves of Fig. 6 apply not only to the 24 -helix array but also approximately to any uniform rectangular broadside array or aperture 40 wavelengths across.

## Case of Two IIalves of Array in Phase Opposition (SplitLobe Pattern)

If the helices to the right of the center of the array are reversed in phase with respect to those to the left, the total field pattern is given closely by

$$
\begin{equation*}
E_{A}=\frac{1}{n^{\prime}} \frac{\sin ^{2} \frac{n^{\prime} \psi}{2}}{\sin \frac{\psi}{2}}=\frac{1}{n} \frac{\left[1-\cos \frac{n \psi}{2}\right]}{\sin \frac{\psi}{2}} ; \tag{7}
\end{equation*}
$$

where
$\psi=$ same as in (4), $n^{\prime}=12$, and $n=24$.

Squaring (7) gives the power pattern. Introducing this in (2) and integrating yields, for $n$ even,

$$
\begin{align*}
G\left(\phi_{n}\right)= & \frac{1}{n}-\frac{2}{n^{2}} \sum_{m=1}^{n-1} m \frac{\sin \left[\frac{(n-m)}{2} d_{r} \alpha\right]}{\frac{n-m}{2} d_{r} \alpha} \\
& \cdot \cos \left\lfloor(n-m) d_{r} \phi_{0}\right\rfloor \\
& +\frac{8(n-2) / 2}{n^{2}} \sum_{m=1}^{m} m \frac{\sin \left[\frac{(n-2 m)}{4} d_{r} \alpha\right]}{\frac{(n-2 m)}{4} d_{r} \alpha} \\
& \cdot \cos \left[\frac{(n-2 m)}{2} d_{r} \phi_{0}\right] . \tag{8}
\end{align*}
$$

livaluating (8) for various values of $\alpha$, curves of Fig. 7 were obtained. These show effect of uniform source distributions of width $\alpha$ on observed pattern. Curves have been normalized. Note that the value of curves in Fig. 7 at minimum occurring at $\phi_{0}=0$ is an effective indicator of angular extent $\alpha$ of source for $\alpha$ values between about $\frac{1}{2}$ degree and 2 degrees.


Fig. 7-Split-lobe patterns that would be olserved with 40 -wavelength broadside array with two halves in phase opposition for assumed uniform source distributions of various angular extent ( $\alpha$ ).

For a source of small extent more large lobes of reduced beamwidth will be obtained as the spacing between the two halves of the antenna is increased. Under these conditions the two halves of the antemna act as the two units of a Michelson interferometer. If the spacing is increased sufficiently the minimum to maximum ratio of the observed lobe amplitude will tend to approach unity. Further increase in spacing will cause fluctuations in the ratio. By using an interferometer with a number of different spacings ${ }^{3,4}$ it is theoretically: possible to deduce the source distribution with approximately the same accuracy as an array (with single-lobe pattern) extending continuously over a distance equal to the largest spacing between the interferometer units.

[^25]
## Case of Comparison Arrangement

In this case the antenna and receiving system are operated in such a manner ${ }^{5}$ that the difference of the single-lobe and split-lobe patterns is obtained. ${ }^{6}$ Fig. 8 shows the curves for this case for various values of $\alpha$. lach curve is obtained by applying a scale factor to a curve of Fig. 7 and then subtracting this curve from corresponding curves of Fig. 6. The scale factor adjusts corresponding curves of Figs. 6 and 7 to the same scale.


Fig. 8-Patterns that would be observed with 40 -wavelength broadside array for the comparison method of operation.

The location of the zeros for the curves of Fig. 8 is independent of receiver gain. This is on advantage. Furthermore the slope is a maximum at the zero points so that these can be used to obtain an accurate time of transit (average of the two zero points), giving precise position data for the celestial source.

## Direct Solution

Finally, in the most general situation where both antenna and source patterns are two-dimensional distributions on a sphere surface, the problem can be stated

$$
\begin{equation*}
G\left(\phi_{0,}, \theta_{0}\right)=\frac{1}{\Omega_{0}} \iint F(\phi, \theta) \int\left(\phi, \phi_{0}, \theta, \theta_{0}\right) d \Omega ; \tag{9}
\end{equation*}
$$

where

$$
\begin{aligned}
G\left(\phi_{0}, \theta_{0}\right) & =\text { olserved or resultant distribution, } \\
F(\phi, \theta) & =\text { true antenna pattern, } \\
f\left(\phi, \phi_{0}, \theta, \theta_{0}\right) & =\text { source distribution, } \\
\Omega_{0} & =\text { equivalent solid angle subtended by } \\
& \text { source. }
\end{aligned}
$$

In (9), the observed and true antenna patterns are usually known, while the source distribution is unknown and is the quantity desired. The source distribution can be obtained indirectly, as done in the special cases above,

[^26]by assuming various source distributions and calculating the corresponding observed distributions. Then by comparison with the actual observed distribution, the source distribution or its equivalent can be decuced.

The source distribution can also be obtained by a direct solution of (9) for $f\left(\phi, \phi_{0}, \theta, \theta_{0}\right)$. This may be done by expanding the patterns into sets of orthogonal functions. For example, in the one-dimensional case the distributions can be expanded into Fourier series. ${ }^{7}$ Thus, the source distribution can be written as

$$
\begin{align*}
f\left(\phi-\phi_{0}\right)= & a_{0}+\sum_{m=1}^{\infty} a_{m} \cos m\left(\phi-\phi_{0}\right) \\
& +d_{m} \sin m\left(\phi-\phi_{0}\right) \tag{10}
\end{align*}
$$

the observed distribution as

$$
\begin{equation*}
G\left(\phi_{0}\right)=b_{0}+\sum_{n=1}^{\infty} b_{n} \cos n \phi_{0}+e_{n} \sin n \phi_{0} \tag{11}
\end{equation*}
$$

and the antenna pattern (symmetrical) as

$$
\begin{equation*}
F(\phi)=c_{0}+\sum_{p=1}^{\infty} c_{p} \cos p \phi \tag{12}
\end{equation*}
$$

Then

$$
\begin{align*}
G\left(\phi_{0}\right)= & \int_{-\pi}^{\pi} f\left(\phi-\phi_{0}\right) F(\phi) d \phi=\int_{-\pi}^{\pi} a_{0} c_{0} d \phi \\
& +\int_{-\pi}^{\pi}\left[\sum c _ { p } \operatorname { c o s } p \phi \sum \left(a_{m} \cos m\left(\phi-\phi_{0}\right)\right.\right. \\
& \left.\left.+d_{m} \sin m\left(\phi-\phi_{0}\right)\right)\right] d \phi \tag{13}
\end{align*}
$$

Integrating,

$$
\begin{align*}
G\left(\phi_{0}\right)= & 2 \pi a_{0 G_{0}} \\
& +\pi \sum_{n=1}^{\infty} c_{m}\left(a_{m} \cos m \phi_{0}+d_{m} \sin m \phi_{0}\right) \tag{14}
\end{align*}
$$

Equating (14) and (11) term by term, the coefficients of the source distriloution are found to be

$$
\begin{equation*}
a_{0}=\frac{b_{0}}{2 \pi c_{0}}, \quad a_{n}=\frac{b_{n}}{\pi c_{n}}, \quad d_{n}=\frac{e_{n}}{\pi c_{n}} \tag{15}
\end{equation*}
$$

A necessary and sufficient condition for there to be a solution for the source distribution is that $b_{n}$ and $e_{n}$ be zero for all $n$ for which $c_{n}$ is zero and that ${ }^{8}$

$$
\begin{equation*}
\sum_{0}^{\infty}\left|\frac{b_{n}}{c_{n}}\right|^{2}+\left|\frac{e_{n}}{c_{n}}\right|^{2}<\infty \tag{16}
\end{equation*}
$$

The significance of the $b_{n}$ and $e_{n}$ coefficients being zere for all $n$ for which $c_{n}$ is zero is to limit the components of the lourier series of the actual observed distribution to those components which are contained in the series expansion of the antenna pattern $F(\phi)$. That

[^27]is, if the series expansion of the source distrilution contains terms which do not appear in the series expansion of $F(\phi)$, then the solution is not unique. Actually the antenna will not respond to variations of the source distribution whose period is less than that recpuired for $\frac{1}{2}$ the beamwidth between first nulls (approximately equal to the half-power beamwidth).

As an example of this method it will be applied to find the source distribution where the antenna pattern $F(\phi)$ and the observed distribution $G(\phi)$ are given. Let the antenna pattern be that of the O.S.U. antenna as shown by the $\alpha=0$ curve of Fig. 6 , and let $G(\phi)$ be the $\alpha=2$ degree pattern in the same figure. Proceeding by the above analytical (Fourier) method the source distribution shown by the solid curve in Fig. 9 was obtained. The assumed distribution is given by a step function (dashed line).


Fig. 9-Assumed source distribution compared with calculated distribution.

The lourier series for $F(\phi)$ and $G(\phi)$ in the above calculation were obtained graphically by the 24 -ordinate method over a range of $\phi$ between +3 degrees and -3 degrees. The calculated distribution ploted in lig. 9 contains terms to the 4 th harmonic of the Fourier series. The higher harmonic terms introduced by the rectangular distribution cannot be determined because of the lack of response of the antenna to these terms. Thus, the antenna tends to smooth out the source variations. ${ }^{9,10}$ However, amplitude and equivalent rectangular extent, as measured at half-power points, are indicated properly by calculated distribution.

It should be noted that for sources of small angular extent where the observed and antenna patterns differ almost imperceptibly it becomes impractical to deduce the source distribution with any certainty (except to state its maximum possible extent).

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[^28]
# Nonsaturating Pulse Circuits Using Two Junction Transistors* 

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#### Abstract

Summary-Junction transistors have been supposed to be too slow for many pulse applications. However, if they are used in a way in which their collector voltage is never permitted to become zero, saturation does not occur and the switching times achieved may be as low as several times the reciprocal of $\omega_{0}$, the radian cut-off frequency of $\alpha$. This time will be less than a microsecond for junction triodes presently available.

Saturation is prevented through use of breakdown diodes which terminate switching transients by their breakdown. They may also serve other functions in the circuit. A binary counter described has stable points dependent upon the breakdown diodes and passive components, these points being virtually independent of the transistor or its temperature-sensitive $I_{c 0}$.

A class of two-transistor pulse circuits is described including a binary counter which with one kind of junction triode operates at 1.25 mc . Monostable and astable circuits of the same general nature are shown. The pulse requirements for switching of the binary counter are analyzed in some detail. It is shown that the switching charge is the significant quantity and that a crude estimate of the charge required is $1 / \omega_{0}$ times the difference in conduction currents of the two transistors in the stable state.


## Introduction

CNCTION transistors have many attributes which are of importance in pulse applications. Their greater designability as devices over point-contact units is important in pulse work as it is elsewhere. The low values of saturation current obtainable in junction units is of particular interest since it is the significant factor determining smallness of power level which can be employed. The principal question usually raised about the application of junction mits in switching is the question of their supposed low speed. Analysis indicates and experiment verifies that pulse circuits using junction transistors currently available can swith in fractions of a microsecond provided that one prevents the collector voltages from going to zero at the terminus of each switching operation. It has been shown ${ }^{2}$ that this condition of saturation floods the base region with minority carriers. The recovery from saturation can be time consuming and can enormously slow operating speeds. R. L. Wallace suggested that one could avoid saturation with two breakdown diodes placed in the circuit in such a way that the switching transient is terminated by their breakdown rather than transistor saturation. ${ }^{2}$ This technique has been employed in a binary counter which operates on pulses occurring at 1.25 me with experimental $n-p-n$ alloy-junction ${ }^{3}$ tran-

[^29]sistors. A similar circuit using grown-jurction ${ }^{4}$ triode transistors operated at a 600 kc rate. In this circuit both transistors are continuously in their active region; they never reach either zero emitter current or zero collector voltage. Thus the role played by the transistors is essentially that of a linear amplifier, the necessary nonlinear functions being performed by breakdown diodes.

Through the use of two additional breakdown diodes one can fix the stable points of a bistable circuit virtually inclependent of transistors or of the temperature sensitive $I_{r 0}$ (saturation current) of the transistors. Moreover, similar techniques to those applicable to bistable circuits can be applied to monostable and astable circuits with the result that similar transition times are obtained and the pulse height can be accurately set.

## Two-Transistor Pulse Circuits

The elements shown in Fig. 1 are all found in the vast majority of two-transistor pulse circuits. Additional elements may be found in some embodiments but those shown are essential.


Fig. 1-Usual form of two-transistor pulse circuit.
All pulse circuits have one characteristic in common. It is that in certain conditions they are unstable. This simply means that for these conditions transient currents and voltages exhibit a growing or self-perpetuating behavior. In pulse circuits the attribute of instability is associated with the self-completion of a switching operation once the circuit is triggered appropriately. The process of triggering to initiate the switching operation successfully involves bringing the circuit into its unstable region and leaving initial conditions at the end of the switching pulse such that the circuit completes whatever is left of the switching transient "on its own."

[^30]With reference to Fig. 1, the instability mentioned above is physically apparent. The circuit, like all trigger circuits of the Eccles-Jordan variety, is essentially a two-stage amplifier with positive feedback.

A second characteristic common to all pulse circuits is that growing transients are always terminated by a change in the characteristics of some component. In the circuits to be described here this component is a breakdown diode (not shown in Fig. 1). The volt-ampere characteristic of an idealized breakdown diode is shown in Fig. 2. In region $b$ the diode approximates an open circuit, but as the voltage is brought to the breakdown point $V_{b d}$, the diode's incremental resistance approximates a slort circuit. In the circuits described in the following it is the diode voltage attaining $V_{b d}$ which terminates the unstable transient. For these circuits when the operating point of the diode is in range " $a$ " the circuits are stable, that is, all of their natural modes


Fig. 2-Volt-ampere characteristic of idealized breakdown diode.
correspond to decaying transients. When the operating point is in range " $b$," at least one of the natural modes corresponds to a growing transient. The remainder of the circuit is essentially linear. In fact two sets of linear analyses, one for the diodes as an open circuit and one for the diodes as short circuits give results in substantial agreement with experiment.

With respect to the pulse circuit shown in Fig. 1, breakdown diodes connected from the collector terminals $c_{1}$ and $c_{2}$ to $B+$ or ground can serve to terminate a growing transient. In general the growing transients of circuits of this type tend to increase one collector voltage and decrease the other. Thus a diode connected to ground from a collector prevents the increase of collector voltage beyond its breakdown point. 'This prevention ordinarily limits the corresponding decrease in collector voltage of the companion transistor. In a similar way a breakdown diode connected from the collector to $B+$ in the proper orientation terminates the fall of collector voltage at $V_{B}-V_{b d}$. The connection of two breakdown diodes in the arrangement shown in Fig. 3 placed between $c_{1}$ and $c_{2}$ in Fig. 1 terminates switching transients when the difference in collector voltages rises to $V_{b d}$ in either direction. (This suggestion was due to R. L. Wallace.) This particular arrangement has the desirable feature that it can also be used for the additional function (described in detail later) of "pulse routing" in a binary counter.

The diagram of Fig. 1 can represent either bistable, monostable or astable circuits. Moreover, it is possible to employ breakdown diodes in all of these types to
prevent the transistors from going into the saturated condition. The distinction between these is simply whether or not the circuit maintains the broken-down diode in its low impedance condition once the switching has taken place. Bistable circuits can remain permanently in either of two states, monostable circuits can remain permanently in only one state and astable circuits have no stable condition in which they will remain.


Fig. 3-Connection of two breakdown diodes and volt-ampere characteristic.

In any case switching, or unstable, transients occur which terminate with the breakdown of a diode. Until the diode breaks down, at least one of the natural frequencies of the circuit corresponds to a growing transient. Once the diode has broken down to terminate the switching, all of the natural frequencies correspond to decaying transients. If the circuit is not to remain in this stable condition, the current in the diode decays to zero and at this point a new switching transient ensues.


Fig. 4-A nonsaturating binary counter.

## Analysis of Binary Counter

A two-transistor monsaturating binary counter of the general type described in the foregoing section is shown in Fig. 4. In the two stable states the circuit will maintain different collector voltages by the amount of $V_{b_{s}}$, that is, one of the two diodes $D_{1}$ and $D_{2}$ will be broken down, the other conducting in the forward direction. Diodes $D_{3}$ and $D_{4}$ are continuously broken down and hence with their bypass condensers maintain constant
potential drops between the points $c_{1}$ and $b_{2}$ and the points $c_{2}$ and $b_{1}$. The collector voltage of the high-conducting transistor will be $l^{\prime} b_{c}-V_{b s}(>0)$ and the collector voltage of the low-conducting transistor will be $V_{b c}+V_{b s}$. When a switching pulse of either polarity is applied at point $B$, transistors $T_{1}$ and $T_{2}$ exchange the roles of high and low conduction. 'lous the circuit acts as a binary connter in that either transistor assumes a given state alternately as pulses are applied. Binary conters ordinarily require auxiliary diode circuits called ronting circuits which direct an incoming pulse to the appropriate point to turn on the low-conducting unit and to turn off the high-conducting unit. The diodes $D_{1}$ and $D)_{2}$ serve this function in addition to their function of preventing saturation of transistors $T_{1}$ and $T_{2}$.

The interesting properties in the application of himary counters are the switching time, pulse requirements for switching and the dependence of the counter's behavior on changes in the transistor with temperature or exchanges of transistors. These properties can be evaluated for the circuit of Fig. 4 quite directly. Moreover the same analysis with only minor change applies to monostable and astable circuits as well.

## Analysis of Switching Time of Binary Countiers

The speed with which the switching of a binary comoter is effected is substantially dependent upon the natural frequencies of the circuit when it is in the unstable state. A switching pulse will have been applied and will have cansed the diodes $D_{1}$ and $D_{2}$ (Fig. 4) both to become open circuits in a manner considered in detail in the next section. At the point that they heome open circuits (arrive at range $b$, Fig. 2) the circuit will have certain initial conditions which don't change instantaneously. 1 lence the evahation of the switching time of the circuit of lig. 4 is really a computation of its transient response employing the initial values of the pertinent variables. 'This analysis can be particularly simple if one makes approximations in the representation of the components. Subsequent to the analysis in the simple form the accuracy of the approximations can be assessed. An approximant to the circuit of Fig. 4 for transients is shown in Fig. 5. In Fig. 5, the transistor is approximated ats having negligible base and emitter resistance and collector conductance. The differential equation applied to the collector current [see Fig. $5(\mathrm{~b})$ ] is a first approximation to the diffusion equation for the tramsistor. The alpha cut-off frepuency is $\omega_{0}$ radians per second. From the equation in Fig. 5 one sees that collector currents cannot change instantaneously with finite emitter current. Hence for the circuit starting a transient, the pertinent initial conditions are the collector currents, $i_{c 1}$ and $i_{c 2}$ and the voltages on the condensers connected to the emitters.

The circuits shown in both Fig. 4 and Fig. 5 are symmetrical. In such circuits one can effect an economy in analysis through considering symmetrical and anti-
symmetrical components of the variables. 'Thus we define:

$$
\begin{align*}
& v_{c s}=\frac{v_{c 1}+v_{c 2}}{2}
\end{align*} \quad v_{c h}=\frac{v_{c 1}-v_{c 2}}{2},\left\{\begin{array}{l}
i_{c 1}+i_{c 2} \\
i_{c b}=\frac{i_{c 1}-i_{c 2}}{2} \tag{1}
\end{array}\right\}
$$

Noreover, all initial conditions and excitations are split into symmetrical and anti-symmetrical components in the same way. The economy arises since symmetrical excitations, initial conditions and variables are independent of the anti-symmetrical ones and the equilibrium equations can be written separately for each, each involving only half the total number of variables. In addition, in considering switching in pulse circuits one is principally interested in the anti-symmetrical components as these are the only ones involved when one side of the circuit changes with respect to the other.


Fig. 5-Approximant of circuit of Fig. 4.
Expressing the fact that current leaving each emitter flows to the adjacent $G-C$ circuit of Fig. 5. one has,

$$
\begin{equation*}
v_{c 1} G_{1}+C_{1} \frac{d v_{c 1}}{d t}=-\frac{i_{c 1}}{\alpha_{0}}-\frac{d i_{c 1}}{d t} \frac{1}{\alpha_{0} \omega_{0}} \tag{2a}
\end{equation*}
$$

and

$$
\begin{equation*}
\hat{i}_{c 2} G_{1}+C_{1} \frac{d v_{c 2}}{d t}=-\frac{i_{c 2}}{\alpha_{0}}-\frac{d i_{c 2}}{d t} \frac{1}{\alpha_{0} \omega_{0}} . \tag{2~b}
\end{equation*}
$$

Subtracting, using (1),

$$
\begin{equation*}
v_{c a} G_{1}+C_{1} \frac{d v_{c a}}{d t}=-\frac{i_{c a}}{\alpha_{0}}-\frac{d i_{c a}}{d t}-\frac{1}{\alpha_{0} \omega_{0}} \tag{2c}
\end{equation*}
$$

Expressing the fact that the sum of currents leaving the cross-coupling wires must add to zero, one has,

$$
\begin{equation*}
i_{1}\left(C_{1}+G_{2}\right)+C_{1} \frac{d i_{1}^{\prime}}{d t}-i_{c 2}+i_{L_{1}}=0 \tag{3a}
\end{equation*}
$$

and

$$
\begin{equation*}
i_{c 2}^{\prime}\left(G_{1}+G_{2}\right)+C_{1} \frac{d v_{c 2}}{d t}-i_{c 1}+i_{c 2}=0 \tag{3b}
\end{equation*}
$$

Subtracting, using (1),

$$
\begin{equation*}
\vartheta_{c a}^{\prime}\left(G_{1}+G_{2}\right)+C_{1} \frac{d v_{c a}}{d t}+2 i_{c a}=0 \tag{3c}
\end{equation*}
$$

From (2c) and (3c) one obtains the characteristic equation

$$
\begin{align*}
p^{2}+p\left[\left(1-2 \alpha_{0}\right) \omega_{11}+\frac{G_{1}+G_{2}}{C_{1}}\right]+\frac{\omega_{0} G_{1}\left(1-2 \alpha_{0}\right)}{C_{1}} \\
+\frac{\omega_{0} G_{2}}{C_{1}}=0 . \tag{4}
\end{align*}
$$

If $G_{2}$ were zero, the roots of (4) become $\left(2 \alpha_{0}-1\right) \omega_{0}$ and $-\left(G_{1} / C_{1}\right)$, the first being associated with the growing transient, the latter being simply associated with the discharge of the re circuit at the emitter. However, if $G_{2}$ is present but much smaller than $G_{1}$, a normal situation, one finds the roots to be approximately:

$$
\left.\begin{array}{l}
p_{1} \cong\left(2 \alpha_{0}-1\right) \omega_{0}-\frac{\frac{2 \alpha_{0} \omega_{0} G_{2}}{C_{1}}}{\omega_{0}\left(2 \alpha_{0}-1\right)+\frac{G_{1}}{C_{1}}} \\
\text { (associated with growing component) } \\
p_{2} \cong-\frac{G_{1}}{C_{1}}+\frac{\frac{G_{2}}{C_{1}}\left(\omega_{0}-\frac{G_{1}}{C_{1}}\right)}{\left(2 \alpha_{0}-1\right) \omega_{0}+\frac{G_{1}}{C_{1}}} \\
\text { (associated with decaving component). }
\end{array}\right\}
$$

Finally where $\left(G_{1} / C_{1}\right) \ll \omega_{0}$ and $\alpha_{0} \rightarrow 1$, the roots of the characteristic equation become approximately

$$
\left(2 \alpha_{0}-1\right) \omega_{0}-2 \frac{G_{2}}{C_{1}} \quad \text { and } \quad-\frac{G_{1}}{C_{1}}+\frac{G_{2}}{C_{1}}
$$

The switching time is primarily dependent upon the exponent of the growing component of the transient, $\left(2 \alpha_{0}-1\right) \omega_{0} t$ or about $\omega_{0} t$. The time for the growing component of the transient to multiply itself by $\epsilon$ from the value left by the switching pulse is about $1 / \omega_{0}$ second, which is less than $0.1 \mu \mathrm{sec}$ for grown-junction triode transistors and smaller by several times in some of the fastest units. The time required for the growing exponential to increase by 10 times is $2.3 / \omega_{0}$ seconds since $\epsilon^{23}$ is 10 .

At the termination of the switching period, the growing component of the transient is ordinarily the largest component. The time required for it to grow to its size at the termination from 0.1 of that value is a good approximation of the rise time of the output pulse or the switching time. Thus the rise time should be about
$2.3 / \omega_{0}$ seconds. Further, the analysis points up the fact that $G_{2}$ should be as small as possible and that the emitter resistors should be bypassed, conclusions toward which one is led if he considers the circuit as a two-stage a mplifier with positive feedback in which loop amplification and bandwidth should be maximized.

The circuit representation shown in Fig. 5 is optimistic in that base and emitter resistances are approximated by short circuits. The influence of these resistors is to slow the transients a bit. The resistors $R_{3}$ in Fig. 4 are essentially in parallel with the two $R_{2}$ 's and these parallel combinations should be identified with $G_{2}$ in Fig. 5.

The switching time is dependent not only upon the natural frequencies of the circuit in its unstable state, but also upon the initial conditions left by the driving pulse. A strong pulse naturally causes the circuit to switch somewhat faster than a weak one, simply because the larger initial value of the switching transient requires less time to grow to the size required to establish the opposite stable state. More than this it is essential that the driving pulse leave initial conditions which will ultimately result in the opposite stable state to that which was in effect before the pulse. Here the final state is associated with the growing transient which with time dominates the decaying one, and one must insure that the sign of the growing component is in the proper direction. Referring to Fig. 5, one recognizes from the earlier discussion that $G_{2}$ should be small in comparison with $G_{1}$.

If one considers the limiting ase when $i_{2}=0$,

$$
\begin{equation*}
-i_{\epsilon 1}=i_{r 1}-i_{c 2}=\frac{i_{c 1}}{\alpha_{0}}+\frac{d i_{c 1}}{d l} \frac{1}{\alpha_{0} \omega_{0}}, \tag{6a}
\end{equation*}
$$

and

$$
\begin{equation*}
-i_{\epsilon 2}=i_{c 2}-i_{c 1}=\frac{i_{c 2}}{\alpha_{0}}+\frac{d i_{c 2}}{d t} \frac{1}{\alpha_{0} \omega_{0}} \tag{6b}
\end{equation*}
$$

Subtracting and using (1),

$$
\begin{equation*}
-i_{e a}=2 i_{c a}=\frac{i_{c a}}{\alpha_{0}}+\frac{d i_{c a}}{d l} \frac{1}{\alpha_{0} \omega_{0}} \tag{6c}
\end{equation*}
$$

Since (6) includes only the variable $i_{c a}$ one can solve it separately finding that the transient is of the form

$$
\begin{equation*}
i_{c a z}=A \epsilon^{\left(2 \alpha_{0}-1\right) \omega_{0} t} \tag{7}
\end{equation*}
$$

On the basis of (7) one concludes that the asymmetrical component of collector current will, after the beginning of the transient, simply increase; it always keeps the direction of the initial value. The current fed into the rc branch at the emitter terminals is proportional to $i_{c a}$. The voltage appearing there includes a component growing with this current.

The requirements on the switching pulse are now clear. It must simply cause transistor to be switched to high conduction to carry a larger collector current at termination of pulse than the companion transistor.

The excess grows as the transient develops. The approximation of $G_{2}$ by an open-circuit is somewhat optimistic. Since the current taken by it is small compared to that taken loy $G_{1}$, the simple result obtained above is not seriously incorrect.

## The Mechanism of Switching

The preceding section has identified the role of the switching pulse to be the estallishment of higher conduction in the transistor which after the switching transient will be the high conducting one. A complete analysis of the switching mechanism is quite complicated, but a semi-quantitative analysis of simplified circuits can show what the important factors are and attach some numerical values to them.

The circuit of Fig. 4 can be approximated by that shown in lig. 6. The cross-coupling breakdown diodes are represented by batteries. The anti-saturation diodes are approximated by open circuits when not broken down and by batteries of voltage $V_{b s}$ when broken down. A simplified model of the transistor is employed as shown in Fig. 6.


Fig. 6-Approximant of binary counter.
Before the switching pulse is applied at point $B$, one of the breakdown diodes $D_{1}$ or $D_{2}$ will be broken down, the other will be carrying a current in its forward direction. It is assumed that the pulse source supplies no current to point $B$ in the quiescent condition. For this simple model one can evaluate the stable point voltages and currents easily. It is convenient to consider symmetrical and asymmetrical components separately. The only asymmetrical voltage source is that represented by the diode which is broken down. For definiteness we take this to be $D_{1}$. The symmetrical component of collector current is

$$
\begin{equation*}
i_{c s}=\frac{V_{b}-V_{b c}}{\frac{1}{G_{1}}+\frac{1}{G_{2}}} . \tag{8}
\end{equation*}
$$

The asymmetrical component of collector current is

$$
\begin{equation*}
i_{c a}=\frac{V_{b z} G_{1}}{2} \tag{9}
\end{equation*}
$$

The current flowing in $D_{1}$ is

$$
\begin{equation*}
i_{D 1}=\frac{V_{b_{s}}}{2}\left(G_{1}-G_{2}\right)=-i_{D 2} . \tag{10}
\end{equation*}
$$

The discussion of the switching mechanism at the application of a pulse is easier after one considers what happens as the voltage $e_{p}$ is slowly changed. If $e_{p}$ is slowly increased, $i_{p}$ simultaneously increases and this current divides equally between diodes $D_{1}$ and $D_{2}$. The admittance seen by the source $e_{p}$ is $2\left(G_{1}+G_{2}\right)$. Thus one sees that total current taken by diode $D_{1}$ increases while that taken by $D_{2}$ decreases. This situation persists as $e_{p}$ is continuously raised until diode $D_{2}$ is carrying no current, whereupon it opens, severing the connection between $B$ and $C$. Further increases in $e_{p}$ beyond this point result in a decrease of $i_{p}$, since the admittance seen at $B$ w:th $D_{2}$ open is $-\left(G_{1}{ }^{2} / G_{2}\right)+G_{2}$. If one increases $e_{p}$ further, the current $i_{p}$ goes to zero and diode $D_{1}$ opens removing the pulse source from the circuit. At that point the two collectors carry equal currents. The circuit is in an unstable condition at this point. In the normal switching case one desires the unstable transient which ensues to cause $D_{2}$ to break down. However, for this "slow-motion" case at the severing of the switching source $e_{p}$, both transistors carry the same current and the direction of the subsequent switching transient is uncertain. The presence of appropriate condensers $C_{1}$ modifies the above description in the case of faster changing switching voltage, making it possible for a switching pulse to leave higher current in the transistor which initially carried the smaller current.

It is simple to trace the sequence of events which occur when with $D_{1}$ broken down a step-of voltage is applied at $B$ of Fig. 6. The condensers $C_{1}$ begin to charge. $D_{2}$ opens as the current impulse brings its total current to zero. Thus $i_{b 2}$ exhibits an impulse as does $-i_{c 2}$ charging the right condenser. The impulse of emitter current causes $i_{c 2}$ to jump to a larger value. Since $G_{2}$ is a small conductance, the principal part of $i_{c 2}$ is drawn from the base of the left transistor. The result is that $i_{e 1}$ exhibits a positive increment and $i_{c 1}$ develops a negative increment which decreases the current in the diode $D_{1}$ ultimately bringing it to zero. At this point the circuit is free of the switching source and executes its own natural transient as described earlier.

Under the assumption that $G_{2}$ approaches zero, one finds the transients sketched in Fig. 7 for a range of different circuit parameters and switching pulses. The transients shown there are evaluated from the beginning of the switching pulse for a circuit representation which does not change during the pulse. In the actual circuit the switching transient is always terminated by the breakdown of one or both of the diodes, $D_{1}$ and $D_{2}$, whereupon the circuit is no longer unstable and a new set of decaying transients occurs. The description at the bottom of the columns in Fig. 7 indicates when the switching transient terminates. If the variables are considered in the order presented, the relationships given in


Fig. 7-Switching pulse shapes.

Fig. 7 are easily verified. From a study of Fig. 7 and consideration of other similar plots one can draw several interesting conclusions.

1. The switching pulse must supply a sufficient impulse to the low-conducting transistor's base to increase its collector current to a value which exceeds that in the companion transistor. A unit impulse of current, a coulomb of charge, applied to a transistor base causes a step in collector current of $\omega_{0}$ amperes. An increase in the size of the emitter condenser $C_{1}$ makes it possible to supply this impulse with lower values of switching voltage and less energy. However, increasing $C_{1}$ while one maintains $G_{1}$ fixed increases the time constant of the emitter branch and results in slower decay of some transient components. Transistors with high values of alpha-cut-off frequency $\omega_{0}$ are inherently easier to switch than poorer units. A simple but very approximate relationship for the minimum charge required for switching is

$$
\begin{equation*}
q=\frac{V_{b_{s}} G_{1}}{\omega_{0}} \tag{11}
\end{equation*}
$$

From (9) one recognizes that this is the amount of charge which if passed through the low-conducting transistor's base will bring its collector current to the value possessed by the high-conducting transistor.
2. The source of the switching pulse may be severed from the counter either by its own high impedance after its impulse has been delivered or it may be severed loy the opening of the second anti-saturation diode. The self-severing is enhanced by values of $C_{1} / G_{1}$ which are several times larger than $1 / \omega_{0}$ or by values of switching pulse greater than $V_{b s} / 2$ which gives a linear decrease in $i_{p}$ with time.
3. Switching failures can result from two kinds of improper pulsing. Too large pulses, greater than $V_{b s}$, for example, may cause both $D_{1}$ and $D_{2}$ to break down and after the pulse is over the counter becomes uncertain as to which way to switch. A similar kind of difficulty is experienced with smaller pulses if during the pulse duration $v_{d 2}$ (see Fig. 7) decreases far enough to induce breakdown of $D_{2}$. The second type of failure results from an unsufficient pulse which after initially opening $D_{2}$ (Fig. 7) fails to get the second diode open and permits $J_{2}$ to begin conducting again in the forward direction.

## Efficts of Temperature Changes on Circuit of Fig. 4

An interesting and useful property of the binary counter shown in Fig. 4 is that the stable voltages with respect to ground of points $c_{1}, \epsilon_{1}, c_{2}$ and $\epsilon_{2}$ are essentially independent of the transistor parameters and saturation current $I_{c 0}$, being dependent upon the resistors and
breakdown diodes and the applied voltage. This would be exactly so if there were no voltage drops between base and emitter of a conducting transistor and if the breakdown diodes have zero incremental impedances, but it is nearly so in any practical case since the emitterbase impedance of a transistor and the conduction impedance of a breakdown diode are small fratctions of the other impedances in the circuit. The reason for the independence is easily explained. An examination of Fig. 6 in which the emitter-base voltages are neglected and the voltages of the diodes $D_{1}$ and $D_{2}$ are taken as $V_{b s}$ and zero reveals that the voltages across the resistors are determined without any further characterization of the transistors. The transistors simply determine how the current divides between them and the crosscoupling diodes without influencing conditions in the resistors at all.

## Experimental Results

A number of different two-transistor pulse circuits have been built and operated. Oscillograms of interesting voltages of a number of them are presented in the remaining figures. In particular the circuits described include a binary counter operated by a blocking oscillator, a two-stage binary counter, the second stage being driven by the first and a monostable or astable twotransistor pulse circuit which is also nonsaturating.


Fig. 8-13inary comnter driven by blocking oscillator.
Fig. 8 gives the circuit arrangement of a binary counter driven by a blocking oscillator which operated at a maximum of 600 kilocycles using grown-junction triodes. Fig. 9 shows a number of oscillograms of the pertinent voltages when the counter is driven by a 200kilocycle wave from the blocking oscillator. Figs. 9(a), 9 (b) and $9(\mathrm{~d})$ are taken with the same attenuator setting on the oscilloscope and with the sweep synchronized with the blocking oscillator is about 2 volts with the peak current from the blocking oscillator being about 3 ma as seen from Fig. 9(b). The switching charge


Fig. 10-Circuit arrangement and oscillograms for two-stage binary counter.
junction triodes the counter will not operate with $150 \mu \mu \mathrm{f}$ condensers and with the alloy-junction transistors the counter cannot be driven at 1 mc with $470 \mu \mu \mathrm{f}$ condensers. Naturally it operates with less driving pulse amplitude at lower frequencies with the larger condensers than is possible with the $150 \mu \mu \mathrm{f}$ condensers.

Fig. 10 shows the circuit arrangement and oscillograms for a two-stage binary counter of the type shown in Fig. 8. The counters are coupled by a diode network which prevents operation of the second binary counter on down swings of the voltage at $B$.

is very crudely estimatrd to be $1.5 \times 10^{-9}$ coulombs which compares with $3.2 \times 10^{-10}$ coulombs obtained from the formula of (11) if one uses the value of $10^{7}$ as the radian cut-off frequency of the grown-junction transistors. Fig. 9(d) applies to a binary counter in which grown-junction transistors are replaced by $n-p-n$ alloyjunction transistors and the blocking oscillator frequency is increased to 1 megacycle. Fig. 9 (c) has a 2 to 1 change in oscilloscope attenuator. It is necessary to observe also that the $470 \mu \mu \mathrm{f}$ condensers shown in Fig. 8 were replaced by $150 \mu \mu \mathrm{f}$ condensers. With grown-

(a)


SCOPE SETtings SAME for (b), (c), (d), (e), AND (f)
FOR (b) AND (c)
$V_{1}=8, \quad V_{2}=10 V \quad V_{3}=15 \mathrm{~V} \quad C_{2}=100 \mu \mu \mathrm{~F}$

FOR (d) ANO (e)
$V_{1}=8, \quad V_{2}=10 \mathrm{~V} \quad V_{3}=15 \mathrm{v} \quad C_{2}=200 \mu \mu \mathrm{~F}$

FOR ( $f$ ) BLOCKING OSCILLATOR REMOVED
$V_{1}=9 \mathrm{~V} \quad V_{2}=10 \mathrm{~V} \quad V_{3}=15 \mathrm{~V} \quad C_{\bar{c}}=200 \mu \mu \mathrm{~F}$
THE GIRCUIT BECOMES FREE FUNNING

(e)

(f)

Fig. 11-Circuit diagram and oscillograms for nonsaturating astable or monostable circuits.

Most of the foregoing material has been applied directly to binary counters. The same general ideas, nonsaturation, switching, speed, etc., apply nearly unaltered in the case of monostable and astable circuits of similar configuration. Fig. 11 (preceding page) shows a circuit which can be made monostable or astable by proper selection of the voltage $V_{1}$. This circuit, which is not symmetrical, differs from the binary counter in two principal respects. First, one of the cross couplings is capacitive, and this accounts for the maximum of one stable state. Second, saturation is prevented by two diodes, both point diodes in the example illustrated, connected to one of the collector terminals. One, $D_{1}$, prevents the collector voltage on the adjacent transistor from dropping to zero. The second, $D_{2}$, prevents a sufficient rise of the collector voltage of the adjacent transistor to drive the collector voltage of the companion transistor off. Both of these diodes could have been breakdown diodes of the appropriate breakdown voltage. Both would have been reversed in orientation, $D_{1}$ being connected to $B^{+}$, $D_{2}$ being connected to ground. The third diode, $D_{3}$, is used simply as a coupling element, which, in combina-
tion with $D_{1}$ and $D_{2}$, assures that the output voltage wave at $c$ has tops and bottoms independent of the transistors. Since at the time of switching, the equivalent representation of this circuit is essentially the same as that of the binary counter, one expects similar switching times and experimental results verify this.

When $V_{1}$ is set to zero or to a value close to zero, diode $D_{2}$ will be broken down and the circuit is in a stable state with

$$
\begin{align*}
& V_{C G}=V_{3}-V_{D 3} \text { and } \\
& V_{D 3}=8.5 \mathrm{v} \text { for diodes } D_{3} \text { used. } \tag{13}
\end{align*}
$$

A positive pulse at $A$ opens $D_{2}$ and the circuit switches $T_{1}$ on more heavily, switching being terminated by breakdown of $D_{1}$. The condenser $C_{2}$ cannot maintain the voltage of point $D$ indefinitely and when the voltage at $D$ falls to a value which permits $D_{1}$ to open, the circuit switches to its original stable state.

When the voltage $V_{1}$ is raised sufficiently that $D_{2}$ is opened the circuit becomes astable with an output wave of amplitude $V_{3}-V_{2}$ and with pulse length being principally determined by $C_{2}$ and the connected resistors.

# A Two-Emitter Transistor with a High Adjustable Alpha* 

R. F. RUTZ $\dagger$


#### Abstract

Summary-The current amplification, $\alpha$, of a point contact transistor can be increased to values in excess of 20 by the addition of a third point contact which is biased so as to act as an emitter. The amount of increase in $\alpha$ can be adjusted by varying the second emitter current. A qualitative explanation of the $\alpha$ enhancement is discussed which involves an internal positive feedback action that varies the hole transport factor, $\beta$, associated with the second emitter as the first emitter current is changed. The effect of varying the second emitter-to-collector spacing is discussed and experimental results are given.


TIHE CURRENT amplification, $\alpha$, of a point contact transistor can be materially increased and controlled by adding a third whisker placed far from the collector and biased so as to act as an emitter. This paper will present a discussion of the mechanism whereby this $\alpha$ enhancement is achieved and describe the characteristics of some experimental two-emitter transistors. Occasionally, in the literature, mention has

[^31]been made of other types of $\alpha$ variations brought about by the use of additional electrodes or whiskers. ${ }^{1}$

The operation of a conventional point contact transistor has been discussed by Shockley. ${ }^{2}$ This transistor consists of a small block of N -type germanium and three electrodes attached to it. These electrodes are an ohmic connection known as the base and two-point contacts, placed a few thousandths of an inch apart, known as the emitter and collector. In normal operation, the emitter is biased positively and injects holes into the germanium, and the collector is biased negatively and collects these holes. The collection mechanism is such

[^32]

Fig. 1-Collector $V$ - $I$ characteristics for a typical transistor with a second emitter spaced close to the collector.
that holes arriving in the neighborhood of the collector may allow additional electrons to flow from the collector to the base. Thus, a given change in the emitter current may cause a greater change in collector current. The current amplification factor $\alpha$ is defined by:

$$
\alpha=-\left(\frac{\partial I_{c}}{\partial I_{e}}\right)_{V c=\text { const. }}
$$

where the subscripts $e$ and $c$ refer to the emitter and collector respectively and $V_{c}$ is the collector to base voltage. Since only the holes injected by the emitter and arriving at the collector are effective in changing the collector current, it is convenient to express $\alpha$ as the product of three factors:

$$
\alpha=\alpha^{*} \beta \gamma
$$

where

$$
\alpha^{*}=\frac{\partial I_{c}}{\partial I_{c p}}, \quad \beta=-\frac{\partial I_{c p}}{\partial I_{e p}}, \quad \gamma=\frac{\partial I_{e p}}{\partial I_{e}}
$$

and the subscript $p$ means that part of the current carried by holes. Here $\gamma$ is the hole injection efficiency of the emitter, $\beta$ is the transport efficiency, and $\alpha^{*}$ is the intrinsic $\alpha$ of the collector. The emitter of a point contact transistor is normally placed so close to the collector that $\beta$ is virtually unity. The $\gamma$ will not in general exceed unity and may be less. The intrinsic $\alpha$ may be considerably greater than unity.

With this background, we are in a position to consider the effect on the $\alpha$ of a transistor of adding a second emitter. It has been found that the effect depends upon the spacing of the second emitter contact from the collector contact. A closely-spaced second emitter also will have a $\beta$ factor of nearly unity and there will be little direct interaction between the emitters. This has been discussed by Itaegele. ${ }^{3}$ The collector hole current will be the sum of the two-emitter hole currents. The average $\alpha$ of the transistor will not be much affected by current in the second emitter. The $\alpha$ referred to here and subsequently in this paper is defined as:

$$
\alpha=-\left(\frac{\partial I_{e}}{\partial I_{e 1}}\right)_{V_{e}, I_{e 2}=\operatorname{const}}
$$

${ }^{3}$ R. W. Haegele, "A crystal tetrode mixer," Syliania Tech., vol. 2, pp. 2-4; October, 1949.


Fig. 2- $\alpha$ vs $I_{\text {el }}$ for a transistor with a second emitter spaced close to the collector.
where the subscripts $e 1$ and $e 2$ refer to the first and second or added emitter respectively.

Fig. 1 shows collector $V-I$ characteristics of a typical transistor with a second emitter close to the collector. The solid lines are for the case where $I_{o 2}=0$ and correspond to the conventional two-whisker transistor characteristics. The dashed lines are for the case where $I_{e 2}=1.0 \mathrm{ma}$. It can be seen that outside of the region of small collector voltage, a displacement of the lines of constant $I_{e 1}$ takes place, but there is no significant change in separation of adjacent lines. Hence, the average $\alpha$ of the original transistor has not been altered appreciably. It has been observed in cases where the curve of $\alpha$ vs $I_{\theta I}$ has a peak, that this peak may be shifted and somewhat modified when a constant current is applied through the second emitter. Fig. 2 shows an example of this. Here $\alpha$ is shown as a function of $I_{\text {el }}$ for two different values of $I_{e 2}$, namely, $I_{e 2}=0$ and $I_{o 2}=0.5 \mathrm{ma}$. It will be noticed that the peak in the $I_{e 2}=0$ curve is displaced to the left by approximately 0.5 ma in the $I_{e 2}=0.5 \mathrm{ma}$ curve. This may be explained on the basis that the peak in $\alpha$ is due to a peak in the intrinsic $\alpha$ of the collector which occurs at some definite value of collector hole current. When some of the hole current is supplied by the second emitter the peak occurs at a lower value of first emitter current. If the second emitter hole current is increased beyond the value at which the intrinsic $\alpha$ has its maximum, then the peak will no longer appear in $\alpha$.

If a second emitter is placed far from the collector so that in the absence of first emitter current its transport efficiency is small, a large enhancement of the $\alpha$ of the transistor is possible without greatly changing the other transistor characteristics. Fig. 3 shows a typical example of this effect in the form of a set of collector $V-I$ characteristics for an experimental transistor of this type for the two cases of $I_{e 2}=0$ (conventional two-whisker transistor) and $I_{c 2}=8$ ma. It is evident that the main effect, outside of the region of small collector voltage, is an increased separation of the lines of constant $I_{\text {el }}$, which means that the average $\alpha$ of the original transistor has been noticeably increased. The effect is due to an increase in the transport efficiency of the second emitter
as a result of an increase in current through the first emitter.

The following cuualitative explanation of how this may be accomplished has been suggested by R. W. Landauer. ${ }^{4}$ Consider l"ig. 4 which is a schematic representation of the hole and electron flow in the transistor with a second emitter placed a relatively large distance from the collector. The holes injected by the second emitter


Fig. 3-Collector V-I characteristics for a transistor with a second emitter spaced far from the collector.


Fig. 4-1 Iole injection by second emitter spaced far from the collector.
are indicated symbolically by + and those injected by the first emitter indicated by $\oplus$. Electrons are indicated by -. The region near the collector where the amplifying mechanism is concentrated is indicated by the region $A$ and the rest of the germanium block is labelled region $B$. Let a constant current $I_{e 2}$ be flowing into the second emitter and let the collector voltage be held at some constant negative value. Now consider what happens when the first emitter current is turned on. If the intrinsic $\alpha$ of the collector is sufficiently high, then the increase in collector current due to holes from the first emitter arriving in region $A$ must be accompanied by an

[^33]increased electric field in region $B$. Holes from the fringe of the reservoir of holes created by $I_{82}$ will be drawn to the collector by this increased field. The collecting of these holes in region $A$ liberates even more electrons to flow into region $B$. This further increases the electric field there so that still more holes are drawn from the reservoir. In this way, we have a positive feedback mechanism which, for a given $I_{e 1}$, will make the current gain greater than would have been present without the reservoir of additional holes. The depletion of the reservoir by the feedback action is equivaleat to an increase in the $\beta$ factor of the second emitter.

It is observed that the back resistance of the transistor is lowered by a current flowing into the second emitter. This is because a small number of holes will arrive at the collector from the second emitter even when there is no current through the first emitter, and the number arriving will depend on the collector voltage. This means that $\beta$ factor associated with the second emitter depends slightly upon the collector voltage. A typical example of the change in back resistance due to second emitter current is shown by the change in slope of the $I_{e 1}=0$ line in Fig. 3.


Fig. 5-Average $\alpha$ and back resistance as a function of second emitter spacing for different values of second emitter current.

The curves in Fig. 5 show how the increase in average $\alpha$ varies with the second emitter to collector spacing in a typical transistor for three different values of second emitter current. Also shown in Fig. 5 is the decrease in collector back resistance for the same transistor. The average alpha, $\bar{\alpha}$, is the alpha associated with the collector and first emitter averaged over the interval
of first emitter current from 0 to 2 ma , with the collector voltage held constant at -10 v and the second emitter current held at some constant value. The back resistance, $R_{c}$, is the ratio of the collector voltage to collector current when $V_{c}=-20 \mathrm{v}$ and $I_{e 1}=0$ and the second emitter current is held constant. For the special cases where $I_{A 2}=0$ the average $\alpha$ and back resistance are designated as $\bar{\alpha}_{0}$ and $R_{c 0}$ respectively.

It is apparent that for second emitter to collector spacings of less than four-thousandths of an inch, $\bar{\alpha}$ actually decreases. This is because the second emitter gives sufficient collector hole current so that $\alpha^{*}$ has passed its peak. Except for this close spacing, $\bar{\alpha}$ increases as second emitter current increases. At the same time back resistance decreases. In general, a compromise must be made between these two effects. As might be expected, for a given second emitter current, the $\alpha$ enhancement falls off for large spacings.


Fig. 6- $\alpha$ vs $I_{\text {e }}$ for various values of second entitter current for a transistor with a second emitter spaced far from the collector.

Fig. 6 shows a detailed picture of the relationship between $\alpha$ and $I_{e 1}$ for different values of $I_{e 2}$ for a single transistor. The back resistance, $R_{e}$, associated with each value of the $I_{e 2}$ is also shown. In this particular transistor, as has been found to be the case generally, the $\alpha$ enhancement due to $I_{02}$ is greatest in the region of low first emitter current and falls off for high values of $I_{\text {el }}$. This clearly must happen since the $\beta$ of the second emitter cannot increase above unity.

To show what can be realized in practice, the characteristics of four encapsulated transistors with widely spaced second emitters are given in Table I. In these transistors back resistance decreases approximately fifty
per cent and average $\alpha$ about doubles when the second emitter current is raised from zero to 5 ma.

TABLE I
Average $\alpha$ and Coliector Back Resistance for Four Transistors with Widely-Spaced Second Emitters

| Unit No. | Back Resistance (ohms) |  | Average $\alpha$ |  |
| :---: | :---: | :---: | :---: | :---: |
|  | $R_{c 0}\left(I_{e 2}=0\right)$ | $R_{c}\left(I_{e 2}=5 \mathrm{ma}\right)$ | $\begin{gathered} \alpha_{\mathrm{avg}}\left(I_{e 2}\right. \\ =0) \end{gathered}$ | $\begin{aligned} & \alpha_{\mathrm{agy}}\left(I_{o 2}\right. \\ & =5 \mathrm{ma}) \end{aligned}$ |
| A | 20,000 | 10,000 | 2.5 | 5.0 |
| B | 21,000 | 10,000 | 1.8 | 3.7 |
| C | 29,000 | 14,000 | 3.0 | 5.9 |
| D | 50,000 | 35,000 | 2.6 | 4.6 |

In general, it has been found that the greatest $\alpha$ enhancement occurs in transistors which have an initially high value of $\alpha$. This result might be expected since a large $\alpha$ implies a large intrinsic $\alpha$ and lience large sweeping fields. The grestest $\alpha$ enhancement is oltained when the second emitter is so placed that the first emitter lies on a line between it and the collector. It has also been found that best results are obtained with short lifetime germanium. This means that the reservoir of holes created by the second emitter can be located relatively close to the collector and hence easily be affected by the sweeping fields. Finally, it has been found desirable to make the base connection farther from the second emitter than is the collector. A nearer base would presumably tend to drain holes from the reservoir.

In circuit applications, the high $\alpha$ two-emitter transistor has the advantage of permitting higher current gain than is normally obtainable from conventional point contact transistors. Also, it is useful for applications requiring the control or modulation of the gain of an amplifier. Measurements on one transistor in a circuit with a collector load resistor of a few hundred ohms indicate that the rise time of a rectangular current pulse amplified by the transistor is substantially the same with moderate second emitter currents as it is with zero second emitter current (i.e., as it is for the original transistor comprised of the first emitter and collector alone). The rise time of an output pulse at the collector when a rectangular positive current pulse switches the second emitter current from zero to 5 ma , for a stearly first emitter current of a few milliamperes, has been found to be less than a half-microsecond in an experiment on one transistor. Thus, the modulating properties of the second emitter are apparently not limited to excessively low frequencies.

## Acknowlfdgment

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# Internal Feedback and Neutralization of Transistor Amplifiers* 

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#### Abstract

Summary-Transistors are nonunilateral amplifying devices. The most important effects of internal feedback are reflected immittances and potential instability of amplifiers in certain frequency ranges. These phenomena are undesirable in many applications.

Considering the various matrix representations of a two-terminal pair one can show that by connecting appropriate networks in a suitable manner to the active nonunilateral element, the internal feedback of the latter can be removed (neutralized).

Several neutralized transistor arrangements have been investigated experimentally. The circuits are based on an analysis of the properties of transistor feedback parameters at "higher" frequencies. Neutralization throughout relatively wide bands of frequencies can be achieved using simple feedback arrangements. The maximum available power gain of transistor amplifiers is only moderately affected by neutralization.


## Intronuction

ATWO-TERMINAL pair network is unilateral if an excitation applied to one of its terminal pairs produces a response at the second terminal pair, whereas an excitation applied to the second terminal pair does not result in a response at the first terminal pair, or vice versa. Networks with biclirectional transmission between terminal pairs are nonunilateral; they may be bilateral (if they obey the theorem of reciprocity) or nonbilateral. ${ }^{1}$

It is well known that vacuum tubes operated at low frequencies can be considered as unilateral devices. Transistors, however, are nonunilateral: the three transistor configurations (common base, emitter and collector) exhibit bidirectional transmission between input and output terminals pairs. Although, in general, the "backward" transmission is considerably lower than the "forward" transmission, its effects on the circuit properties of the transistor cannot be neglected. The major effects of the backward transmission (or "internal feedback") on circuit behavior are:

1. The input and output driving point immittances are functions of the load and source immittances respectively.
2. Internal feedback may lead to instability of amplifiers and circuits may oscillate even in the absence of an external feedback loop.
These phenomena are undesirable in numerous transistor applications and of ten represent major diffi-

[^34]culties to the circuit designer and user. In other cases, these properties may be of no consequence or are even desirable.

The problem of internal feedback has been studied by Mason, ${ }^{2}$ who has shown that unilateralization can be achieved by lossless reciprocal coupling. The approach of this paper is somewhat different. The discussed methods of neutralization involve both resistive and reactive elements. This may result in a sacrifice of power gain, but of ten "simplifies" the neutralizing network and makes it easier to design an amplifier neutralized throughout a relatively wide band of frequencies. The discussion is mainly concerned with transistor amplifiers; the principles can, however, be applied to other active devices.

## Reflected Immittances

The behavior of any transistor configuration, considered as an active linear two-terminal pair, can be described by one out of six possible sets of two linear equations. Using the "series-parallel" representation for the generalized transistor amplifier of Fig. 1 one can write:


Fig. 1-Schematic representation of transistor amplifier.

The source impedance is $Z_{G}=R_{G}+j X_{G}$ and the load admittance is $Y_{L}=G_{L}+j B_{L}$. The parameters $h_{i j}$ are functions of the $d c$ operating point, the signal frequency and are, of course, different for different transistor configurations. The existence of internal feedback manifests itself by $h_{12} \neq 0$. The input impedance $Z_{i}$ is a function of the load admittance:

$$
\begin{equation*}
Z_{i}=h_{11}-\frac{h_{12} h_{21}}{h_{22}+Y_{L}} \tag{2}
\end{equation*}
$$

In a similar manner, the output admittance $Y_{0}$ is a function of the source impedance:

[^35]\[

$$
\begin{equation*}
Y_{0}=h_{22}-\frac{h_{12} h_{21}}{h_{11}+Z_{G}} \tag{3}
\end{equation*}
$$

\]

Due to "reflected immittances" a transistor amplifier stage cannot be designed as an isolated unit. To obtain proper performance, the designer must take into account the effect of adjacent and of ten even that of more remote amplifier stages. The internal feedback of transistors makes it difficult to use them in various applications, e.g., in certain types of laboratory equipment.

The problem of reflected immittances is particularly serious in the case of high frequency amplifier stages, where transistor parameters and external circuit elements are complex quantities. Fig. 2 shows the variation of input and output impedances of a common emitter amplifier with tuned input and output as functions of frequency. The complicated nature of the reflected immittances results in a distortion of the bandpass characteristic of tuned amplifiers (nonsymmetrical bandpass).


Fig. 2-Input and output impedances of the common emitter stage as functions of frequency with load and source tuned to 500 kc respectively.

Furthermore, it is not only difficult to design and compute the performance of a multistage tuned amplifier for desired bandshape and gain, but even the alignment of such an amplifier may be a laborious task. The design and alignment of multistage staggered tuned amplifiers present particularly difficult problems.

It can, therefore, be stated that reflected immittances are undesirable in many circuit applications and their elimination could solve numerous design problems

## Stablelty Considerations

The transducer gain of the transistor amplifier (defined as the power delivered to the load divided by the available power of the source) as calculated from (1), is:

$$
\begin{equation*}
G=\frac{4\left|h_{21}\right|^{2} R_{G} G_{L}}{\left|\left(h_{11}+Z_{Q}\right)\left(h_{22}^{\prime}+Y_{L}\right)-h_{12} h_{21}\right|^{2}} \tag{4}
\end{equation*}
$$

$h_{12}$ being different from zero, internal feedback may lead to instability even without an external feedback loop. A well-known example of such behavior in vacuum tube circuitry is the tuned plate-tuned grid oscillator, in which the grid-to-plate capacitance of the vacuum tube furnishes the internal feedback that is required for oscillation.

Analogous phenomena exist in transistor circuits. At higher frequencies the transistor parameters $h_{i j}$ are complex quantities:

$$
\begin{aligned}
h_{11} & =h_{11}^{(R)}+j h_{11}^{(I)} \\
h_{22} & =h_{22}^{(R)}+j h_{22}^{(I)} \\
h_{12} h_{21} & =I I=I H_{R}+j H_{I} .
\end{aligned}
$$

It can be shown ${ }^{3}$ that, restricting $R_{G}$ and $G_{L}$ to positive values, instability may occur, provided that:

$$
\begin{equation*}
H_{i}{ }^{2} \geqq 4 h_{11}^{(R)} h_{22}(R)\left(h_{11}^{(R)} h_{22}^{(R)}-I_{R}\right) \tag{5}
\end{equation*}
$$

In inequality (5) does not hold, the transistor configuration is unconditionally stable, whatever load and source immittances are connected to it. Condition (5) can also be written in the form

$$
\begin{equation*}
|H|+H_{R} \geqq 2 h_{11}^{(R)} h_{22}^{(R)} \tag{5a}
\end{equation*}
$$

An analysis of the three transistor configurations in the light of conditions (5) or (5a) shows that they exhibit potential instability throughout wide frequency ranges. This potential instability can be eliminated by neutralization of the internal feedback.

## Fundamental Aspects of Internal Ferdback

A two-terminal pair network is characterized by two voltages, $E_{1}$ and $E_{2}$, and two currents, $I_{1}$ and $I_{2}$. Consequently, depending on which of these four quantities are considered as independent and dependent variables, the behavior of the network can be described by any one of six possible sets of two-linear equations. ${ }^{4}$ Two of these sets (those separating input variables from output variables) are useful if one considers cascaded networks, whereas the four other sets are important from the point of view of internal feedback. Using the matrix notation, these are:

$$
\begin{align*}
& {\left[\begin{array}{l}
E_{1} \\
E_{2}
\end{array}\right]=\left[\begin{array}{ll}
z_{11} & z_{12} \\
z_{21} & z_{22}
\end{array}\right] \times\left[\begin{array}{l}
I_{1} \\
I_{2}
\end{array}\right]}  \tag{6}\\
& {\left[\begin{array}{l}
I_{1} \\
I_{2}
\end{array}\right]=\left[\begin{array}{ll}
y_{11} & y_{12} \\
y_{21} & y_{22}
\end{array}\right] \times\left[\begin{array}{l}
E_{1} \\
E_{2}
\end{array}\right]}  \tag{7}\\
& {\left[\begin{array}{l}
E_{1} \\
I_{2}
\end{array}\right]=\left[\begin{array}{ll}
h_{11} & h_{12} \\
h_{21} & h_{22}
\end{array}\right] \times\left[\begin{array}{l}
I_{1} \\
E_{2}
\end{array}\right]}  \tag{8}\\
& {\left[\begin{array}{l}
I_{1} \\
E_{2}
\end{array}\right]=\left[\begin{array}{ll}
g_{11} & g_{12} \\
g_{21} & g_{22}
\end{array}\right] \times\left[\begin{array}{c}
E_{1} \\
I_{2}
\end{array}\right]} \tag{9}
\end{align*}
$$

[^36]If the matrix elements corresponding to one of the representations (6) to (9) are known, those pertaining to the other sets can be computed. Eqs. (6) to (9) can be written in a symbolic form:

$$
\left[\begin{array}{l}
D_{1}  \tag{10}\\
D_{2}
\end{array}\right]=\left[\begin{array}{ll}
k_{11} & k_{12} \\
k_{21} & k_{22}
\end{array}\right] \times\left[\begin{array}{l}
J_{1} \\
J_{2}
\end{array}\right]
$$

where $D_{i}$ and $J_{i}$ symbolize dependent and independent variables pertaining to terminal pair $i$ respectively. The matrix elements of (10) have the following significance:
$k_{11}$ is an input immitance,
$k_{12}$ is a backward transfer ratio or immittance indicative of intermal feedhack,
$k_{21}$ is a forward transfer ratio or immittance,
$k_{22}$ is an output immittance.
The network is unilateral if the feedlack parameter $k_{1:}=0$ ). Using a matrix conversion table one sees that if $k_{12}=0$ in one of the representations (6) to (9), it is zero in the others. For example, if $h_{12}=0$, than $z_{12}=y_{12}=g_{12}=0$. This, of course, is not surprising, because $h_{12}=0$ implies that an output excitation results in zero input response.

Athough the feedback parameters in the representations (6) to (9) vanish simultancously, it is useful to note that they symbolize different aspects of the internal feedback mechanism.

The transfer impedance $z_{12}$ represents series feedlack proportional to the output current;
The transfer admittance $y_{12}$ represents parallel feed-
back proportional to the output voltage;
The voltage transfer ratio $h_{12}$ represents series feedback proportional to the output voltage; and
The current transfer ratio $g_{12}$ represents parallel feed-
back proportional to the output current.

## Themry of Neutralization

Let the transistor $T$ in a given configuration be represented by

$$
\left[\begin{array}{l}
D_{1}  \tag{10}\\
D_{2}
\end{array}\right]=\left[\begin{array}{ll}
k_{11} & k_{12} \\
k_{21} & k_{22}
\end{array}\right] \times\left[\begin{array}{l}
I_{1} \\
I_{2}
\end{array}\right]
$$

One can then consider a network $N$ ( $N$ is not necessarily passive) describerl hy:

$$
\left[\begin{array}{c}
D_{1}{ }^{\prime}  \tag{11}\\
D_{2}{ }^{\prime}
\end{array}\right]=\left[\begin{array}{ll}
k_{11}{ }^{\prime} & k_{12}{ }^{\prime} \\
k_{21}{ }^{\prime} & k_{22}{ }^{\prime}
\end{array}\right] \times\left[\begin{array}{c}
J_{1}^{\prime} \\
J_{2}^{\prime}
\end{array}\right] .
$$

If $T$ and $N$ are connected to each other in a manner that forces $J_{1}{ }^{\prime}$ and $J_{2}{ }^{\prime}$ to be equal to $J_{1}$ and $J_{2}$ respectively, the composite network $C$ is defined by:

$$
\left[\begin{array}{l}
D_{1}^{\prime \prime}  \tag{12}\\
D_{2}^{\prime \prime}
\end{array}\right]=\left[\begin{array}{ll}
k_{11}+k_{11}^{\prime} & k_{12}+k_{12}{ }^{\prime} \\
k_{21}+k_{21}^{\prime} & k_{22}+k_{22}{ }^{\prime}
\end{array}\right] \times\left[\begin{array}{l}
J_{1} \\
J_{2}
\end{array}\right]
$$

If the feedback parameters of $T$ and $N$ satisfy the condition

$$
\begin{equation*}
k_{12}^{\prime}=-k_{12} \tag{13}
\end{equation*}
$$

the composite network $C$ is unilateral: one can say that the intermal feedthack of $T$ is nentralized.

These considerations can be applied to all representations (6) to (9) and result in the four basic neutralizing arrangements shown schematically in Fig. 3. The met hod of interconnection of $T$ and $N$ in these arrangements is consistent with the significance of the feedback parameters outlined at the end of the previous section.


Fig. 3-Masic neutralizing arrangements.

If a practical neutralizing circuit is designed on the hasis of the schematic arrangements of Fig. 3, caution must be exercised to interconnect $T$ and $N$ in a "permissible" manner ${ }^{5}$ insuring the validity of (12). Prob)lems of interconnection can be solved, as is well-known, using isolating transformers. Such a solution, however, is often undesirable because of circuit performance or cost. Consequently, the most practical circuits will be those which circumvent the difficulties of interconnection.

(a)

(b)

(c)

Fig. 4-h-type neutralization with and without transformer.

For example, one may consider the case of "h-type," ("series-parallel") neutralization. One transistor terminal is common to input and output, the transistor being in reality a three-terminal device. Neutralizing network $N$ may also be of the three-terminal type [Fig. 4(a)].

The two networks can be connected in series-parallel using an isolating transformer $X$ as shown in Fig. $4(\mathrm{~h})$. If the transformer is close to "ideal" with unity transformation ratio and negligible phase shift, the composite network is unilateral for $h_{12}=-h_{12}$; if the transformer provides 180 -(legree phase shift, neutralization

[^37]occurs for $h_{12}=h_{12}{ }^{\prime}$. However, the transformer is not necessary for interconnecting the networks. $T$ and $N$ can be directly connected in series-parallel, withont transformer, in eight different ways. Only one of these connections is "permissible" and is shown in Fig. 4(c). For this arrangement neutralization occurs provided $h_{12}^{\prime}=h_{12}$. This "bridge" circuit, is, of course, more desirable than its equivalent using a transformer.

Similar considerations apply to the "parallel-series" type interconnection of $T$ and $N$, for " $g$-type" neutralization. The use of a transformer can be avoided, as shown schematically in Fig. 5.


Fig. 5-g-type neutralization without transformer.
The composite networks of Fig. 4(c) and 5 are of the "four-terminal" type; no terminal is shared by input and output. Consequently, if such amplifying stages are cascaded, interstage transformers are usually required for reasons of common ground and power supply. If, on the other hand, a transformer is used within the neutralizing network, the composite network is of the threeterminal type and stages can be cascaded without an interstage transformer, if interstage impedance matching is considered minecessary. It is, however, often preferable to use an interstage transformer rather than a transformer inside the neutralizing loop. The reason for this is that the network $N$ required for neutralizing a transistor throughout a reasonable frequency range is usually simple, provided no transformer is used (or for the case of an "ideal transformer"). With an actual transformer in the neutralizing loop the design of $N$ must be modified to account for the deviation of the transformer from an "ideal" one and this may complicate the design of $N$ considerably. This problem does not arise if the transformer is external to the neutralizing loop and is used solely for interstage coupling.

(a)

(b)

(c)

Fig. 6-s-type neutralization.
'The problem of interconnecting $T$ and $N$ has somewhat different aspects in the case of $z$-type ("series") or $y$-type ("parallel") neutralization. The circuit of Fig. 6(a) is a "permissible" and transformerless series connection of $T$ and $N$ for $z$-type neutralization. The com-
posite network is neutralized provided $z_{12}{ }^{\prime}=-z_{12}$. For different transistor configurations the real part of $z_{12}$ is usually positive and consequently $z_{12}$ ' must have a negative real part of prescribed magnitude to match $z_{12}$. This can be realized easily for one given frequency, but it is generally difficult to obtain the desired $z_{12}{ }^{\prime}$ throughout a band of frequencies without using relatively complicated neutralizing networks consisting of four or more elements. The number of elements, however, should be kept minimum to avoid difficulties of alignment, and consequently, the transformerless circuit can be neutralized easily only in the neighborhood of one frequency. This may be satisfactory in certain applications, but if neutralization throughout a wider band is desired the circuit of Fig. 6 (b) [or its practical version of lig. 6(c)], using a transformer, would be preferable. The transformerless circuit may, however, be adequate for certain narrowband tuned amplifiers.


Fig. 7-y-type neutralization.

Similar considerations apply to $y$-type neutralization. ln Fig. 7 (a) no phase-inverting transformer is used and consequently neutralization will occur for $y_{12}{ }^{\prime}=-y_{12}$. In general, for transistors, $y_{12}$ has a negative real part [due to (7) and the sign of the currents and voltages in Fig. 1] and consequently $y_{12}{ }^{\prime}$ is required to have a prescribed magnitude with a positive real part. Without a transformer this can be realized only for one frequency using relatively simple networks. In most tuned amplifiers an interstage transformer is used anyway for interstage matching and the interstage transformer can simultaneously be used to provide neutralization in conjunction with a feedback network, as shown in Fig. 7(b). A practical realization is shown in Fig. 7 (c).

The discussion can be summarized by stating that while no transformer is necessary for $h$ - and $g$-type neutralization (except for interstage coupling purposes), in the case of $z$ - or $y$-type neutralization a transformer within feedback network may be desirable. For wide band neutralization, $h$ and $g$-methods are often superior.

## Discussion of the Feedback Parameters

## Equivalent Circuits and Approximations

The low frequency values of the transistor feedback parameters $h_{12}, g_{12}, y_{12}$ and $z_{12}$ can be easily determined using any of the well-known low frequency equivalent circuits of the transistor. 'This section is concerned mainly with the feedback parameters at "higher" frequencies.

The $T$-type equivalent circuit of the common emitter configuration [Fig. 8(a)] is derived from a simplified version of the well-known equivalent of the common base amplifier [Fig. 8(b)]. Using these representations the following approximations and assumptions are made, in order to simplify the expressions obtained:

1. The common base transistor is composed of a simplified "ideal transistor," derived from the diffusion equation, collector barrier capacitance "base spreading impedance" $Z_{b}{ }^{\prime}$ [Fig. 8(b)]. The "ideal transistor" is considered as having no internal feedback, feedback in the actual transistor being due to $Z_{b}{ }^{\prime}$. At higher frequencies, this is a justifiable simplification.


Fig. 8-Equivalents of common emitter and common base transistor circuits.
2. The frequency dependence of $a$ (short-circuit current amplification of the common base amplifier) is represented as

$$
\begin{equation*}
a=\frac{a_{0}}{1+j \omega / \omega_{a}}, \tag{14}
\end{equation*}
$$

where $\omega_{a}$ is the " $\alpha$-cut-off frequency" and $a_{0}$ the low frequency value of $a$.
3. The diffusion impedance of the emitter junction $Z$ 。 is a function of frequency

$$
\begin{equation*}
Z_{\theta}=\frac{r_{0}}{1+j \omega / \omega_{a}}, \tag{15}
\end{equation*}
$$

where $r_{\theta}$ is the low frequency value of $Z_{6}$. This means that $Z_{0}$ is considered as the parallel connection of $r_{0}$ and of a capacitance

$$
\begin{equation*}
C_{0}=\frac{1}{r_{0} \omega_{a}} . \tag{16}
\end{equation*}
$$

4. The collector impedance $Z_{c}$ is considered as the parallel connection of the collector resistance $r_{c}$ and the collector capacitance $C_{c}\left(C_{e}\right.$ being the sum of collector diffusion and barrier capacitances):

$$
\begin{equation*}
Z_{c}=\frac{r_{c}}{1+j \omega C_{c} r_{c}}=\frac{r_{c}}{1+j \omega / \omega_{c}} . \tag{17}
\end{equation*}
$$

At frequencies considerably higher than $\omega_{c} / 2 \pi$ :

$$
\begin{equation*}
Z_{c} \cong \frac{1}{j \omega C_{c}} \tag{18}
\end{equation*}
$$

5. The base spreading "resistance" $r_{b}$ ' is considered complex ( $Z_{b}{ }^{\prime}$ ) for grown junction $n-p-n$ transistors. This is due to the distributed nature of transistor parameters and $r_{b}{ }^{\prime}$, as shown by Pritchard and Coffey. ${ }^{8}$
The simplifying assumptions 1 to 4 are justified in the range of intermediate and higher frequencies and do represent the behavior of the transistor adequately up to a considerable fraction of $\omega_{a}$.

## Common Emitter Configuration

To explain the frequency dependence of the feedback parameters in common emitter configuration, it is necessary to consider the effect of a capacitance $C_{b c}$ connected between base and collector. The effect of this capacitance is particularly important in the case of grown junction $n-p-n$ transistors, as Pritchard has shown ${ }^{7}$ (base overlap capacitance). In the case of other transistor types $C_{b c}$ is smaller, but, of course, still exists, due to transistor lead, socket and wiring capacitances.
Using the equivalent circuit of Fig. 8(a), the $h$-parameters of the common emitter amplifier are found to be approximately:

$$
\begin{align*}
h_{11} & \cong Z_{b^{\prime}}+\frac{Z_{e}}{1-a}  \tag{19}\\
h_{12} & \cong \frac{Z}{(1-a) Z_{c}}+j \omega C_{b c}\left[Z_{b}^{\prime}+\frac{Z_{e}}{1-a}\right]  \tag{20}\\
& =h_{11^{\prime}}+h_{12}{ }^{\prime \prime} \\
h_{21} & \cong \frac{a}{1-a}  \tag{21}\\
h_{22} & \cong \frac{1}{(1-a) Z_{c}}+\frac{j \omega C_{b c}}{1-a} . \tag{22}
\end{align*}
$$

It must be remembered that, even if $\omega \gg \omega_{c}$, in view of the phase shift of $(1-a), h_{22}$ is not purely capacitive (it is, in fact, almost purely resistive throughout a wide range of frequencies).

The feedback parameter $h_{12}$ consists of two components $h_{12}{ }^{\prime}$ and $h_{12}{ }^{\prime \prime}$. According to (14) and (17)

$$
\begin{equation*}
h_{12}^{\prime}=\frac{Z_{c}}{(1-a) Z_{c}}=\frac{r_{e} /\left(1-a_{0}\right)}{r_{c}} \frac{1+j \omega / \omega_{c}}{1+j \omega /\left(1-a_{0}\right) \omega_{a}} . \tag{23}
\end{equation*}
$$

$h_{12}{ }^{\prime}$ has the constant value $r_{e} /\left(1-a_{0}\right) r_{c}$ at low frequencies, increases between $\omega_{c}$ and $\omega_{\alpha}\left(1-a_{0}\right)$ and at higher frequencies is equal to $r_{0} \omega_{a} C_{c}$. The phase shift of $h_{12}{ }^{\prime}$ vanishes at low and high frequencies and has a peak between $\omega_{c}$ and $\omega_{a}\left(1-a_{0}\right)$. The schematic frequency response of $h_{12}{ }^{\prime}$ is shown in Fig. 9 (opposite), whereas $h_{12}{ }^{\prime \prime}$ increases with frequency and has a positive phase shift.

[^38]The measured frequency dependence of $h_{12}$ of a grownjunction $n-p-n$ transistor is shown in Fig. 10. At low frequencies $h_{12}{ }^{\prime}$ prevails. At higher frequencies $h_{12}{ }^{\prime \prime}$ takes over and $h_{12}$ increases with frequency. The phase response corresponds to the amplitude response. The behavior of $h_{12}$ proves that the effect of $C_{b c}$ is considerable.


Fig. 9-Schematic frequency response of $h_{12}{ }^{\prime}$ (common emitter).


Fig. 10-Measured frequency response of $h_{12}$ for $n-p-n$ grown junction transistor (common enitter).


Fig. $11-h_{12}$ as function of $I_{s}$ and $V_{c}$ (common emitter).
$h_{12}$ is function of the de operating point. The dependence of $h_{12}$ on emitter current $I_{e}$ and collector voltage $V_{c}$ is shown in Fig. 11.

The other three feedback parameters can be calculated from (19) to (22). Current feedback parameter is:

$$
\begin{aligned}
&\left(-g_{12}\right)=\frac{h_{12}}{\Delta^{h}} \\
& \cong \frac{Z_{e}+j \omega C_{b c} Z_{c}(1-a)\left[Z_{b}^{\prime}+\frac{Z_{e}}{1-a}\right]}{Z_{b}^{\prime}+Z_{e}+j \omega C_{b c} Z_{c}(1-a)\left[Z_{b^{\prime}}+\frac{Z_{e}}{1-a}\right]}
\end{aligned}
$$

If $Z_{b}{ }^{\prime}$ is purely resistive and $C_{b c}$ is very small, the phase shift of $\left(-g_{12}\right)$ is negative. This, however, is not in agreement with measurements made on all rate grown $n p n$ and some $p n p$ transistors. In the case of grown junction $n p n$ transistors the discrepancy is clue mainly to the previously mentioned capacitive component of $Z_{b}{ }^{\prime}$ discussed by Pritchard and Coffey. They also measured the frequency response of $Z_{b}{ }^{\prime}$ and found that the magnitude and phase of $Z_{b}{ }^{\prime}$ change rather slowly with frequency. ${ }^{8}$ Fig. 12 shows the measured variation of $g_{12}$ with frequency for a rate grown $n-p-n$ transistor. Considering that in (24) usually $Z_{\ell} \ll Z_{b}{ }^{\prime}$, the response of $g_{12}$ seems to corroborate the measurements of Pritchard and Coffey. $g_{12}$ is function of $I_{e}$ and $V_{c}$ (Fig. 13).


Fig. 12-Measured frequency response of $g_{12}$ for an $n-p-n$ grown junction transistor (common emitter).


Fig. 13- $g_{12}$ as function of $I_{s}$ and $V_{c}$ (common emitter).
The feedback impedance parameter is:

$$
\begin{align*}
z_{12} & =\frac{h_{12}}{h_{22}} \cong Z_{e}+j \omega C_{b c} Z_{c}(1-a)\left[Z_{b}{ }^{\prime}+\frac{Z_{e}}{1-a}\right] \\
& =z_{12}{ }^{\prime}+z_{12}{ }^{\prime \prime} . \tag{25}
\end{align*}
$$

Here, as in the case of $h_{12}$, the effect of $C_{b c}$ is very important. If $C_{b c}$ could be neglected $z_{12}$ would be equal to $Z_{\text {e }}$ and would have a negative phase shift ( $Z$. being capacitive). The portion $z_{12}{ }^{\prime \prime}$ contributed by $C_{b c}$ is inductive and thus $z_{12}$ has a positive phase shift.

Feedback admittance parameter can be computed as:

$$
\begin{equation*}
\left(-y_{12}\right)=\frac{h_{12}}{h_{11}} \cong \frac{1}{Z_{c}} \frac{Z_{e}}{Z_{b}^{\prime}(1-a)+Z_{e}}+j \omega C_{b c} . \tag{26}
\end{equation*}
$$

${ }^{8}$ Pritchard and Coffey, loc. cit. (Fig. 8).

If $Z_{b}^{\prime}$ were zero, $\left(-y_{12}\right)$ would be purely capacitive at frequencies excecding $\omega_{c}$. In actual transistors, however, the feedback admittance has a considerable conductive component. The phase of $y_{12}$ depends, of course, also on the phase of $Z_{b}{ }^{\prime}$.

The expressions derived for the feedback parameters can be simplified considerably and their significance visualized if one considers frequencies exceeding $\left(1-a_{0}\right) \omega_{a} / 2 \pi$, but smaller than $\omega_{a} / 2 \pi$. In this frequency range one may use the approximations:

$$
\begin{align*}
& a \cong a_{0}  \tag{27a}\\
& Z_{c} \cong r_{a}  \tag{27b}\\
&(1-a)=\left(1-a_{0}\right) \frac{1+j \omega / \omega_{a}\left(1-a_{0}\right)}{1+j \omega / \omega_{a}} \cong j \frac{\omega}{\omega_{a}}  \tag{27c}\\
& 1 / Z_{c} \cong j \omega C_{c}  \tag{27d}\\
& Z_{b}^{\prime}>\frac{Z_{a}}{1-a} \tag{27e}
\end{align*}
$$

With these simplifications the feedback parameters can be written as:

$$
\begin{align*}
h_{12} & \cong r_{e} \omega_{a} C_{c}+j \omega C_{b c} Z_{b}^{\prime}  \tag{28}\\
\left(-g_{12}\right) & \cong \frac{r_{e}+j \frac{\omega}{\omega_{a}} \frac{C_{b c}}{C_{c}} Z_{b}^{\prime}}{Z_{b}^{\prime}+r_{c}+j \frac{\omega}{\omega_{n}} \frac{C_{b c}}{C_{c}} Z_{b}^{\prime}}  \tag{29}\\
z_{12} & \cong r_{e}+j \frac{\omega}{\omega_{a}} \frac{C_{b c}}{C_{c}} Z_{b^{\prime}}^{\prime}  \tag{30}\\
\left(-y_{12}\right) & \cong \frac{r_{a} \omega_{a} C_{c}^{\prime}}{Z_{b}^{\prime}}+j \omega C_{b c}^{\prime} \tag{31}
\end{align*}
$$

These approximations are very coarse [especially (27e) applies only to certain transistors having relatively large $Z_{b}^{\prime}$ and being biased in a restricted range of ac operating points] and do not permit accurate design, but give an impression of the order of magnitude of the neutralizing network components and the relative phases involved. The expressions can be easily correlated with the neutralized amplifier circuits discussed in the next section. If $Z_{b}{ }^{\prime}$ is resistive, the neutralizing network required can be determined with particular ease.

## Common Base Configuration

The effect of parasitic capacitances can be neglected in most cases when analyzing the common base configuration. The approximate expressions of the $h$-parameters are:

$$
\begin{align*}
& h_{11} \cong Z_{e}+Z_{b}^{\prime}(1-a)  \tag{32}\\
& h_{12} \cong j \omega C_{c} Z_{b}^{\prime}  \tag{33}\\
& h_{21} \cong a  \tag{34}\\
& h_{22} \cong \frac{1}{Z_{c}} \cong j \omega C_{c} \tag{35}
\end{align*}
$$

$\| Z_{b}{ }^{\prime}$ is purely resistive $\left(=r_{b}{ }^{\prime}\right)$

$$
\begin{equation*}
h_{12} \cong j \omega C_{c} r_{b}^{\prime} \tag{36}
\end{equation*}
$$

and has a positive phase shift of approximately 90 degrees.

The other three feedback parameters can be calculated from the $h$ parameters. The current feedback parameter is:

$$
\begin{equation*}
\left(-g_{12}\right)=\frac{h_{12}}{\Delta^{h}}=\frac{Z_{b}^{\prime}}{Z_{e}+Z_{b}^{\prime}} \tag{37}
\end{equation*}
$$

If $Z_{b}{ }^{\prime}=r_{b}{ }^{\prime}$ the phase shift of $\left(-g_{12}\right)$ is positive; if $Z_{b}{ }^{\prime}$ is capacitive, phase shift of $\left(-g_{12}\right)$ may become negative.

The feedback impedance parameter is:

$$
\begin{equation*}
z_{12}=\frac{h_{12}}{h_{22}}=Z_{b}^{\prime} \tag{38}
\end{equation*}
$$

The feedback admittance parameter is:

$$
\begin{equation*}
\left(-y_{12}\right)=\frac{h_{12}}{h_{11}}=\frac{j \omega C_{c} Z_{b^{\prime}}^{\prime}}{Z_{e}+Z_{b}^{\prime}(1-a)} \tag{39}
\end{equation*}
$$

At frequencies higher than $\left(1-a_{0}\right) \omega_{a} / 2 \pi$ but lower than $\omega_{a} / 2 \pi$, one can write in first approximation, using (27e):

$$
\begin{equation*}
\left(-y_{12}\right)=\omega_{a} C_{c} \tag{40}
\end{equation*}
$$

According to (40), in this frequency range, the feedback admittance is almost purely conductive, in reality, however, a reactive component is to be expected.

The above expressions can be easily correlated with the neutralizing circuits shown in the following section.

## Neutralized Ampiffier Circuits

During the experimental phase of the investigation mainly General Electric grown junction $n-p-n$ transistors (type 7.J6) were used. In many experiments, however, General Electric fused junction $p-n-p$ transistors (types $2 \mathrm{~N} 43,2 \mathrm{~N} 44,2 \mathrm{~N} 45$ ) were employed.

Due to the great variety of slightly different feasible neutralizing arrangements, a report describing all experiments completed would be impractical and repetitious in substance. Consequently, only a limited number of useful circuits will be discussed in this section. The experimentation was done to determine the practical value of the equations discussed in the preceding section and to answer the following questions:

1. Which types of neutralized circuits are most independent of dc collector voltage ( $V_{c}$ ) and emitter current ( $I_{e}$ ) variations? The transistor small-signal parameters being functions of the de operating point, dependence of the neutralized condition on $\Gamma_{c}$ and $I_{e}$ can, of course, be expected.
2. Which types of neutralized circuits are most appropriate for "wide-band neutralization" of transistors? 'The feedback parameters $h_{12}, g_{12}, z_{12}, y_{12}$ are not efually complicated functions of frequency. Consequently, some of them can be represented by simpler networks than the others.
3. How does neutralization affect the maximum available power gain of transistor amplifiers?

## Audio Frequency Circuits

In audio frequency applications the problem of stability does not arise. Neutralization may still be desired in certain applications to make input and output immittances independent of the terminal immittances.

Using the familier $T$-type equivalent circuit of the common base transistor [Fig. 14(a)], the " $h$-neutralized" common base amplifier circuit can be constructed by inspection [Fig. $14(\mathrm{~b})$ ]. In principle, the value of $k$ is arbitrary, but $k$ does, of course, influence the immittance levels and the maximum power gain of the stage. If $k$ is of the order of 0.1 reasonable unilateral performance can be expected: low input impedance, relatively high output impedance and, in most applications, practically unchanged power gain.


Fig. 14-(a) Low frequency equivalent of the common base amplifier, and (b) neutralized circuit.

U'sing the 7 -type equivalent circuit of the common collector stage [lig. 15(a)] a " $g$-neutralized" circuit can le derived [1Fig. 15(b)]. In this circuit

$$
\begin{aligned}
& R_{1}=k r_{c}(1-a) \\
& R_{2}=k\left(r_{b}+a r_{c}\right)
\end{aligned}
$$

The value of $k$ influences the immittance levels and the power gain of the neutralized amplifier, but is otherwise arbitrary. At very low frequencies this circuit exhibits very high input impedance and low output impedance. Including appropriate capacitances across $R_{1}$ and $R_{2}$ the frequency range of neutralization can be extended.

(a)

(b)
lijg. 15-(a) L.ow frequency equivalent of the common collector amplifier and (b) neutralized circuit.

Using a $p-n-p$ transistor with $R_{2}=440 \mathrm{~K}$ ohms and $R_{1}=24 \mathrm{~K}$ ohms an input impedance of 250 K ohms and output impedance of 30 ohms was measured at 500 cps , independently of source and load impedances. Maintaining the ratio $R_{1} / R_{2}$ constant and varying $R_{1}$, the
maximum available power gain of the amplifier varied as shown in Fig. 16. This circuit is useful as a high input impedance amplifier.

Other neutralized audio frequency circuits can be designed in a similar manner and also have attractive performance characteristics: their disadvantage lies in the necessary interstage transformer the cost of which may be prohibitive in many applications.


Fig. 16-Gain of a neutralized common collector stage vs $R_{1}$.

## High Frequency Circuits

A considerable number of neutralized amplifiers were built and tested in the frequency range between 100 kc and 3 mc . A few typical circuits are shown in Figs. 17 to 21. The investigation was mainly concerned with the common emitter amplifier, this configuration being the most important one from a practical point of view. The common emitter configuration exhibited a strong tendency toward becoming unstable with commonly used values of terminating impedances.


Fig. 17-h-neutralized common emitter circuits.
Fig. 17 (a) represents an $h$-neutralized common emitter amplifier. The structure of the neutralizing network corresponds to (28). With grown junction $n-p-n$ transistors the frequency performance of the circuit shown in lig. 17(b) was somewhat superior to that of Fig. 17(a): the amplifier could be neutralized throughout a wider frequency range, using a given neutralizing network. The reason for this hehavior is connected with the complex nature of $Z_{b}{ }^{\prime}$. The circuits of lig. 17 (b) and 17 (c) are equivalent. Using an inductive neutralizing network one can, by proper choice of the circuit elements, make the input impedance of the composite amplifier purely resistive at one frequency and the reactive component is reduced through a wide range of frequencies. Neutralization in these circuits varies moderately with $I_{e}$ and $V_{c}$.

Fig. 18 shows $g$-neutralized common emitter amplifier circuits. Figs. 18(a) and 18(b) are representative of most neutralized $n-p-n$ and $p-n-p$ stages, whereas Fig. 18(c) shows a $p-n-p$ stage with very small $C_{b c}$. The sign of the imaginary component of $g_{12}$ in 18(c) is different from that in 18(a) and 18(b). This phenomenon can be explained by considering the effect of the complex base spreading impedance of grown junction transistors and that of $C_{b c}$. The circuit of Fig. 18(b) was considerably superior to that of Fig. 18(a) from the point of view of dc operating point and frequency variations.


Fig. 18-g-neutralized common emitter circuits.
A $y$-neutralized common emitter stage is shown in Fig. 19. Due to the existence of $Z_{b^{\prime}},\left(-y_{12}\right)$ is not purely capacitive, as indicated by (31). Neutralization in this circuit was fairly independent of variations of the de operating point. Fig. 20 represents a $z$-neutralized common emitter amplifier. The structure of the circuit corresponds to (30).


Fig. 19-y-neutralized common emitter circuit.


Fig. 20-z-neutralized common emitter circuit.
Neutralization of the common base stage can also be achieved in many ways. The $h$-neutralized common base circuit of Fig. 21(a) has been described by Angell and Keiper. ${ }^{9}$ They have shown that the neutralizing network follows directly from the high frequency $T$-type equivalent circuit of the common base transistor. At lower frequencies perfect neutralization can be obtained by adding a resistance $r$ in parallel to the capacitor $C$ ( $r$ can be omitted if $\omega \gg \omega_{c}$ ). The neutraliza-

[^39]tion in this circuit is practically independent of the frequency, but varies with $I_{s}$ and $V_{c}$. A $z$-neutralized common base circuit is shown in Fig. 21(b).

If viewed closely, in all circuits described [with the exception of Fig. 21 (a)] the elements of the neutralizing network required for perfect neutralization varied with the operating frequency and consequently perfect neutralization was possible only at one given frequency. However, with most circuits approximate neutralization was achieved over a relatively wide band. For example, using the circuits of Figs. 17(c) and 18(b) and adjusting the circuit elements to give "perfect" neutralization at 500 kc , no noticeable dependence of the input and output immittances on the terminal immittances was experienced at 300 and 800 kc . Consequently from a practical viewpoint wide band neutralization can easily be achieved.


Fig. 21-Neutralized common base circuits.


Fig. 22-Circuit with "perfect" neutralization throughout very wide frequency range.
"Perfect" neutralization throughout a wide range of frequencies could be achieved with the aid of complicated feedback networks. For example, in view of the $\pi$-type equivalent circuit of the common emitter stage, ${ }^{10}$ the circuit of Fig. 22 could be adjusted for "perfect" neutralization throughout a very wide frequency band. However, the alignment of such a complicated circuit

[^40]would be exceedingly difficult. Consequently the simple circuits discussed above are definitely preferable, as they contain a maximum of three elements in the neutralizing network.

## Neutralization and Gain

The unneutralized transistor amplifier may become unstable provided that condition (5) is fulfilled. Consequently, the maximum unneutralized gain may become infinite and one must be cautious when referring to the maximum available power gain of an unneutralized transistor amplifier Pritchard and Coffey have shown ${ }^{11}$ that introducing certain constraints (e.g. matching at the output and tuning out the short circuit input impedance $h_{11}$ ) a constrained "maximum power gain" can be established for reference:

$$
\begin{equation*}
G_{u n n}=\frac{\left|h_{21}\right|^{2}}{h_{11}^{(R)} h_{22}(R)\left[1+\sqrt{1-I_{R} / h_{11}\left(R / h_{22}(R)\right.}\right]^{2}} \tag{41}
\end{equation*}
$$

The maximum available power gain of the neutralized amplifier is

$$
\begin{equation*}
G_{n}=\frac{\left|h_{21}^{\prime \prime}\right|^{2}}{4{h_{11}^{\prime \prime}(R)}_{h_{22}}{ }^{\prime \prime(R)}} \tag{42}
\end{equation*}
$$

Assuming now that the neutralizing network does not modify $h_{21}, h_{11}^{(R)}$ and $h_{22}{ }^{(R)}$ considerably, one can see that if $H_{R}>0$, neutralization will cause a decrease in gain, whereas if $H_{R}<0$ an increase in gain may result from neutralization.

In practical circuits, however, the neutralizing network modifies $\left|h_{21}\right|^{2 /} h_{11}^{(R)} h_{22}{ }^{(R)}$ considerably, and

$$
\frac{\left|h_{21}^{\prime \prime}\right|^{2}}{{h_{11}^{\prime \prime}(R)}^{\prime \prime} h_{22}^{\prime \prime(R)}}<\frac{\left|h_{21}\right|^{2}}{h_{11}^{(R)} h_{22}^{(R)}}
$$

Consequently, even if $H_{R}<0$, neutralization usually results in moderate, if any, additional gain. Gain decreases of few db , due to neutralization, have been measured using the common emitter configuration. Small decrease or increase in gain has been measured in the case of the common base circuit.

[^41]
## Measurement of Feedback Parameters

Neutralization, i.e., bridge methods, can be used advantageously to measure certain feedback parameters of transistors. A simplified version of a measuring arrangement used throughout this study for measuring $h_{12}$ is shown schematically in Fig. 23.


Fig. 23-Arrangement for measuring $h_{12}$.
The signal generator feeds the output 2 of the transistor and the signal appearing at the input 1 is amplified and displayed on an oscilloscope. If the neutralizing network is properly adjusted, the backward transmission becomes zero and $h_{12}$ can be computed from the values of the neutralizing network elements. Similar methods can be used for the measurement of $g_{12}$.

## Conclusion

Internal feedback of transistor amplifiers can be cancelled using different methods of neutralization. Neutralized transistor circuits are stable with input and output immittances independent of terminal immittances.

The elements of the circuits required for neutralization depend on the dc operating point and the frequency. Neutralization throughout relatively wide bands of frequencies can be achieved using simple feedback arrangements. The gain of the amplifier may be increased or decreased as a result of neutralization, depending on the nature of the internal feedback of the unneutralized amplifier.

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# Backward-Wave Oscillator Efficiency* 

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#### Abstract

Summary-The theoretical and experimental results of a study of the factors which determine the efficiency of backward-wave oscillators are described. The dependence of power output upon space charge, circuit loss, beam thickness, velocity spread, and circuit mismatches is determined by a combination of theoretical and empirical means. In addition, the effect of circuit mismatches upon the starting current and frequency is discussed. The study shows that efficiency can be increased by increasing either the gain parameter $C$ or the space-charge parameter $\omega_{q} / \omega$. Circuit loss, beam thickness, and velocity spread are found to decrease the efficiency. The use of the results reported here permits the designer to control the factors affecting the power output and predict the efficiency of a tube with reasonable accuracy before it is built.


## Introduction

1IIE I)ISCOVERY of backwarl-wave oscillations ${ }^{1}$ in traveling-wave devices has led to a new type of microwave tube the backward-wave oscillator, or "carcinotron" as it is called by French workers." Backward-wave oscillations occur when an electron beam interacts with a periodic structure which is in general shorter than that used for a traveling-wave amplifier. This fact explains the occasional occurrence of backward-wave oscillations in conventional travelingwave amplifiers. When the electron velocity is synchronized with the phase velocity of a backward wave of a slow-wave structure, the device behaves as a back-ward-wave amplifier with internal positive feedback and will amplify for values of current below a critical value. dhove the critical value of beam current the device will oscillate. Since backward waves are dispersive with a phase velocity which is a function of frequency, the frequency of oscillation may be changed continuously by changing the electron velocity. Thus the voltage of the electron beam may be used to tune the frequency of oscillation of a backward-wave oscillator. This brief description shows that backward-wave oscillators are both new and interesting microwave devices. It shows too that such oscillators have characteristics which are not common to other tubes.

The purpose of this paper is to add to the existing knowledge an understanding of some of the factors which determine the level of oscillation and hence the power output or efficiency to be expected. Other workers have described their experimental and theoretical re-

[^42]sults. ${ }^{3-12}$ In general, the definitions and notation introduced by P'ierce ${ }^{3}$ will be used and it should be noted that all of the analysis contained in the paper concerns the extension of linear theory ${ }^{6}$ to the region of oscillation. No attempt has been made to carry out a nonlinear theor $y^{13}$ of backward-wave oscillators because of the complexity of such a study. The objective has been to obtain theoretical and experimental information that could be applied easily to predict the efficiency of oscillators with reasonable accuracy. The factors which are shown to influence the efficiency may be listed as follows: space charge, circuit loss, beam thickness, and velocity spread. A study of the effect of imperfect matches on the starting. conditions and on the efficiency is also included. The results contained in this paper make it possible to design backward-wave oscillators with a fairly clear idea of the factors that influence the efficiency so that the amount of rf output power to be obtained may be controlled in the design. We will assume that the efficiency of a back-ward-wave oscillator can be written in the following form
\[

$$
\begin{equation*}
\eta=\eta_{0}(Q C) F_{1}\left(L_{d b}\right) F_{2}\binom{\beta t}{\beta b} F_{3}(S) F_{4}^{\prime}(R) \tag{1}
\end{equation*}
$$

\]

where each of the factors on the right is defined and discussed in the following sections. The form of (1) may not be valid if the correction factors $F_{1}-F_{4}$ differ greatly from unity.

The symbols used above and throughout this paper are as follows:

$$
\begin{aligned}
b & =\text { velocity parameter defined by I'ierce, } \\
C & =\text { gain parameter defined by Pierce, } \\
E(z) & =r f \text { circuit electric field at the plane } z,
\end{aligned}
$$

[^43]$F_{1}\left(L_{d b}\right)=$ efficiency reduction factor due to the circuit loss,
$F_{2}\binom{\beta t}{\beta b}=$ efficiency reduction factor due to beam
$F_{3}(S)=$ efficiency reduction factor due to velocity spread of the beam,
$F_{4}(R)=$ efficiency reduction factor due to reflections,
$G=$ efficiency parameter defined l)y (4),
$I_{0}=$ average convection current of the beam,
$i(z)=$ rf convection current of the beam,
$L_{d b}=$ total uniform circuit loss in decibels,
$L=$ length of circuit,
$P(z)=$ power on the circuit at the plane $z$,
$P_{a v}=$ power output obtained from the average value of $C$,
$Q C=$ space-charge parameter defined by Pierce,
$R=$ reflection coefficient defined by (33),
$S=$ electron beam velocity distribution parameter,
$V_{0}=\mathrm{dc}$ beam voltage,
$V(z)=\mathrm{rf}$ circuit voltage at the plane $z$,
$y=2 \pi C N=$ length parameter defined by Pierce,
$z=$ distance parameter of one dimensional model,
$\alpha=$ attenuation constant of the circuit,
$\beta=$ cold circuit propagation constant,
$\beta_{s}=$ propagation constant of the electron beam,
$\xi=$ incremental propagation constant,
$\Delta=$ incremental propagation constant defined by (9)
$\rho=$ reflection coefficients,
$\omega=$ frequency of oscillation in radians per second,
$\omega_{q}=$ reduced plasma frequency in radians per second,
$\eta=$ over-all efficiency,
$\eta_{0}(Q C)=$ basic efficiency as a function of space charge.

## General Method of Solution

The calculation of efficiency of a one-dimensional, lossless backward-wave oscillator is based on the assumption that the oscillation level of the tube is limited by the saturation of the beam current. The degree of saturation is assumed by specifying the relationship between the magnitude of the rf portion of the convection current and the de portion of the convection current at the collector end of the interacting structure. At least two logical values may be assumed. One corresponds to the assumption that the beam has no harmonics and has the appearance of a sinusoid plus a constant value where the components have magniturles at the collector end of the structure given by

$$
i=I_{0}
$$

which is a choice which will be used to a large extent in
the following. Another value of special interest which will be used to a smaller extent is based on the maximum theoretical value which the fundamental rf component can have with reference to the de component. This maximum value occurs when the waveform consists of a series of equally spaced delta functions and has equal-amplitude harmonics related to the steady component by

$$
i=2 I_{0}
$$

In general the value $i=I_{0}$ will be used for reasons which will become evident later.

The definition of $C$ will he used to determine the value of the rf power on the circuit of the backward-wave oscillator. The efficiency may be written thus:

$$
\begin{equation*}
\eta_{0}=\frac{P(0)}{I_{0} V_{0}}=\frac{|E(0)|^{2}}{V_{0}^{2}} \frac{1}{8 \beta^{2} C^{3}} \tag{2}
\end{equation*}
$$

The efficiency is determined from (2) by relating $E(0)$ and $i(L)$ which is easily accomplished by relating the respective waves of the convection current and the circuit electric field from the linear theory. Thus, the convection current may be written in the following manner:

$$
\begin{equation*}
i(L)=-j \frac{I_{0} G}{2 V_{0} C^{2} \beta_{e}} E(0) \tag{3}
\end{equation*}
$$

where $G$ has the definition given below

$$
\begin{align*}
G= & {\left[\frac{e^{i \nu \xi_{1}}}{\left(\xi_{1}-\xi_{2}\right)\left(\xi_{1}-\xi_{3}\right)}+\frac{e^{i \nu \xi_{2}}}{\left(\xi_{2}-\xi_{1}\right)\left(\xi_{2}-\xi_{3}\right)}\right.} \\
& \left.+\frac{e^{j y \xi_{3}}}{\left(\xi_{3}-\xi_{1}\right)\left(\xi_{3}-\xi_{2}\right)}\right] e^{-i \beta_{e} L_{0}} . \tag{4}
\end{align*}
$$

Then letting the current at the collector be equal to the dc current $I_{0}$ and solving for $E(0)$ from (3) we obtain

$$
\begin{equation*}
E(0)=j \frac{2 V_{0} C^{2} \beta_{e}}{G} \tag{5}
\end{equation*}
$$

Eq. (5) may be substituted into the efficiency relation given by (2) with the following result:

$$
\begin{equation*}
\eta_{0}=\frac{C}{2|G|^{2}} \tag{6}
\end{equation*}
$$

Eq. (6) represents the general expression for efficiency. The efficiency may be found by substituting in to (6) the value of $G$ as defined by (4) and as determined from the start oscillation condition. We should note that the efficiency relation will hold for various conditions providing that the correct values of the roots for any given degree of space charge are substituted into (4). As an example for the space-charge case the $\xi$ 's are the three roots of the equation

$$
\begin{equation*}
\xi^{3}+b \xi^{2}-4 Q C \xi-4 Q C b+1=0 \tag{7}
\end{equation*}
$$

Thus it should le observed that the $\xi$ 's used here differ from Pierce's $\delta$ 's by the quantity $j$. 'The general equations (4) and (6) will be used throughout this paper.

## Efficiency for Large Space Charge

In order to solve (6) for large space charge and small $C$ it is necessary to determine the behavior of the roots and to evaluate both $2 \pi C N$ and $|G|^{2}$. One important equation for large space charge which will be demonstrated in this section for start oscillation is given as

$$
\begin{equation*}
b \cong \sqrt{4 \varrho C} \tag{8}
\end{equation*}
$$

The roots of (7) for this value of $b$ all have approximately the same magnitude for large space charge so that a simplification results from the substitution

$$
\begin{equation*}
\Delta=b+\xi, \tag{9}
\end{equation*}
$$

which will separate the two equivalent roots. Substituting (9) into (7) yields

$$
\begin{equation*}
\Delta^{3}-2 b \Delta^{2}+\left(b^{2}-4 \varrho C\right) \Delta+1=0 . \tag{10}
\end{equation*}
$$

If the real root of (10) is called $\Delta_{1}$, then the other two roots may be written approximately as follows:

$$
\begin{align*}
& \Delta_{2} \cong-\Delta_{1} \\
& \Delta_{3} \cong 2 b . \tag{11}
\end{align*}
$$

Substitution of these approximate roots into the equation for start oscillation leads to the conclusion that

$$
\begin{equation*}
y \Delta_{1} \cong \frac{\pi}{2} \tag{12}
\end{equation*}
$$

which result depends on the assumption that $\Delta_{1}$ is much smaller than $b$. This assumption will be justified later. However, substituting $\Delta_{1}$ into (10) and applying the same assumption along with (8) yields the result

$$
\begin{equation*}
1=2 b \Delta_{1}{ }^{2} . \tag{13}
\end{equation*}
$$

Combining (8), (12), and (13) leads to the conclusion

$$
\begin{equation*}
(2 \pi C N)^{2}=y^{2}=\frac{\pi^{2}}{2} \sqrt{4 \varrho C} \tag{14}
\end{equation*}
$$

The validity of (8) may now be verified by constructing (10) from the roots of (11) and noting that

$$
\begin{equation*}
b^{2}-4 \varrho C=\Delta_{1}^{2} . \tag{15}
\end{equation*}
$$

Eqs. (12), (13), and (14) are consistent only where $\Delta_{1}$ decreases with an increase of $b$. Thus the basic assumption given by ( 8 ) is consistent with the oscillation condition. By means of (8), (11), and (14) the evaluation of $|G|$ for large space charge may be carried out and the corresponding efficiency becomes

$$
\begin{equation*}
\eta_{0}=\sqrt{4 \varrho C^{3}}=\frac{\omega_{Q}}{\omega}, \tag{16}
\end{equation*}
$$

where $\omega_{q}$ is the reduced plasma frequency and $\omega$ is the frequency of oscillation. We should realize that (16) is based on the conditions which exist at start oscillation where the start oscillation conditions are satisfied. At this point $b$ has the value given by (8). It should be pointed out that although the same result is obtained
with a similar saturation assumption for a travelingwave tube, ${ }^{14}$ a backward-wave oscillator with large space charge is found experimentally to adhere much closer to this theoretical value.

It is interesting to note that the waveforms which exist on a backward-wave oscillator for large space charge may be written in analytic form. The current waveform is given by

$$
\begin{equation*}
i(z)=i(L) \sin \frac{\pi z}{2 L} e^{-j \theta z}, \tag{17}
\end{equation*}
$$

whereas the circuit field waveform is given by

$$
\begin{equation*}
E(z)=E(0) \cos \frac{\pi z}{2 L} e^{-j \beta \varepsilon} . \tag{18}
\end{equation*}
$$

Examination of these functions shows that the beam is bunched to the maximum extent when the field is zero and conversely the convection current is zero when the field is maximum. This situation is entirely different from that which exists in a traveling-wave tube where the convection current and the circuit field both increase exponentially at the same rate. For this reason the efficiency of a backward-wave oscillator may be expected to be less than a comparable traveling-wave tube operating at the point of maximum gain.

## Theoretical Efficiency for Smail. Space Charge

From (4), (6), (7), and the oscillation condition, it is possible to solve for the efficiency for any value of space charge. For example, at zero space charge where

$$
\begin{aligned}
b & =1.52 \\
y & =2 \pi C N=1.97
\end{aligned}
$$

we obtain the efficiency

$$
\begin{equation*}
\eta_{0}=0.21 C . \tag{19}
\end{equation*}
$$

Fig. 1 was obtained by plotting this point and others where the assumption of equivalence between the rf convection current and the dc beam current is made. It is consistent with Pierce ${ }^{14}$ to suppose that at zero space charge the efficiency will more nearly be given by multiplying the result given by (19) by four to obtain

$$
\begin{equation*}
\eta_{0}=0.84 C, \tag{20}
\end{equation*}
$$

which is equivalent to the assumption that $i=2 I_{0}$. The ripples shown in Fig. 1 occur because of ripples which exist in the magnitude of the convection current as functions of distance and space charge. The process of setting the rf convection current equal to the dc beam current causes ripples to exist in the efficiency curve. For $\omega_{q} / \omega C>1$ the ripples are seen to vary about the value given by (16). It is important to note that the presence of the ripples is a result of the linear theory and could hardly be expected to occur in actual tubes.

[^44]

Fig. 1-Theoretical curve of $\eta / C$ vs $\omega_{q} / \omega C$ obtained from linear theory with the assumption that the electron beam at the collector saturates at a value of $i=I_{0}$.

## Probable Efficiency and <br> Exprermental Results

Summarizing the results of the previous sections we find that for zero space charge

$$
\eta_{0} \simeq C
$$

and for the large space charge

$$
\eta_{0} \simeq \frac{\omega_{q}}{\omega},
$$

In the region between zero space charge and large space charge we should expect the efficiency to depend on both $C$ and $\omega_{q} / \omega$.

Experimental verification of the large space charge relation was obtained with a tube designed by J. I., Putz and W. R. Luebke of this laboratory. Fig. 2 shows plots of efficiency vs frequency for beam currents of 250 and 150 ma . The theoretical values of $\omega_{q} / \omega$ are also shown in Fig. 2. The close agreement which was achieved between the theoretical curves and the experimental curves was gratifying.

Points which were obtained from a number of tubes at various frequencies are shown in Fig. 3. The cluster of points leads to the combined theoretical and empirical curve shown also in Fig. 3. This curve will be subsequently called the "lasic" efficiency curve and the procedure used to predict efficiency will be based upon this curve and upon the correction factors subsequently described. Experimentel results indicate that the "basic" efficiency curve may be used when the beam current is more than two times the starting current.


Fig. 2-Comparison of theoretical and measured efficiency of Putz and Luebke's tube operating at 150 and 250 ma.


Fig. 3-Comparison of the experimental results obtained for the efficiency of a large number of backward-wave oscillators and theoretical results presented in this paper. A probable efficiency curve is drawn for the best agreement between the theory and experiment and is called the "basic" efficiency curve.

## Efficiency Rebuction Due to Circuit Loss

On a lossy structure it is reasonable to suppose that the output of a backward-wave oscillator is reduced hecause the energy is not transferred to the circuit at the output terminal but must travel through at least a portion of the lossy structure. The reduction in effaciency of an oscillator with uniformly distributed circuit loss is determined in this section by reducing each increment of power according to the amount of loss which it encounters. The magnitude of the field on a structure with large space charge is given by (18) which is not greatly different from the waveform for small space charge. From (18) power on the structure can be expressed as

$$
\begin{equation*}
P(z)=P_{\text {out }}(0) \cos ^{2} \frac{\pi z}{2 L} . \tag{21}
\end{equation*}
$$

The rate of change of the power with $z$ is found by


Fig. 4-Theoretical efficiency reduction factor for the case of uniform circuit loss. The experimental points were furnished by W. A. Harman.
differentiating (21) with respect to $z$ :

$$
\begin{equation*}
\frac{d P(z)}{d z}=\frac{\pi}{2 L} P_{\text {out }}(0) \sin \frac{\pi z}{I .} \tag{22}
\end{equation*}
$$

The ratio of transfer of power from the beam to the line is seen to be a sinusoidal function with a zero occurring at both ends of the structure and most of the transfer taking place at the midpoint. This condition may actually be suspected from the nature of the field and current waveforms. A rough approximation may be obtained by considering that all the power orginates in the center of the backward-wave oscillator and traverses one half the length of the structure in reaching the output terminal. This approximation yiclds the result

$$
\begin{equation*}
F_{1}=\exp \left(-0.115 L_{d b}\right) . \tag{2.3}
\end{equation*}
$$

We should expect (23) to approximate the true result. To analyze the problem more exactly we can attenuate all elements of power according to their origin on the line. This procedure yields the result

$$
\begin{equation*}
F_{1}=\exp \left(-0.115 L_{d b}\right) \frac{\cosh \left(0.115 L_{d b}\right)}{1+\left(\frac{0.230 L_{d b}}{\pi}\right)^{2}} \tag{24}
\end{equation*}
$$

which is the exact solution. It is apparent that the approximate result appears with a modifying correction to take account of the fact that all the elements of power do not originate at the midpoint of the structure. The resulting function of $F_{1}$ versus loss $L_{d b}$ has been plotted in Fig. 4. Also shown are experimental points provided by W. A. Harman of this laboratory. The experimental points were positioned to obtain the best fit to the theoretical curve inasmuch as the loss could not be completely removed from the oscillator so that the no-loss point could not be determined. The good agreement between the measured points and the theoretical curve is indicative of the fact that the effect of distributed loss on efficiency can be predicted by (24) or Fig. 4.

## Efficiency Renuction for a Thick Beam

In a backward-wave oscillator with a thick beam such that the rf electric field varies across the beam cross sec-


Fig. 5-Theoretical efficiency reduction factor for a thin hollow beam where fields vary as $I_{1}(\beta r)$ or $I_{0}(\beta r)$ and $\beta b \gg \beta t$.
tion we might expect the magnitucle of the saturation convection current of an element of the beam to be proportional to the electric field which acts upon this beam element. It is consistent with the assumption of current saturation to assume that the saturation of the element of the beam which lies in the region of highest impedance will control the degree of saturation of the rest of the beam. 'These two assumptions have been used in the analysis which follows. The application of these assumptions leads to the relationship given below in which each element $n$ of the beam is assumed to have an effeciency $\eta_{n}$ which may le written as follows

$$
\begin{equation*}
\eta_{n}=\eta_{0}(Q C) \frac{K_{n}}{K_{\max }}, \tag{2.5}
\end{equation*}
$$

where the uncorrected efficiency $\eta_{0}(Q C)$, as determined from the space charge parameter $\omega_{q} / \omega$ and the average value of $C$, is the same for all elements $n$ of the beam.

It is convenient for a thick heam to use Fig. 3 to determine a value for the eff ciency based on the average value of $C$ and then correct this value by an amount depending on the beam thickness. This procedure makes it possit)le to use the same value of impedance to compute both the starting conditions and the uncorrected efficiency. 'The correction factor $F_{2}$ is deffned in the following manner:

$$
\begin{equation*}
F_{2}=\frac{P}{P_{a \vartheta}}=\frac{\sum_{n} \eta_{n} I_{n}}{\eta_{0}(Q C) I_{T}} . \tag{26}
\end{equation*}
$$

The combination of (25) and (26) leads to the general conclusion that

$$
\begin{equation*}
F_{2}=\frac{K_{a v}}{K_{\max }} \tag{27}
\end{equation*}
$$

Eq. (27) may be solved for both a hollow beam and a solid beam. The efficiency correction factor for a hollow beam with thickness $t$ and $\beta b \gg \beta t$ may be written as

$$
\begin{equation*}
F_{2}(\beta t)=\frac{1-e^{-2 \beta t}}{2 \beta t} \tag{28}
\end{equation*}
$$

This expression has been plotted in Fig. 5. It is equally valid for a field variation of $I_{0}(\beta r)$ or $I_{1}(\beta r)$ as long as
$\beta b \gg 1$. A similar procedure for a solid beam with ' radius $b$ and with the electric field varying as $I_{0}(\beta r)$ yields the relation

$$
\begin{equation*}
F_{2}(\beta b)=1-\frac{I_{1}^{2}(\beta b)}{I_{0}^{2}(\beta b)} \tag{29}
\end{equation*}
$$

which has been plotted in Fig. 6. These reduction factors can be used to compute the efficiency when the beam is thick and where the $C$ has been computed in the usual small-signal manner. Fig. 5 for a thick hollow beam has been successfully used to compute the output power of a hollow-beam backward-wave oscillator built by L. A. Roberts of this laboratory. A comparison of the theoretical and measured power output for this oscillator is shown in Fig. 7. The agreement between the curves is well within the known accuracy of the tule parameters. It should, however, be noted that the theoretical curve of Fig. 7 includes the correction for the uniformly distributed circuit loss as given by Fig. 4.


Fig. 6-Theoretical efficiency reduction factor for solid beam where the field varies as $I_{0}(\beta r)$.


Fig. 7-Comparison of measured and theoretical power output where correction was made for the thickness of a hollow beam. The experimental data were furnished by L. A. Roberts.

## Effect of the Velocity Spread of the Beam

In an electron beam focused by means of an axial magnetic field uniform through the cathode, the dc velocity of the electrons varies over the beam cross sec-
tion. This variation is caused by the potential depression produced by the charge of the electrons. Thus some elements of the beam travel faster than other elements and the synchronous beam voltage is not clearly defined. Under these conditions Watkins and Rymn ${ }^{15}$ have shown that velocity spread in traveling-wave devices has an effect similar to space charge in the region of operation where linear theory applies. It was suspected that linear theory could not be extended in this case to the region of nonlinearity. No reasonable theory has been developed to predict the effect of velocity spread of the electrons of the beam on the efficiency of backwardwave oscillators. This section concerns an experimental approach to the problen. An actual backward-wave oscillator was modified to make possible a measurement of the effect of velocity spread on efficiency:


Fig. 8-I'hotograph of the $500-1,000 \mathrm{mc}$ backward-wave oscillator used to measure the effect of velocity spread on the efficiency. The oscillator used a hollow electron beam placed close to the helix.

This measurement was performed on the backwardwave oscillat or shown in Fig. 8. The beam was 0.005 inch thick and was spaced about 0.005 inch from the helix which was ahout 1.3 inches in diameter. This normalarrangement was modified by inserting a 0.5 inch diameter cylinder down the axis of the tube, as suggested by P. D. Lacy ${ }^{16}$ of the Hewlett-Packard Co. The cylinder could be operated at any desired voltage with respect to the helix in order to introduce artificially a velocity spread to the electrons of the beam. In order to obtain meaningful data all of the parameters of the tube were held constant except the helix voltage and the cylinder voltage. The frequency was maintained constant by adjusting both the helix and cylinder voltages to keep the "average" beam voltage constant. Measurements were made at five different frequencies across the band. The measured data are shown in Fig. 9 (next page) where the reduction in power output expressed by $F_{3}(S)$ is plotted against the parameter

$$
S=\left(\frac{\Delta V}{4 C V_{0}}\right)^{2}, \quad C \cong 0.06 \text { (for these measurements) }
$$

which was introduced by Watkins and Rymn. ${ }^{15}$ Since the

[^45]gain parameter $C$ did not vary for these measurements, it could not easily be determined whether the parameter $S$ is the correct one to use for the reduction of efficiency.

The examination of the experimental points led to their presentation on semi-log paper where they are seen to trace a straight line. The magnitude of the reduction down to 10 per cent, although seemingly large, corresponds to a relatively large voltage variation. Thus, at $S=0.71$ the voltage drop across the beam was about 8 v . This is quite large when compared to the average value of 36 v . The straight line, drawn for the best fit, represents an empirical curve showing efficiency reduction.


Fig. 9-The experimental measurement of efficiency reduction caused by velocity spread in the $500-1,000 \mathrm{mc}$ backward-wave oscillator. The empirical curve was drawn for the hest fit to the experimental data.

## Effect of Reflections on Starting Conditions

The gain expression for small $C$, zero space charge, and no loss is given by the relation

$$
\begin{align*}
\frac{V(L)}{V(())}= & \frac{e^{i \nu \xi_{1}}}{\left(\xi_{1}-\xi_{2}\right)\left(\xi_{1}-\xi_{3}\right)\left(\xi_{1}+b\right)} \\
& +\frac{e^{i y \xi_{2}}}{\left(\xi_{2}-\xi_{1}\right)\left(\xi_{2}-\xi_{3}\right)\left(\xi_{2}+b\right)} \\
& +\frac{e^{i y \xi_{3}}}{\left(\xi_{3}-\xi_{1}\right)\left(\xi_{3}-\xi_{2}\right)\left(\xi_{3}+b\right)} . \tag{.30}
\end{align*}
$$

The method used to solve for the start oscillation condition is to set $V^{\prime}(L)=0$ and solve (30) simultaneously with the root equation

$$
\begin{equation*}
\xi^{3}+b \xi^{2}+1=0 \tag{31}
\end{equation*}
$$

A different approach was used to solve for the effect of reflections. The equations were solved as a function of $V(I) / V(0)$ to obtain the starting parameters when $V(L) / V(0)$ is not zero. The quantity $V(L) / V^{\gamma}(0)$ is the ratio bet ween the voltage that is applied at the collector
end of the transmission line and the output voltage. For convenience let us define the ratio as

$$
\begin{equation*}
R e^{\prime \theta}=\frac{V(L)}{V(0)} \tag{32}
\end{equation*}
$$

The quantity $R$ is a real number which depends on the matches at the ends of the tube and the loss of the circuit as defined by the relation

$$
\begin{equation*}
R=\rho_{1} \rho_{2} e^{-\alpha L} \tag{3.3}
\end{equation*}
$$

where $\rho_{1}$ and $\rho_{2}$ are the reflection coefficients and $\alpha$ is the attenuation constant. $R$ will be zero when the reflection coefficients are zero or the loss is infinite. The method of solution was to expand about the point $R=0$ for deviations of $y$ and $b$ in the form of a Taylor series:

$$
\begin{align*}
R e^{i \theta}= & f(y, b) \\
= & f(1.97,1.52)+f_{y}(1.97,1.52) \Delta y \\
& +f_{b}(1.97,1.52) \Delta b \tag{34}
\end{align*}
$$

The first term of the series is identically zero, and only the first-order effects have been inclurded to make the problem soluble. For convenience let us define

$$
\begin{align*}
& f_{y}(1.97,1.52)=A e^{i \alpha} \\
& f_{b}(1.97,1.52)=B e^{i \beta} . \tag{.35}
\end{align*}
$$

Then separating (34) into a real and an imaginary equation and solving simultaneously yields the result

$$
\begin{align*}
& \Delta y=\frac{R \sin (\beta-\theta)}{A \sin (\beta-\alpha)}  \tag{36}\\
& \Delta b=\frac{R \sin (\alpha-\theta)}{B \sin (\alpha-\beta)} . \tag{.37}
\end{align*}
$$

It has been demonstrated here that if we find the value of $f_{y}(1.97,1.52)$ and $f_{b}(1.97,1.52)$ then we can write $\Delta y$ and $\Delta b$ in analytical form to give the variation for values of $R>0$. From (36) and (3i) we see that $\Delta y$ and $\Delta b$ are sinusoidal functions of $\theta$ with amplitudes which are proportional to the value of $R$. Thus we see that the effect of reflections, at least on the starting conditions, should be sinusoidal in mature and should vary about the mean position which would exist in the alsence of the reflections. The evaluation of $f_{\nu}(y . b)$ at the point $y=1.97$ and $b=1.52$ where the $\xi^{\prime}$ 's are given by' (31) yields the result

$$
\begin{align*}
A & =1.53 \\
\alpha & =223 \text { degrees } \tag{.38}
\end{align*}
$$

and the evaluation of $f_{b}(y, b)$ at the point $y=1.97$, $b=1.52$ yields the result

$$
\begin{align*}
B & =1.045 \\
\beta & =296.7 \text { degrees. } \tag{39}
\end{align*}
$$

The results given by (38) and (39) may be sulbstituted into (36) and (37) to give the absolute magnitude of $\Delta y$ and $\Delta b$

$$
\begin{align*}
|\Delta y| & =1.68 R  \tag{40}\\
|\Delta b| & =-1.00 R \tag{41}
\end{align*}
$$

These values represent the mathematical solution of the problem. From (40) we can obtain

$$
\begin{equation*}
\frac{\left|\Delta I_{0}\right|}{I_{0}}=R \tag{42}
\end{equation*}
$$

which is the desired result expressing the change of starting current as a function of $R$. From (36) and (42) we can see that $\Delta I_{0} / I_{0}$ is a sinusoidal function with an amplitude $R$ which shows that the starting current $I_{0}+\Delta I_{0}$ varies in a simusoidal manner about a mean position $I_{0}$ corresponding to the starting current in the absence of reflections.

A similar expression for the frequency variations may be obtained from the definition of $b$ as given below:

$$
\begin{equation*}
\frac{|\Delta f|}{f}=2 C R \frac{V_{0}}{f}\left|\frac{d f}{d V_{0}}\right| \tag{43}
\end{equation*}
$$

where $C$ is the gain parameter, $V_{0}$ is the beam voltage, and $d f / d V_{0}$ is the tuning rate. On a helical structure (43) may be written as follows:

$$
\begin{equation*}
\frac{|\Delta j|}{f}=C R(1-k a) \tag{44}
\end{equation*}
$$

It has thus been possible to find analytic relationships

$$
\begin{equation*}
I_{4}=\frac{1}{1+\frac{2 R}{\sin (\alpha-\beta)}\left[\frac{B_{3}{ }^{2}}{B^{2}}+\frac{A_{3}{ }^{2}}{A^{2}}+2 \frac{A_{3} B_{3}}{A B} \cos (\alpha-\beta)\right] \cos \left(\theta+\theta_{0}\right)} \tag{51}
\end{equation*}
$$

for the magnitude of the variations of starting current and frequency as a function of the reflection parameter $R$. Both (42) and (44) have been verified experimentally by measurement on the tube of Fig. 8 with the value of $C$ at start oscillation used in the comparison with (44).

## Effect of Rbfiections on tine Efficiency

It has been observed that poor matches on a back-ward-wave oscillator produce variations in the power output and hence in the efficiency. In a nanner similar to that used to determine the effect of reflections on the starting conditions the effect on the efficiency will be determined in this section. For this case the efficiency must be determined as a function of $\Delta b$ and $\Delta y$.

Let us define the relation

$$
\begin{equation*}
G=G_{0}+G_{b} \Delta b+G_{y} \Delta y \tag{45}
\end{equation*}
$$

where $G_{0}$ is the value of $G$ without reflections corresponding to zero space charge, and $G_{b}$ and $G_{y}$ are the partial derivatives of $G$ with respect to $b$ and $y$, respectively, evaluated at $y=1.97$ and $b=1.52$. Let us define a correction factor $F$ as follows

$$
\begin{equation*}
F_{4}=\frac{\eta}{\eta_{0}}=\frac{\left|G_{0}\right|^{2}}{|G|^{2}} \tag{46}
\end{equation*}
$$

Then substituting from (45) we obtain

$$
\begin{equation*}
I_{4}=\frac{1}{\left|1+\frac{2 G_{b}}{G_{0}} \Delta b+\frac{2 G_{y}}{G_{0}} \Delta y\right|} \tag{47}
\end{equation*}
$$

By substituting (36) and (37) into (47) we obtain (48) which expresses $F$ as a function of $R$ and $\theta$

$$
\begin{equation*}
F_{4}=\frac{1}{\left|1+\frac{2 R}{G_{0}}\left[\frac{G_{b}}{B} \frac{\sin (\alpha-\theta)}{\sin (\alpha-\beta)}+\frac{G_{y}}{A} \frac{\sin (\beta-\theta)}{\sin (\beta-\alpha)}\right]\right|} \tag{48}
\end{equation*}
$$

The maximum or minimum value of $F$ can be found easily when the second term of the denominator is less than unity. For this case it is approximately true that the components which are at right angles with unity are negligible and only the real part of the second term is important. To simplify the expression let us define

$$
\begin{align*}
& B_{3}=\text { Real Part }\left[\frac{G_{b}}{G_{0}}\right]  \tag{49}\\
& A_{3}=\text { Real Part }\left[\frac{G_{y}}{G_{0}}\right] . \tag{50}
\end{align*}
$$

The definitions given by (49) and (50) may be substituted into (48) and the correction factor may be expressed in the following manner:

路
where $\theta_{0}$ is the phase at $\theta=0$ and is unnecessary to the present treatment. The efficiency is seen to vary in a sinusoidal manner about the mean position with the maximum efficiency and the minimum efficiency occurring when $\cos \left(\theta+\theta_{0}\right)= \pm 1$. Since values of $A$ and $B$ are already available, it is only necessary to evaluate $A_{3}$ and $B_{3}$. The parameters of (49) and (50) may be evaluated at the point $y=1.97$ and $b=1.52$ to yield the results

$$
\begin{align*}
G_{y}(1.97,1.52) & =1.227 e^{j 297.3 \text { degrees }}  \tag{52}\\
G_{b}(1.97,1.52) & =1.00 e^{j 191.5 \text { degrees }}  \tag{53}\\
G_{0}(1.97,1.52) & =1.53 e^{j 318 \text { degrees }} \tag{54}
\end{align*}
$$

Using (49) and (50) we may then evaluate $A_{3}$ and $B_{3}$ by substituting the results contained in (52), (53), and (54). Finally, from (51) we obtain the equation for the reduction factor

$$
\begin{equation*}
F_{4} \cong 1+1.42 R \cos \left(\theta+\theta_{0}\right) \tag{55}
\end{equation*}
$$

The reduction factor does indeed vary in a sinusoidal manner about the zero reflection case. All of the previous work has been based on the assumption that

$$
1.42 R \ll 1
$$

ereby restricting the region of validity of the result ven by (55). No quantitative experimental verification (55) has been made although qualitative agreement s been observed.

## Conclusion

The results presented in this paper enable the designer a backward-wave oscillator to predict the efficiency dl output power at the time of the initial design. Thus ckward-wave oscillators can now be designed not only t the basis of whether they will oscillate but also on the sis of how much rf power they will produce. The effect the important parameters which control the effiency has been presented. The effect of space charge has en considered to be the primary factor and all other ctors have been considered to produce corrections to is basic efficiency. The basic efficiency $\eta_{0}$ may be obined from Fig. 3 when the value of the space-charge rameter $Q C$ is calculated. The correction factors, $F_{1}$, , and $F_{3}$ for circuit loss, beam thickness, and velocity read are presented on Figs. 4, 5, 6, and 9, respectively.

By making use of these curves the designer can obtain a reasonably accurate value for the efficiency. If the matches are imperfect then the magnitude of the variations in efficiency can be obtained from (55). Eqs. (42) and (43) also permit the designer to ealculate the magnitude of the variations in the starting current and the frequency. These results therefore permit the design of backward-wave oscillators with considerably more confidence than was formerly possible.

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# Che Effects of Junction Shape and Surface Recombination on Transistor Current Gain-Part II* 

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Summary-Previous work demonstrated the importance of sure recombination and junction shape on the transistor current aplification factor $\alpha$ by means of a two-dimensional conducting per analog. ${ }^{1}$ This is now extended theoretically and experimentally other cases which are also of practical importance.
Exact analytical solutions have been obtained for the collector--base current amplification factor, $\alpha_{c b}$, for plane-parallel (grown action) transistors of rectangular and round cross section including rface and volume recombination. For the case in which the surface combination velocity $s$ is small and the volume lifetime $\tau$ is large, ese equations reduce to

$$
\alpha_{c b}=\left(\frac{s}{K}+\frac{T}{\tau}\right)^{-1}
$$

tere $K$ and $T$ are geometrical constants simply related to the base dth $W$ and the cross-sectional area. The range of validity of this uation has been investigated.
For the more complex geometries usually found in alloy transis$: s, \alpha_{c b}$ is found to be of the form:

$$
\alpha_{c b}=\left(\frac{s}{K}+\frac{T}{\tau}\right)^{-1}[1+F(\tau, s)],
$$

dere $K$ and $T$ are again geometrical constants. $F(\tau, s)$ is a small sitive correction term which goes to zero as $s$ approaches zero and tpproaches infinity.
Although the constants $K$ and $T$ are not readily evaluated ana:ically, in most practical cases volume recombination can be neg:ted, and then

$$
\alpha_{c b}=\frac{K}{s}[1+f(s)] .
$$

By means of a three-dimensional electrolytic conductance analog, $K$ and $f(s)$ have been evaluated for various geometries typical of alloy junction transistors. Since $f(s)$ is small, the constant $K$ serves as a geometrical figure of merit against which various emitter-collector configurations can be judged. The most striking result of this study is that for a given minimum junction spacing $W$ and given emitter and collector diameter, the value of $K$, and hence $\alpha_{r b}$, can be varied over a considerable range by changes in emitter and collector penetration. In particular, the highest figure of merit is obtained by combining essentially nonpenetrating emitters (i.e., lying on the wafer surface) with collectors which penetrate as far as necessary to give the required minimum spacing $W$.

## Introduction

$\Gamma$HIS PAPER presents an extension, both theoretical and experimental, of a study reported in a previous paper. ${ }^{1}$ In the theoretical portion the general analytical approach to the determination of the fate of injected minority carriers in semiconductor devices, in the steady state, is formulated. It is carried to explicit solution for certain simple geometries, corresponding to transistors of the grown junction type with
rectangular or circular cross section. In the case of geometrically more complex structures, such as the alloy junction TA-153 transistor, ${ }^{2}$ an implicit perturbation treatment is applied to deduce the general form to which the dependence of current gain on the surface recombination velocity and bulk lifetime must reduce for small values of the former and large values of the latter. These results prove of considerable use in analyzing the experimental findings.

On the experimental side, the conductance analog introduced in Part I is made three-dimensional by an adaptation of the familiar electrolytic tank. In this way it is possible to deduce the transistor current gain $\alpha$ directly from measured tank currents, and field plotting is then unnecessary. This represents a large saving in effort and an increase in accuracy, once the tank has been properly constructed. Presented here are the results of such an analog survey of the TA-153 alloy transistor, carried out to establish the dependence of $\alpha$ on surface recombination velocity and certain geometrical factors in this device. It is found that the shape of the junctions is quite important in this regard; for a given minimum emitter-collector separation, flatness (small penetration) of emitter offers considerable advantage in limiting loss of minority carriers due to surface recombination.

## Formulation of the Problem

The fate of injected minority carriers in semiconductors is of concern in numerous important instances, both from the theoretical and the practical points of view. This paper deals with the particular case of injected minority carriers in junction transistors when the system is in the steady (time independent) state and when the motion of the carriers is diffusion controlled. For simplicity of expression, the text will speak of holes in $p-n-p$ devices and will denote by $P$ the excess hole density in the $n$-regicn. The results will he equally valid in the $n-p-n$ case, by suitable interchanging of electrons for holes and use of the applicable diffusion coefficients.

In general, the emitter current of the $p-n-p$ transistor shown schematically in Fig. 1 consists of holes injected into the base and electrons flowing from the base into the emitter. Since in alloy transistors the emitter section has very high conductivity compared to the base, the emitter current consists almost entirely of holes. ${ }^{3}$ It will be assumed that the number of holes injected per second into the base at the emitter constitutes the entire emitter current $I_{e}$. Those holes which survive the trip to the collector comprise the collector current $I_{c}$. Some, however, are lost by volume recombination in the base region. Their total number per second is designated as the volume current, $I_{v}$. Still others are lost by surface recombination at the free surface of the base region. This

[^46]total number per second is designated as the surface current, $I_{s}$. These currents are related by
\[

$$
\begin{equation*}
I_{s}=I_{c}+I_{v}+I_{s} \tag{1}
\end{equation*}
$$

\]

and the sum of the volume and surface currents constitutes the base current, $I_{b}$.

$$
\begin{equation*}
I_{b}=I_{v}+I_{s} \tag{2}
\end{equation*}
$$



Fig. 1-Schematic representation of an alloy transistor.
Two particular ratios of currents are of interest in transistor performance:

$$
\begin{align*}
& \alpha_{c b}=\frac{I_{c}}{I_{e}}=1-\frac{I_{b}}{I_{e}}=\frac{\alpha_{c b}}{1+\alpha_{c b}}  \tag{3}\\
& \alpha_{c b}=\frac{I_{c}}{I_{b}}=\frac{\alpha_{c \theta}}{1-\alpha_{c e}}, \tag{4}
\end{align*}
$$

and it is the purpose here to extend the study, begun in Part I, of the dependence of these quantities on surface recombination velocity, bulk lifetime, and geometry.

In the diffusion controlled steady state, the lehavior of the system is contained in the differential equation ${ }^{4}$

$$
\begin{equation*}
D \nabla^{2} P-\frac{P}{\tau}=0 ; \text { in the base volume, } \tag{5}
\end{equation*}
$$

subject to the boundary conditions:

$$
\begin{align*}
& -\overrightarrow{D \nabla P} \cdot \vec{n}=s P ; \text { at the free base surface, }  \tag{6}\\
& P=P_{\bullet} ; \text { at the emitter junction surface, }  \tag{7}\\
& P=0 ; \text { at the collector junction surface. } \tag{8}
\end{align*}
$$

Here $D$ is the diffusion constant for holes, $\tau$ the bulk lifetime, $s$ the surface recombination velocity, and $\vec{n}$ the unit surface vector. The desired currents for evaluating (3) and (4) are then certain surface and volume integrals of the solution of these equations, i.e.,
$I_{e}=q D \int|\nabla P| d \sigma ;$ over emitter junction surface,
$I_{c}=q D \int|\nabla P| d \sigma ;$ over collector junction surface,

- Moore and Pankove, loc. cit.; and Shockley, op. cit., p. 320.
$I_{v}=\frac{q}{\tau} \int P d V$; over the base volume,
$I_{s}=q s \int P d \sigma ;$ over free base surface.
Here $q$ is the charge of a hole.


## Solution for Special Cases

Eq. (5) with its boundary conditions (6), (7), and (8) can be solved explicitly for certain geometrically simple cases. ${ }^{5}$ Here the computed $\alpha_{c e}$ for two of these is presented, leaving the mathematical details for the Appendices. These geometries correspond to grown (flat) junction types with circular and rectangular cross sections, respectively.

For the circular case, with the emitter-collector spacing $W$ and cross-sectional radius $R$,

$$
\begin{equation*}
\alpha_{c e}=1-\frac{\sum_{n} \alpha_{n} \tanh \left[\frac{W}{2 R} \sqrt{n^{2}+R^{2} / D \tau}\right]}{\sum_{n} \alpha_{n} \operatorname{coth}\left[\frac{\Pi}{R} \sqrt{n^{2}+R^{2} / D \tau}\right]} ; \tag{13}
\end{equation*}
$$

where the sums are over all the successive positive roots of an equation involving the zeroth-order Bessel function, $J_{0}$;

$$
\begin{equation*}
n J_{0}{ }^{\prime}(n)+\frac{R s}{D} J_{0}(n)=0 \tag{14}
\end{equation*}
$$

and where

$$
\begin{equation*}
\alpha_{n}=\frac{\sqrt{n^{2}+R^{2} / D_{\tau}}}{n^{2}\left(n^{2}+R^{2} s^{2} / D^{2}\right)} . \tag{15}
\end{equation*}
$$

Similarly, for the rectangular case, with cross sectional dimensions $2 a$ and $2 b,{ }^{6}$


Fig. 2-Comparison of theoretically computed current gain and measurements in the tank of Fig. 3. The abscissa gives surface recombination velocity $s$ in terms of the linear scale-up factor $k$.

$$
\begin{equation*}
\alpha_{m}=\frac{\sin ^{2} m b}{m^{2} b+\frac{s}{D} \cos ^{2} m b} \tag{18}
\end{equation*}
$$

As a special case of (16), let $s \rightarrow 0$ on the face $y=b$, and let $\tau \rightarrow \infty$ (no volume recombination). 'This gives the result for a "two-dimensional" transistor in which vol-

$$
\begin{equation*}
\alpha_{c e}=1-\frac{\sum_{n} \sum_{m} \alpha_{n} \alpha_{m} \sqrt{n^{2}+m^{2}+1 / D_{\tau}} \tanh \left[\frac{W}{2} \sqrt{n^{2}+m^{2}+1 / D_{\tau}}\right]}{\sum_{n} \sum_{m} \alpha_{n} \alpha_{m} \sqrt{n^{2}+m^{2}+1 / D_{\tau}} \operatorname{coth}\left[W \sqrt{n^{2}+m^{2}+1 / D_{\tau}}\right]} ; \tag{16}
\end{equation*}
$$

where the sums are over all the positive roots of the equations

$$
\begin{equation*}
n \tan n a=\frac{s}{D}, \quad m \tan m b=\frac{s}{D} \tag{17}
\end{equation*}
$$

and where

$$
\alpha_{n}=\frac{\sin ^{2} n a}{n^{2} a+\frac{s}{D} \cos ^{2} n a}
$$

[^47]ume recombination is negligible:
\[

$$
\begin{equation*}
\alpha_{c e}=1-\frac{\sum_{n} n \alpha_{n} \tanh \left(\frac{n \|}{2}\right)}{\sum_{n} n \alpha_{n} \operatorname{coth}(n W)} . \tag{19}
\end{equation*}
$$

\]

In the experimental part of this bulletin, the result (19) is compared with measurements in an electrolytic tank. The behavior of this series is depicted in Fig. 2, above, for a particular choice of $W$ and $a$.

## Perturbation Approach

In geometrically complex arrangements, such as the TA-153 alloy junction transistor, explicit and analytical
solution has not been achieved. ${ }^{7}$ Nevertheless, interesting general results can be deduced for such structures by a somewhat implicit use of the perturbation (iteration) approach. Assume that when there is no volume or surface recombination ( $\tau=\infty, s=0$ ) the solution to (5)-(8) is $P^{0}$. When both recombinations are present, a first approximation, whose validity is related to the largeness of $\tau$ and the smallness of $s$, consists of substituting $P^{0}$ into (9)-(12) to compute the currents. The integrals are then independent of $\tau$ and $s$, and lead to

$$
\begin{equation*}
\alpha_{c s} \approx 1-\frac{T}{\tau}-\frac{s}{K} \tag{20}
\end{equation*}
$$

or

$$
\begin{equation*}
\alpha_{c b} \approx\left(\frac{T}{\tau}+\frac{s}{K}\right)^{-1} \tag{21}
\end{equation*}
$$

where

$$
\begin{equation*}
T=\frac{\int_{v}^{P^{0} d V}}{D \int_{e}\left|\nabla P^{0}\right| d \sigma} \tag{22}
\end{equation*}
$$

and

$$
\begin{equation*}
K^{-1}=\frac{\int_{s}^{P^{0} d \sigma}}{D \int_{e}\left|\nabla P^{0}\right| d \sigma} \tag{23}
\end{equation*}
$$

are certain geometrical factors.
The approximate expressions (20) and (21) can be made formally correct in the following way. The actual hole concentration in the volume and on the surface will always be smaller than $P^{0}$ when recombination is taking place, since $P^{0}$ neglects recombination. Hence the insertion of $P^{0}$ into (11) and (12) progressively overestimates the volume and surface currents as $\tau$ decreases and $s$ increases. Thus it can be expected, e.g., that

$$
\begin{equation*}
\alpha_{c b}=\left(\frac{T}{\tau}+\frac{s}{K}\right)^{-1}[1+F(\tau, s)] \tag{24}
\end{equation*}
$$

where $F(\tau, s)$ is some positive function which vanishes as both $\tau \rightarrow \infty$ and $s \rightarrow 0$. The corrected $\alpha_{c e}$ is similarly obtained by using (24) in (3).

In simple cases one can evaluate the constants $T$ and $K$ analytically. As a demonstration one can consider the structures for which explicit solutions were obtained in the previous section, i.e., plane parallel junctions of circular and rectangular cross section. The zero-order hole density is

$$
\begin{equation*}
P^{0}(Z)=P \cdot \frac{W-Z}{W} \tag{25}
\end{equation*}
$$

[^48]where $Z$ is the distance from the emitter. Then the integrals required in (22) and (23) are
\[

$$
\begin{align*}
& \int_{V} P^{0} d V=P_{e} \frac{W}{2} \times \text { cross-sectional area, }  \tag{26}\\
& \int_{e}\left|\nabla P^{0}\right| d \sigma=P_{e} \frac{1}{W} \times \text { cross-sectional area }  \tag{27}\\
& \int_{s} P^{0} d \sigma=P_{e} \frac{W}{2} \times \text { cross-sectional perimeter, } \tag{28}
\end{align*}
$$
\]

for both these cases. Consequently, the approximate form (20) becomes

$$
\begin{equation*}
\alpha_{c e} \approx 1-\frac{W^{2}}{2 D \tau}-\frac{W^{2} s}{D R} \tag{29}
\end{equation*}
$$

for the circular cross section; and

$$
\begin{equation*}
\alpha_{c e} \approx 1-\frac{W^{2}}{2 D \tau}-\frac{W^{2} s}{D}\left[\frac{1}{a}+\frac{1}{b}\right] \tag{30}
\end{equation*}
$$

for the rectangular. These expressions are just what one obtains from the series solutions (13) and (16) by neglecting all but the lead terms and by replacing the remaining functions by their small argument power expansions.

In more complicated cases, one can seek to evaluate the various quantities empirically. The remainder of this paper deals primarily with such a study of the TA-153. For the most part, it will be assumed that the bulk lifetime is sufficiently long so that volume recombination is negligible. Then (24) becomes

$$
\begin{equation*}
\alpha_{c b}=\frac{K}{s}[1+F(s)] ; \quad F(s) \rightarrow 0, s \rightarrow 0 \tag{31}
\end{equation*}
$$

and similarly

$$
\begin{equation*}
\alpha_{c e}=1-\frac{s}{K}\left[\frac{1}{1+f(s)+s / K}\right] \tag{32}
\end{equation*}
$$

Both the value of $K$ and the form of $f(s)$ are deduced from analog measurements.

## Electrolytic Tank Analog

In Part I it was shown that the process of diffusioncontrolled minority carrier flow in a semiconductor with surface recombination is analogous to field controlled electrical current flow in a conductive medium whose surface is divided into small, isolated segments (tabs), each leaking current to ground through suitable resistors. ${ }^{8}$ The analogs constructed at that time were twodimensional, and in order to deduce the three-dimensional result, it was necessary to plot fields and to weigh radially the computed local currents. Since that time three-dimensional analogs have been constructed by

[^49]

Fig. 3-Photograph of the electrolytic tank with plane-parallel geometry (grown-junction transistor analog).
adaptation of the familiar electroytic tank. In these the desired total currents can be measured directly.

The analogy between the diffusion and the conductance devices is given by the equation

$$
\begin{equation*}
\frac{s}{D}=k \frac{\rho}{A R} \tag{33}
\end{equation*}
$$

where $\rho$ is the electrical conductivity of the analog medium, $A$ is the area of a surface tab, $R$ is the bleeder resistance, and $k$ is the linear scale-up factor from the semiconductor. Variation in simulated $s$ can be achieved by changing either $\rho$ or $R$. In covering a wide range of values of $s$, it proves advisable to do both in order to keep the over-all resistance of the analog within reasonable bounds.

To imitate a surface of constant $s$, it is not necessary that all the tabs have the same size, but merely that the product of each tab area by its bleeder resistance have the same value. The analogy (33) is strictly valid only when the dimensions of each surface tab approach zero in any direction in which $P$ changes. In practice, of course, one must be satisfied with making such tab dimensions reasonably small. The tabs need not be small, however, along directions of constant $P$. When these directions are apparent, from symmetry or otherwise, the tabs can be taken as long in them as is desirable. This fact is used in the construction of the present a nalogs.


Fig. 4-Photograph of the electrolytic tank with geometry suitable for alloy transistor analog.


Fig. 5-Interior of the electrolytic tank of Fig. 4, showing the surface recombination rings and extra irterchangeable electrodes.

Two such tanks have been built and operated; a rectangular one for which the analytical solution is available, and a more complex one which represents typical alloy transistor geometry. These are shown in Figs. 3, 4, and 5; and are described in somewhat more detail in subsequent sections. Both are lucite tanks whose inside dimensions reproduce the ofter dimensions of the semiconductor base region. Electrolyte solution in the tank provides the conducting medium. Surface electrodes on the plastic walls were made by an initial spray deposit of silver, which was then built up by copper electroplating and finally given a light goldplate to retard corrosion. The tabs were formed by scribing through the metal plate, and each tab was provided with an external electrical contact by means of a small pin through the plastic. Similar plated surfaces served as emitter and collector electrodes in the rectangular tank. The curved emitter and collector electrodes in the

T: $1-153$ were in the form of detachable inserts, either of solid metal or coated plastic depending on their size.

## Measurfment of Currents

Fig. 6 is a schematic diagram of the current measuring circuit used in conjunction with the tank analogs. A 10 -ke sigmal of 2 to 10 volts was applied across $A B$, and $R_{c}$ and $R_{s}$ were each adjusted to give null readings against the equal-armed, pure resistive branch $R_{1} R_{2}$. The capacitors $C_{c}$ and $C_{s}$ were tuned to nullify the reactive components in each of the two tank paths, $E C$ and $E S$. At mutual balance the voltage drop across each path is the same, hence the transistor parameters of interest,

$$
\begin{align*}
& \alpha_{c b}=\frac{I_{c}}{I_{s}}=\frac{R_{s}}{R_{c}}  \tag{34}\\
& \alpha_{c e}=1-\frac{I_{s}}{I_{s}+I_{c}}=1-\frac{R_{c}}{R_{c}+R_{s}}, \tag{35}
\end{align*}
$$

are calculable (lirectly from the bridge readings.
Since electrolyte solutions decrease in resistivity by about $2 \frac{1}{2}$ per cent per degree $C$. rise at room temperature, ${ }^{9}$ it is necessary either to maintain the tank at a fairly constant temperature or to record the solution temperature at balance and compute the correction. The latter procedure was followed in the present work, the correction being applied to the value of $s$ through (33).

## Rectangular Analog and Results

The rectangular tank, Fig. 3 , was constructed for the purpose of gauging the adequacy of such analog measurements by comparison with values which can be computed analytically. The inside dimensions of the device are 10 inches by 2 inches by 2 inches. Emitter and collector are represented by the two plated, vertical faces; their effective size can be varied simply by adjusting the depth of elentrolyte in the tank. The tank bottom, which represents the surface of recombination, consists of 20 tabs ( 0.1 inch wide) cut parallel to the emitter and collector. Each tab is in electrical contact with the insulated terminals of the bleeder resistance plug board on the right side of the device. The terminals on the left are connected to a common bus bar, which represents $S$ in Fig. 6. The illustration shows a set of bleeder resistances in place.

A comparison between measured $\alpha_{c b}$ and that calculated from (19) and (4) is shown in Fig. 2 for a 10 -inch depth of electrolyte. These results were taken with two concentrations of solution having resistivities of 2,000 and $10,000 \mathrm{ohm} \mathrm{cm}$, and with bleeder resistors of from $1 K$ to $330 K$ ohms. Agreement between the analog measurement and theory is quite satisfactory.

[^50]
## TA-153 Analog

Fig. 4 shows the TA-153 analog in operating condition. As in the rectangular tank, the surface tabs are again led to the insulated terminals of the bleeder plug board on the side of the device and thence through the bleeder resistors to the common terminal bar below. The protruding central bolt serves both to hold the emitter (or collector) in place and to provide electrical contact to these electrodes. Thus it represents $E$ (or $C$ ) in Fig. 6. An opened view of the tank is provided by


Fig. 6-Circuit diagram for measurement of $\alpha_{c e}$ and $\alpha_{c b}$ with the electrolytic tank analog.

Fig. 5. It shows the emitter and collector inserts in place and also the surface tabs, which could be taken here in the form of concentric rings in view of the cylindrical symmetry. Additional emitter and collector inserts of various curvatures are to be seen in the foreground.


Fig. 7-Cross section through a typical alloy transistor; $X$ and $Y$ measure the maximum penetration of the emitter and collector respectively, while $W$ is the minimum separation.

The nominal dimensions of the TA-153 transistor, depicted in section in Fig. 7, are assumed to be: wafer thickness, 5 mils; emitter diameter, 15 mils; collector diameter, 44 mils; and the emitter and collector junction surfaces are idealized as spherical sections. The analog was constructed at a scale-up of 200:1; hence the inner width of the tank is 1 inch, and the emitter and collector inserts are 3 inches and 8.8 inches in diameter. The inner length and breadth of the tank were arbitrarily fixed at 16 inches. Measurements in Part I established that the carrier concentration at such distances from the cylindrical axis is quite small. For the same reason, recombination on the thin outer surfaces of the wafer was ignored, which means that these tank surfaces were not plated.

The surface tabs on the emitter side were taken as five rings of 0.1 -inch width from radius 1.5 inches (emitter edge) to radius 2 inches, then ten rings of 0.2 -inch width, and finally ten more of 0.4 -inch width. On the collector side nine rings of $0.4-\mathrm{inch}$ width were cut, starting at radius 4.4 inches (collector edge). The insulating gap produced by the scribing tool was 0.002 to 0.004 inch.

Previous work with the two-dimensional analog had demonstrated the importance of maintaining a large ratio of collector-to-emitter diameter. In the present work the geometrical variation studied most extensively was the emitter-and-collector penetration, i.e., the radii of curvature of the junction surfaces. $X$ and $Y$ denote the maximum emitter and collector penetration respectively, and $W$ the minimum emitter-collector separation, Fig. 7. With these conventions established, the device geometry can be specified by the sequence $X: W: Y$; for example, $1: 2: 2$ denotes such a TA-153 structure with 1 mil maximum emitter penetration, 2 mils minimum gap, and 2 mils maximum collector penetration.

The analog was arranged to represent $n$-type germanium, with $D_{p}$ taken to be $44 \mathrm{~cm}^{2} / \mathrm{sec}$. For convenience a standard set of bleeder resistors, $R^{*}$, was chosen, of such size that the magnitude of $s$ in $\mathrm{cm} / \mathrm{sec}$ equals half the magnitude of the electrolyte resistivity in ohm-cm:

$$
\begin{equation*}
|s|=\left|\frac{\rho}{2}\right| \tag{36}
\end{equation*}
$$

By (33) this required the $R$ for each ring to satisfy

$$
\begin{equation*}
A R^{*}=17,600 \text { ohm }-\mathrm{cm}^{2} \tag{37}
\end{equation*}
$$

$A$ being the ring area. Additional sets of bleeders were prepared so that

$$
\begin{equation*}
R_{t}=t R^{*} ; \quad t=2,5,10,20,50,100,200 . \tag{38}
\end{equation*}
$$

With these,

$$
\begin{equation*}
|s|=\left|\frac{\rho}{2}\right| \times \frac{1}{t} . \tag{39}
\end{equation*}
$$

Measurements were made using all these bleeders in conjunction with two concentrations of electrolyte having resistivities of 20,000 and 3,000 ohm-cm respectively. In this way the range of $s$ up to $10,000 \mathrm{~cm} / \mathrm{sec}$ was surveyed. The results are presented and discussed in the following sections. They are not limited to $n$-type germanium devices, but can easily be applied to a material with minority carrier mobility $D \mathrm{~cm}^{2} / \mathrm{sec}$ by renormalizing $s$ as follows:

$$
\begin{equation*}
s=s\left(n-G_{e}\right) \times \frac{D}{44} \tag{40}
\end{equation*}
$$

## Results

The dependence of $\alpha_{c s}$ and $\alpha_{c b}$ on $s$ has been measured in this analog on fifteen geometrical variations of the

TA-153, comprising all the compatible combinations of integral values of $X$ and $Y$ from 0 to 4 mils. Fig. 8 shows a plot of the experimentally determined values of $\alpha_{c e}$ against ln $s$ for the case, 1:2:2, in both the forward and the reverse directions. Fig. 9 shows $\ln \alpha_{c b}$ against $\ln s$ for the same structure. Similar plots are obtained in all instances.


Fig. 8- $\alpha_{\text {ce }}$ as a function of $s$ for a typical alloy transistor with $X=1 \mathrm{mil}, W=2 \mathrm{mils}$, and $Y=2$ mils.


Fig. 9- $\alpha_{c b}$ as a function of $s$ for the same geometry as in Fig. 8.
The dashed line shows the linear approximation $\alpha_{c b}=K s^{-1}$.

Analysis of the results along the lines suggested by perturbation theory, (31), discloses that in this general geometry and up to the maximum value of $s$ investigated, $10,000 \mathrm{~cm} / \mathrm{sec}$, the dependence of $\alpha_{c b}$ on $s$ in the forward direction is very adequately represented by an equation of the form:

$$
\begin{equation*}
\alpha_{c t}=\frac{K}{s}[1+b \sqrt{s}], \tag{41}
\end{equation*}
$$

where $K$ and $b$ are geometrically dependent. The values of these parameters for the geometries investigated are shown in Table I. A typical comparison between the experimental $\alpha_{c b}$ and those computed from (41), for the case $1: 2: 2$, is contained in Table II.

TABLE I
Current Gain in TA-153; $\mathrm{X}: \mathrm{W}^{\prime}$ : I' (Forward)

$$
\alpha_{c b}=\frac{K}{s}[1+b \sqrt{s}]
$$

K

|  |  |  |  |  |  |
| :---: | ---: | ---: | ---: | ---: | :---: |
|  | 11 | 1 | 2 | 3 | 4 |
|  |  |  |  | 5 |  |
| 0 | 49,470 | 19,170 | 11,230 | 7,820 | 4,490 |
| 1 | 22,500 | 11,160 | 8,640 | 5,150 |  |
| 2 | 14,900 | 9,410 | 5,580 |  |  |
| 3 | 10,730 | 6,230 |  |  |  |
| 4 | 7,370 |  |  |  |  |

$b$

|  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
|  |  |  |  |  |  |
|  |  |  |  |  |  |
| 0 | 0.01211 | 0.01140 | 0.01290 | 0.01722 | 0.02214 |
| 1 | 0.00939 | 0.01263 | 0.01443 | 0.01635 |  |
| 2 | 0.01232 | 0.01430 | 0.01637 |  |  |
| 3 | 0.01566 | 0.01797 |  |  |  |
| 4 | 0.02016 |  |  |  |  |

TABLE II
Current Gain in TA-153; 1:2:2 (Forwari))

$$
\alpha_{c b}=\frac{11,160}{s}[1+0.01263 \sqrt{ } \sqrt{s}]
$$

| $S$ | $\alpha_{c b}$ (measured) |
| :---: | :---: |
| 8,460 | 2.86 |
| 4,260 | 4.79 |
| 1,700 | 9.88 |
| 1,400 | 11.46 |
| 850 | 18.15 |
| 700 | 21.0 |
| 425 | 34.4 |
| 280 | 47.7 |
| 170 | 76.9 |
| 140 | 91.9 |
| 86 | 150 |
| 70 | 174 |
| 43 | 287 |
| 28 | 414 |
| 14 | 8.3 .78 |
| 7 | 1,590 |

For given value of $s$, the bracketed quantity in (41) does not vary greatly from one geometry to another. Thus the geometrical dependence of $\alpha_{c b}$ is essentially contained in $K$; one can take this number as a geometrical figure of merit. For example, with a given emitter shape (given $X$ ) the size of $K$ increases with decreasing $W$, as is to be expected. In addition, however, $K$ is markedly affected by junction shape; and favorable
junction geometry can compensate for considerable disadvantage in $W$. As a particular instance, $1: 2: 2$ is about twice as good as $3: 2: 0$, and is as good as $3: 1: 0$. The general situation with regard to junction geometry can be summed up in the following rule: minority carrier loss through surface recombination in the TA- 153 can be held down by keeping the emitter junction as flat as possible (consistent with good junction properties) and by closing the emitter-collector gap by collector penetration.

By means of (41) one can compute the maximum $s$ tolerable for a given $a_{c b}$, when volume recombination is negligible. 'Table Ill shows the geometrical dependence of this value of $s$, for the particular case $\alpha_{c b}>30$.

TABLE III
Maximum Tolerable sin TA-153; X: И: $Y$ (for)
$\alpha_{c b} \geqq 30$

|  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: |
| $X$ | 1 | 2 | 3 | 4 |
| 0 | 2,700 | 850 | 480 | 340 |
| 1 | 970 | 470 | 370 | 210 |
| 2 | 650 | 400 | 230 |  |
| 3 | 480 | 270 |  |  |
| 4 | 340 |  |  |  |

In the reverse direction, $\alpha_{c b}$ is similarly expressible as

$$
\begin{equation*}
\alpha_{c b}=\frac{K}{s}\left[1+b s^{0.7}\right] . \tag{+2}
\end{equation*}
$$

Table IV contains the values of the observed constants, and Table $V$ (next page) gives a comparison between observed and computed quantities.

TABLE IV
Current Cain in TA-153; $X: W: Y$ (Reverse)
$\alpha_{c b}=\frac{k}{s}\left[1+b s^{0.7}\right]$
K

|  | $W$ | 1 | 2 | 3 | 4 |
| :---: | :---: | :---: | :---: | :---: | :---: |
| 0 |  |  |  | 5 |  |
| 0 | 320 | 200 | 145 | 130 | 120 |
| 1 | 295 | 180 | 165 | 135 |  |
| 2 | 245 | 200 | 150 |  |  |
| 3 | 235 | 200 |  |  |  |
| 4 | 235 |  |  |  |  |

$b$

|  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
| X | W | 1 | 2 | 3 | 4 |
| 0 | 0.0163 | 0.0139 | 0.0172 | 0.0159 | 0.0175 |
| 1 | 0.0139 | 0.0178 | 0.0154 | 0.0169 |  |
| 2 | 0.0176 | 0.0159 | 0.0179 |  |  |
| 3 | 0.0181 | 0.0164 |  |  |  |
| 4 | 0.0203 |  |  |  |  |

TABLE V
Current Gain in TA-153; 1:2:2 (Reverse)

$$
\alpha_{c b}=\frac{180}{s}\left[1+0.0178 s^{0.7}\right]
$$

| $S$ | $\alpha_{r b}$ (measured) | $\alpha_{c l}$ (calculated) |
| :---: | :---: | :---: |
| 8,460 | 0.2 .35 | 0.2 .33 |
| 4,260 | 0.307 | $0 . .303$ |
| 1,700 | 0.460 | 0.451 |
| 1,400 | 0.484 | 0.494 |
| 850 | 0.609 | 0.638 |
| 700 | 0.686 | 0.708 |
| 425 | 0.928 | 0.953 |
| 280 | 1.17 | 1.24 |
| 170 | 1.75 | 1.74 |
| 140 | 3.01 | 2.01 |
| 86 | 3.32 | 3.94 |
| 70 | 7.61 | 3.47 |
| 43 | 14.61 | 7.24 |
| 28 | 28.8 | 14.61 |
| 14 |  | 27.5 |
| 7 |  |  |
|  |  |  |

Interpretation of Geomitrical. Dependence
From (29) and (30) it may be seen that, for the simple geometries discussed in that section,

$$
\begin{equation*}
K \propto \frac{1}{W^{2}} . \tag{43}
\end{equation*}
$$

In the TA- 153 the relation between $K$ and $W$ is more complicated. The perturbation treatment, however, in pointing up the geometrical source of $K$, sheds considerable light on the observed results in Table I.

In the simple geometry the inverse square dependence in (43) arises from two factors. First, the collector current depends on the hole concentration gradient between emitter and collector, and this varies as $W^{-1}$, (27). Second, the surface current depends on the effective free surface area, and this varies with $W$, (28). These factors cannot be changed independently when the cross section is fixed, thus yielding the inverse square dependence of (43).

In the TA-153, $K$ similarly depends on these currents. Here, however, the geometry is of such a nature that considerable independent variation is possible. The collector current depends on an average hole gradient between emitter and collector, and this can be expected to change in some inverse manner with $W$, as in Fig. 10. The relation of surface current to $W$, however, is not direct. As established in Part I, the surface current is confined essentially to a small, effective region of the free surface around the emitter. This region is indicated by $S$ in Fig. 10. The extent of $S$ is determined by competition between the surface and the collector, $C$, for the holes injected near the outer edge of the emitter, $E$. Thus, the surface current is expected to depend largely upon the nearness of $C$ to $S$, and this can change or not with $W$. For example, if $W$ is reduced by moving the emitter to $E^{\prime}$, essentially only the collector current is changed. Hence in such cases $K$ should vary in some
duced to a similar extent by changing the collector to $C^{\prime}$, not only is the collector current changed but also the effective surface, i.e., the nearness of $C$ to $s$. In this case, therefore, a higher order dependence of $K$ on $W$ is expected.


Fig. 10-Cross section through an alloy transistor showing how pursible variations in geometry can change the surface currents alle. $W$ independently.

The values of $K$ given in Table $I$ are plotted vs $W$ in Fig. 11 to show the observed geometrical dependence of this figure of merit in the TA-153. The results are compared there with a grid representing
$K=\frac{35 \times 10^{4}}{(W+6)[1.35(W+X)-0.35]}$, and $W+X=5-Y$.
Now the use of this expression is intended merely to be suggestive, the actual relation between $K$ and geometry being surely more complex. In particular (44) should not be used loosely for any extended extrapolation to smaller $W$. Nevertheless, it is a simple function whose behavior shows good qualitative agreement with the observed results, and one which gives a more mathematical statement to the line of argument presented above. When $W$ is varied with $Y$ constant on this grid, the second term in the denominator of (44) is unchanged, and $K$ varies as $\left(W^{7}+6\right)^{-1}$. Alternatively, when $W$ is varied with $X$ constant, both terms in the denominator change, and $K$ varies inversely as a quadratic in $W$. Thus the two factors in the denominator simulate respectively the influence of geometry on the collector current and on the effective surface area, $S$.

## Volume Recombination in TA-153

The analog measurements of the minority carrier loss in the TA-153 apply strictly to a base material of bulk Refifetime sufficiently long so that surface recombination
completely over-shadows volume recombination. It shall now be shown how one can make a rough estimate of the requirement this puts on $\tau$. For this purpose it is convenient to use the approximate expression for $\alpha_{c e}$, (20). The size of the surface term is known from Table I. The volume term can, in fair approximation, be taken in the form of the volume term in (29), if a suitable choice of an equivalent cross section and emitter-collector spacing is inserted into (26) and (27). For simplicity, take the actual cross-sectional area of the emitter and the minimum spacing $W$. Then,

$$
\begin{equation*}
\alpha_{c e}=1-\frac{s}{K}-\frac{W^{2}}{2 D \tau} \tag{45}
\end{equation*}
$$

or, explicitly for $n$-type germanium,

$$
\begin{equation*}
\alpha_{c}=1-\frac{s}{K}-\frac{0.075 W^{2}}{\tau}, \tag{46}
\end{equation*}
$$

with $W$ in mils and $\tau$ in $\mu$ sec.


Fig. 11-The dependence of $K$ on $W$ with $X$ or $Y$ held fixed.
The use of a $W^{2}$ term for the volume recombination in the TA-153 can be justified, in essence, by arguments similar to those used in the previous section. The square of $W$ in the volume term represents again the product of two factors: an increase in $W$ both decreases the velocity of diffusion (decreases the gradient) and increases the path length. These factors, which determine the transit time, vary simultaneously both in the simple geometry and in the TA-153.

As an example of such an estimate of the relative importance of surface and volume terms, consider TA-153; 1:2:2 with $\tau=100 \mu \mathrm{sec}$. Insertion of the proper figures into (46) shows the surface loss to be about 10 times the volume loss when $s=330$, but to be about equal to the volume loss when $s=33$. Since in most cases practical values of $s$ are in the range of a few hundred to one thousand, the assumption of negligible volume loss is a good one.

## Appendix A

Consider the grown (fat) junction type with emittercollector separation $W$ and circular cross section of radius $R$. A cylindrical co-ordinate system is used, with the $z$-axis perpendicular to the junction faces; and for convenience reduced units with $R$ as the unit length are introduced; abbreviating:

$$
\begin{equation*}
g=\frac{W}{R} \quad h=\frac{R s}{D} \tag{47}
\end{equation*}
$$

Then (5)-(8) become

$$
\begin{align*}
& \frac{1}{r} \frac{\partial}{\partial r}\left(r \frac{\partial P}{\partial r}\right)+\frac{\partial^{2} P}{\partial Z^{2}}-\frac{R^{2} P}{D \tau}=0  \tag{48}\\
& \frac{\partial P}{\partial r}=-h P ; \quad r=1, \quad 0<Z<g  \tag{49}\\
& P=P_{e} ; \quad r \leqq 1, Z=0  \tag{50}\\
& P=0 ; \quad r \leqq 1, Z=g \tag{51}
\end{align*}
$$

This set of equations can now be solved by any of the standard methods; here it was chosen to apply transform techniques. By means of the Hankel Transform, ${ }^{10}$

$$
\begin{equation*}
\bar{P}(n, Z)=\int_{0}^{1} P(r, Z) r J_{0}(n r) d r \tag{52}
\end{equation*}
$$

defined in terms of the zero order Bessel function $J_{0}$, (48) is converted to

$$
\begin{equation*}
\frac{\partial^{2} \bar{P}}{\partial Z^{2}}-\left(n^{2}+\frac{R^{2}}{D \tau}\right) \bar{P}=0 ; \tag{53}
\end{equation*}
$$

provided that by (49), $n$ is chosen such that

$$
\begin{equation*}
n J_{0}^{\prime}(n)+h J_{0}(n)=0 \tag{54}
\end{equation*}
$$

Conditions (50) and (51) become

$$
\begin{align*}
& \bar{P}(n, 0)=P \cdot \frac{J_{1}(n)}{n}  \tag{55}\\
& \bar{P}(n, Z)=0 . \tag{56}
\end{align*}
$$

The solution of (53) which satisfies (55) and (56) is readily found to be

$$
\begin{equation*}
\bar{P}(n, Z)=P_{e} \frac{J_{1}(n) \sinh \left[(g-Z) \sqrt{n^{2}+R^{2} / D_{r}}\right]}{\sinh \left[g \sqrt{n 2+R^{2} / D_{r}}\right]} \tag{57}
\end{equation*}
$$

[^51]Then inversion ${ }^{11}$ gives the required hole density:

$$
\begin{equation*}
P(r, Z)=2 \sum_{n} \frac{h}{n^{2}+h^{2}} \frac{n J_{0}(n r)}{J_{0}(n) J_{1}(n)} \bar{P}(n, Z), \tag{58}
\end{equation*}
$$

where the sum is over the successive positive roots of (54).

The currents required by (9) and (10) are

$$
\begin{equation*}
I_{e}=\left.2 \pi \int_{0}^{1} r \frac{\partial P}{\partial Z}\right|_{Z=0} d r \tag{59}
\end{equation*}
$$

and

$$
\begin{equation*}
I_{c}=\left.2 \pi \int_{0}^{1} r \frac{\partial P}{\partial Z}\right|_{z=g} d r, \tag{60}
\end{equation*}
$$

and when these operations are performed on (58) and the results sulstituted into (3), (13) results.

## Appindix 13

Consider the grown (flat) junction type with emittercollector separation $W$ and rectangular cross section of semi-dimensions $a$ and $b$. A rectangular co-ordinate system with the $z$-axis perpendicular to the junction faces is used. Then (5)-(8) become
$\left(\frac{\partial^{2}}{\partial x^{2}}+\frac{\partial^{2}}{\partial y^{2}}+\frac{\partial^{2}}{\partial z^{2}}\right) P-\frac{P}{D \tau}=0$
$\frac{\partial P}{\partial x}=\frac{-s}{D} P ; \quad x=a, 0<y<b, 0<Z<W$
$\frac{\partial P}{\partial x}=0 ; x=0,0<y<b, 0<Z<W^{*}$
$\frac{\partial P}{\partial y}=-\frac{s}{D} P ; \quad 0<x<a, y=b, 0<\%<W$
$\frac{\partial P}{\partial y}=-\frac{s}{D} P ; \quad 0<x<a, y=0 ; \quad 0<Z<W$
$P=P_{e} ; \quad 0<x<a, 0<y<b, Z=0$
$P=0 ; \quad 0<x<a, 0<y<b, Z=W$.
By means of the double Cosine Transform ${ }^{12}$
$\overline{\bar{P}}(n, m, Z)=\int_{0}^{a} d x \int_{0}^{b} d y P(x, y, Z) \cos n x \cos m y$.
Eq. 61 is converted to

$$
\begin{equation*}
\frac{\partial^{2} \overline{\bar{P}}}{\partial Z^{2}}-\left(n^{2}+m^{2}+\frac{1}{D \tau}\right) P=0, \tag{67}
\end{equation*}
$$

provided that by (62) and (63) $n$ and $m$ are chosen such that

$$
\begin{align*}
& n \tan n a=\frac{s}{D} \\
& m \tan m b=\frac{s}{D} . \tag{68}
\end{align*}
$$

Conditions (64) and (65) become

$$
\begin{align*}
\overline{\bar{P}}(n, m, 0) & =P_{e} \frac{\sin n a}{n} \frac{\sin m a}{m}  \tag{69}\\
\overline{\bar{P}}(n, m, W) & =0 . \tag{70}
\end{align*}
$$

The solution of (67) which satisfies (69) and (70) is

$$
\begin{align*}
P(n, m, Z) & =P_{e} \frac{\sin n a}{n} \frac{\sin m b}{m} \\
& \times \frac{\sinh [W-Z] \sqrt{n^{2}+m^{2}+1 / \overline{D \tau}}}{\sinh \left[W \sqrt{n^{2}+m^{2}+1 / D \tau}\right.} . \tag{71}
\end{align*}
$$

Then inversion gives ${ }^{13}$

$$
\begin{align*}
& P(x, y, Z) \\
& \quad=4 \sum_{n} \sum_{m} \frac{n}{n a+\left(s \cos ^{2} n a / n D\right.} \\
& \quad \times \frac{m}{m b+\left(s \cos ^{2} m b / m D\right.} \cos n x \cos m y \overline{\bar{P}}(n, m, Z) \tag{72}
\end{align*}
$$

where sums are over the successive positive roots of (68).
The currents required by (9) and (10) are

$$
\begin{equation*}
I_{e}=\left.\int_{0}^{a} d x \int_{0}^{b} d y \frac{\partial P}{\partial Z}\right|_{z=0} \tag{73}
\end{equation*}
$$

and

$$
\begin{equation*}
I_{c}=\left.\int_{0}^{a} d x \int_{0}^{b} d y \frac{\partial P}{\partial Z}\right|_{Z=w}, \tag{74}
\end{equation*}
$$

and when these operations are performed on (72) and the results substituted into (3), (16) results.

When $s \rightarrow 0$ on the surface $y=b$, one need only consider the first term in the $m$ series. As $s \rightarrow 0$, this first value of $m$ behaves as

$$
\begin{equation*}
m=\sqrt{\frac{s}{b D}} \tag{75}
\end{equation*}
$$

Then, allowing $\tau \rightarrow \infty$ as well, (72) reduces to
$P(x, y, Z)=2 P_{e} \sum_{n} \frac{\sin n a \cos n x}{n a+s \cos ^{2} n a} \frac{\sinh [n(W-Z)]}{\sinh [n W]}$,
and (19) is obtained by using this hole density in (37) and (74) and thence in (3).

[^52]

# Further Analysis of Transmission-Line Directional Couplers* 

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#### Abstract

Summary-The conditions of infinite directivity for transmissionline directional couplers are derived in a general and rigorous way. Exact expressions valid for any degree of coupling are found.

The case of small coupling and the case of a matched primary line are considered as particular cases of this general analysis. In the case of small coupling, the condition $Z_{11} Z_{44}=1$ and $Z_{22} Z_{33}=1$ given by W. L. Firestone in a recent paper are found to be correct; it is shown also that these simple conditions do not apply anymore when the coupling becomes large.


IIN HIS PAl'ER, ${ }^{1}$ Firestone derived a condition of infinite directivity for transmission-line directional couplers. He expressed this condition in the notation defined further on as follows:

$$
\begin{equation*}
Z_{22} Z_{33}=1 \quad \text { and } \quad Z_{11} Z_{44}=1 \tag{1}
\end{equation*}
$$

It can be shown that conditions (1) are not the most general conditions for infinite directivity, but rather they apply only for the special case of weak coupling. The general conditions for infinite directivity will be derived here.

A rigorous amalysis, based clirectly on Maxwell's equations, of a system of $n$-parallel cylindrical conductors of arbitrary cross section, was published by the author (in joint. authorship) several years ago. ${ }^{2}$ As an application of this analysis, the case of a system of three conductors was investigated. (Such a system may consist of two coupled single wire transmission lines above a common ground or within a common cylindrical ground conductor.) It was found at that time (as Firestone has shown by other methods in his paper) that such a system has the properties of a directional coupler, if properly terminated. We shall base our present derivations upon this analysis.

While Firestone first analyzes a system of four parallel conductors and treats the system of three conductors as a particular case of the system of four conductors, for the sake of simplicity we shall limit our present discussion to the system of three conductors; in doing so we may bear in mind that a similar discussion is also possible for a system of four conductors. Further, let us refer from now on to Firestone's paper as (I) and to our paper as (II).

Let us consider the system shown in Fig. 1 with the notation indicated on this figure (a practical design of such a system for vhf or uhf is shown in Fig. 2). This system may be considered as the junction of four trans-

[^53]mission lines coupled over the length $l$ by the two transmission lines $A$ and $B$, the lines $A$ and $B$ consisting respectively of the conductors $a-0$ and $b-0$.

Let $z$ be the co-ordinate in the direction parallel to the transmission lines $A$ and $B$. I et be $z=0$ at the terminals (1) and (3) and $z=l$ at the terminals (2) and (4), the lines $A$ and $B$ being coupled only in the interval ( $0, l$ ).

$l=$ Distance between terminals (1)-(2) and (3)-(4)
$Z_{L_{1}}$ to $Z_{L_{4}}=$ Impedances seen at terminals (1) to (4)
Conductors $(a)$ and $(o)=$ Transmission line $A$
Conductors $(b)$ and $(o)=$ Transmission line $B$
Conductor (o) = common ground or common cylindrical envelope for lines $A$ and $B$
$C_{a 0}=$ Capacity per unit length between conductors (a) and (o)
$C_{b_{0}}=$ Capacity per unit length between conductors (b) and (o)
$C_{a b}=$ Capacity per unit length between conductors (a) and (b)
( $C_{a b}=$ coupling capacity)
Fig. 1-Coupled transmission lines.
Considering the propagation of 'TEM waves only, we have at any given plane perpendicular to the conductors ( $z=$ constant) uniquely defined voltages $V_{a}$ on line $A$ between the conductors $(a)-(0)$, and $V_{b}$ on line $B$ between the conductors $(b)-(0)$.


Fig. 2-Transmission-line directional coupler for uhf.
According to (II) [matrix equation (15)] $V_{a}$ and $V_{b}$ can be expressed as a sum of forward and backward traveling waves:

$$
\left.\begin{array}{l}
V_{a}=V_{a}^{\prime} e^{j(\omega t-\beta z)}+V_{a}^{\prime \prime} e^{j(\omega t+\beta z)}  \tag{2}\\
V_{b}=V_{b}^{\prime} e^{j(\omega t-\beta z)}+V_{b}^{\prime \prime} e^{j(\omega t+\beta z)}
\end{array}\right\}
$$

where

$$
V_{a}^{\prime}, V_{a}^{\prime \prime}, V_{b}^{\prime} \text { and } V_{b}^{\prime \prime}
$$

are constants and can be complex.

Let us define the complex reflection coefficients $\Gamma_{a 1}$ at terminals (1) of line $A$ and $\Gamma_{b s}$ at terminals (3) of line $B$ as follows:

$$
\left.\begin{array}{l}
\Gamma_{a 1}=\frac{V_{a}^{\prime \prime}}{V_{a}^{\prime}} \\
\Gamma_{b 3}=\frac{V_{b}^{\prime \prime}}{V_{b}^{\prime}} \tag{3}
\end{array}\right\} .
$$

According to (2), (3) and with the notation of Fig. 1, the voltages $V_{3}$ and $V_{4}$ at the terminals (3) and (4) may be expressed as follows:
$V_{s}=\left(V_{b}{ }^{\prime}+V_{b}{ }^{\prime \prime}\right) \cdot e^{j \omega t}=\left(1+\Gamma_{b z}\right) V_{b}{ }^{\prime} e^{j \omega t}$
$V_{4}=\left(V_{b}^{\prime} e^{-j \beta l}+V_{b^{\prime \prime}} e^{+i \beta l}\right) e^{j \omega t}=\left(e^{-i \beta l}+\Gamma_{b 3} e^{+i \beta l}\right) V_{b}^{\prime} e^{i \omega t}$.
Let us define the directivity $D_{1}$ of the system, when it is excited at terminals (1), by:

$$
\begin{equation*}
D_{1}=20 \log \left|\frac{V_{3}}{V_{4}}\right| . \tag{4}
\end{equation*}
$$

The expressions for $V_{3}$ and $V_{4}$ show that, according to (4)

$$
\begin{equation*}
D_{1}=20 \log \left|\frac{1+\Gamma_{b s}}{e^{-j \beta l}+\Gamma_{b s e^{+}}+j B l}\right| . \tag{5}
\end{equation*}
$$

This last relation yields for the condition of infinite directivity:

$$
\begin{equation*}
\Gamma_{b 8}=-e^{-9 ; \beta} . \tag{6}
\end{equation*}
$$

We can derive from (18) of (II) a general and rigorous expression for $\mathrm{I}_{\mathrm{bs}}$. In order to do so let the following quantities be defined in agreement with the notation of Fig. 1:

$$
Z_{\mathrm{a} 0}=\frac{\sqrt{\epsilon \mu}}{C_{0 \mathrm{a}}+C_{\mathrm{ab}}}=\text { characteristic impedance of line } A
$$

$$
Z_{0 b}=\frac{\sqrt{\epsilon \mu}}{C_{b 0}+C_{a b}}=\text { characteristic impedance of line } B
$$

(by introducing the coefficient $1 / \sqrt{\epsilon \mu}=v=p$ hase velocity of the TEM waves, the inductivities per unit length of the conductors considered are eliminated from our expression).

Let us further define, according to the notation of Fig. 1 and in agreement with the notation used by Firestone:

$$
\begin{aligned}
& Z_{1 L}=\text { impedance seen at terminals (1) } \\
& Z_{2 L}=\text { impedance seen at terminals (2) } \\
& Z_{3 L}=\text { impedance seen at terminals (3) } \\
& Z_{4 L}=\text { impedance seen at terminals (4) } \\
& Z_{11}=\frac{Z_{1 L}}{Z_{0 a}} \quad Z_{22}=\frac{Z_{2 L}}{Z_{0 a}} \\
& Z_{33}=\frac{Z_{8 L}}{Z_{06}} \quad Z_{44}=\frac{Z_{4 L}}{Z_{06}} .
\end{aligned}
$$

Using the notation just defined we obtain from (18) of (II) the following expression for $\Gamma_{b 3}$ :
$\mathrm{r}_{b s}=$
$\frac{\left(1+Z_{33}\right)\left(\mathrm{\Gamma}_{a 1} e^{j \beta l}-e^{-j \beta l}\right) Z_{44}-\left(1-Z_{44}\right)\left(1-\mathrm{\Gamma}_{a 1}\right) Z_{33 e^{-}}{ }^{-j \beta l}}{\left(1+Z_{44}\right)\left(1-\Gamma_{a 1}\right) Z_{33} 8^{i \beta l}-\left(1-Z_{33}\right)\left(\mathrm{\Gamma}_{a 1} 1^{i \beta l}-e^{-i \beta l}\right) Z_{44}}$.
(It will be noticed that when $Z_{33}=Z_{44}$, (7) becomes identical with (22) of (II), if the proper change of notation is performed. This had also to be expected.)

Introducing (7) into (6) one obtains for the condition of infinite directivity the following expression:

$$
\begin{equation*}
\mathrm{I}_{a_{1}} e^{2 ; \beta l}=-\frac{Z_{33}-1}{Z_{33}+1} . \tag{8}
\end{equation*}
$$

For some applications, it may be more convenient to express the condition of infinite directivity in terms of impedances rather than of reflection coefficients. For this purpose we express $\Gamma_{a 1}$ as a function of the normalized terminating impedances $Z_{11}$ to $Z_{44}$, and of a coefficient of coupling $K$ which will be defined.

For $\Gamma_{a 1}$ we have, from (28) of (II):

$$
\begin{equation*}
\mathrm{r}_{a 1}=\frac{\mathrm{\Gamma}_{10}-K^{2} \frac{Z_{22}}{1+Z_{22}} e^{-2 ; \beta l}}{1-\frac{K^{2}}{2} \cdot \frac{Z_{22}}{1+Z_{22}}\left(1+e^{-2, \beta l}\right)} \tag{9}
\end{equation*}
$$

with

$$
\begin{gather*}
\mathrm{V}_{10}=\frac{Z_{22}-1}{Z_{22}+1} e^{-3 ; \beta l}  \tag{9a}\\
K=\frac{C_{a b}}{\sqrt{\left(C_{a 0}+C_{a b}\right)\left(C_{b 0}+C_{a b}\right)}} \leqq 1 . \tag{9D}
\end{gather*}
$$

The reflection coefficient $T_{10}$ would exist on line $A$ at terminal (1) in the case of zero-coupling ( $C_{a b} \rightarrow 0$ ).

By means of (8) and (9) the condition of infinite directivity can finally be expressed as follows:

$$
\begin{equation*}
Z_{2 s}=\frac{\frac{1}{Z_{22}}+\frac{K^{2}}{4}\left(1-e^{-2 ; s l}\right)}{1-\frac{K^{2}}{4}\left(3+e^{-2 ; \beta l}\right)} \tag{8a}
\end{equation*}
$$

Eq. (8a) is equivalent to (8).
If we excite the system of Fig. 1 at the terminals (2) and define the directivity $D_{2}$ by:

$$
D_{2}=20 \log \left|\frac{V_{4}}{V_{3}}\right|
$$

the condition for infinite directivity $D_{2}$ is found, in analogy to (8a) to be:

$$
\begin{equation*}
Z_{44}=\frac{\frac{1}{Z_{11}}+\frac{K^{2}}{4}\left(1-e^{-2 ; j g}\right)}{1-\frac{K^{2}}{4}\left(3+e^{-2 ; \beta l}\right)} . \tag{8b}
\end{equation*}
$$

The relations (8a) and (8b) show that infinite directivity can be obtained when the primary line (line $A$ in our case) is mismatched as well as when it is matched. (We call line $A$ "matched" when $\Gamma_{a 1}=0$.) Iaving found the general condition for infinite directivity $[(8 a)$ and (81)) ] let us consider some particular cases of practical interest.

## Case df Small Coupling

Let us define "small coupling" by the condition:

$$
\begin{equation*}
K^{2} \ll 1 \tag{10}
\end{equation*}
$$

In this case, the conditions of infinite directivity given by (8a) and (8b) become:

$$
\begin{align*}
& \lim Z_{22} Z_{83}=1  \tag{11}\\
& \lim Z_{11} Z_{44}=1
\end{align*} \quad \text { for } \quad K^{2} \rightarrow 0
$$

Comparing (11) (which has been rigorously derived) with the conditions $Z_{22} Z_{33}=1$ and $Z_{11} Z_{44}=1$ given by Firestone for infinite directivity, we see that these last conditions hold only in the case of weak coupling. Although the case of weak coupling is of considerable practical importance [e.g. when a transmission-line directional coupler is tised for the measurement of swr, in the way suggested in (II)], cases of strong coupling may also be important; in such cases, the conditions (11) for infinite directivity are not valid any more, and (8a) and (8b) have to be used instead.

## Case of Matched Lines

Let us determine under which conditions the reflection coefficient on line $A$ is zero and the directivity of the system is infinite.

From (9) we find the condition for $\Gamma_{a 1}=0$ to be:

$$
\begin{equation*}
Z_{22}=\frac{1}{1-K^{2}} \tag{12}
\end{equation*}
$$

From (8) we find the condition of infinite directivity $\left(I_{1}=\infty\right)$ when $\Gamma_{a 1}=0$ to be:

$$
\begin{equation*}
Z_{83}=1 \tag{13}
\end{equation*}
$$

One notices in (12) how the coupling $K$ affects the matching condition of line $A$.

As a conclusion of this analysis we may summarize our results as follows:

1. Infinite directivity may le obtained with trans-mission-line directional couplers for mismatched as well as for matched lines. The condition of infinite directivity in its general form is given by the relations (8a) and (8b).
2. In the case of weak coupling, the conditions of infinite directivity reduce to the conditions $Z_{22} Z_{33}=1$ and $Z_{11} Z_{44}=1$.
3. In the case of matched lines and infinite directivity, ( $\Gamma_{a 1}=0$ and $D_{1}=\infty$ ) the matching impedance of the primary line is a function of the coupling coefficient [cf. (13a)].

# Phase Stabilization of Microwave Oscillators* 

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#### Abstract

Summary-A circuit has been developed with which microwave oscillators may be phase-locked to weak but stable reference signals. The circuit was operated with S-band oscillators (707B klystron; 2 C 37 triode oscillator) and a 2 K 50 K -band klystron. It is possible to lock a microwave oscillator directly or through a cascade of such circuits to a quartz-stabilized oscillator. The statistical theory of random noise is used to obtain an analysis of the stabilizing effect of the circuit, and the power spectrum of the stabilized microwave source is calculated. The scheme can also be applied in divider operation. Modifications are discussed. A modified circuit that uses carriersuppressed modulation of the reference signal has also been realized. In another circuit, the oscillator frequency is converted by means of a stable reference, and compared with a second reference that can be of low frequency and tunable. These latter circuits allow elimination of the excess noise introduced by crystal diodes. In the original straight dc circuit this noise cannot be eliminated, but calculation shows that its influence on the output power spectrum is very small.


[^54]
## Introduction

T1HIS P:AIER will discuss the phase stabilization of microwave oscillators. It should be clearly understood at the outset that phase stabilization is quite distinct from frequency stabilization in the conventional form. A frequency discriminator with a $\left(f_{0}-f\right)^{-1}$ control circuit would give essentially a phasestabilization type of control, but such a $\left(f_{0}-f\right)^{-1}$ control is neither realizable physically nor defined analytically for the operating region, i.c., for $f=f_{0}$. If instead of a frequency discriminator a phase discriminator is used at the outset, all necessary components are realizable.

Note also that frequency stabilization allows one to establish a frequency to an accuracy which is constant with time. Phase stabilization establishes a mean frequency with an accuracy directly proportional to the locking time. The interest then in phase stabilizing microwave oscillators is to realize the transference of frequency stability from one frequency region to another with any desired precision.

Recent developments in the techniques of molecular beam measurements and microwave spectroscopy ${ }^{-1}$ make it possible to olserve substances in a state where they absorb electromagnetic energy at one or several extremely sharply defined microwave frequencies. A sul)stance in such a state is therefore analogous to a cavity of very high $Q$ ( $10^{7}$ or better) with a persistently accurate resonance frequency. In order to measure the center frequency of one of these resonances it is desirable to have microwave oscillators whose output power is as monochromatic as possible.

At lower frepuencies, an oscillator controlled ly a quartz crystal may be used to generate a signal with very high stability for a period of hours or days. This low frequency can be multiplied by means of vacuum tuhe or silicon diode multipliers. Conventional multiplication usually yields a high-frecpuency spectrum that is not monochromatic but has sidebands, arising from lower frequency modulations, that remain because of the finite selectivity of the circuits. Furthermore, since multipliers with a gain of less than one (silicon diodes, for example) introduce additional noise into the spectrum, it is not desirable to multiply a frequency by more than a factor of 10 in these diodes.

In this paper we describe a stabilization circuit that allows a microwave oscillator to be locked to a harmonic of a stable reference oscillator. Through iteration of this process, the stability of a quartz-controlled oscillator can, essentially, be transferred to a K-band oscillator $(23,040 \mathrm{mc})$. Description of the experiment is followed by an analysis of the stability and performance of a phase-locking circuit.

## Phase Stablitzation of a K-Band Oscillator

The circuit that has been successfully used to stabilize the frequency of a 2 K 50 klystron is shown in lig. 1. The circuit consists, essentially, of a single, absolutely stable feedlack loop. Any phase modulation in the klystron is detected in the phase discriminator: the resulting signal is amplified in the differential de amplifier and applied to the repeller of the klystron to produce an opposite and staliilizing phase modulation.


Fig. 1-13lock diagram of phase-stabilizing circuit for K-band oscillator,

## Phase Discriminator

The reference signal and oscillator are introduced through the noncoupling arms of a waveguide hybrid junction, or "magic tee." The signals that arrive at the detector crystals on the two remaining arms, 1 and 2,

[^55] duction," Phys. Rev., vol. 94, pp. 1393-1394; June, 1954.
$\epsilon_{S 1} \sin \left(\omega_{c} t+\eta_{1}\right)$ and $\epsilon_{S 2} \sin \left(\omega_{c} l+\eta_{2}\right)$ are shown by the vector diagram in Fig. 2.

These fields are the sum of a signal coming from the reference source, $\epsilon_{R 1} \sin \omega_{c} t$ and $\epsilon_{R 2} \sin \omega_{1} t$, and a signal coming from the oscillator, $\epsilon_{01}\left(\sin \omega_{c} t+\phi_{1}\right)$ and $\epsilon_{02}\left(\sin \omega_{c} t\right.$


ARM ,


ARM 2

Fig. 2-Vector diagram of signals in hybrid junction.
$+\phi_{2}$ ). From symmetry properties of the magic tee it is seen that $\phi_{1}=\phi_{2}+\pi$. Since detected power $P_{D}$ in a silicon diode as a function of input power $P_{i}$ is given by ${ }^{2}$
$P_{D}=S P_{i}^{2}\left(S=\right.$ conversion gain perwatt $\left.\approx 10^{4}(\text { watts })^{-1}\right)$
it follows that the detected output signal is:

$$
V_{j}=-2 \sqrt{R S P_{0 j} P_{R j}} \cos \phi_{j}+\sqrt{R S}\left(P_{0 j}+P_{R j}\right), j=1,2 .
$$

Here, $P_{0 j}, P_{R j}$ are the power of oscillator and reference signal in each arm; $V_{j}$ is the phase discriminator detected output of each of the crystals. It can be seen that if $P_{0}$ and $P_{R}$ are divided equally between arms 1 and 2 , then $V_{1}-V_{2}$ is independent of variations of $P_{0}$ and $P_{R}$ for $\phi_{j} \approx \pi / 2$. Amplitude modulation of the two signals is therefore proportional to the balanced-out control signal, and hence second order, being negligitle in the limit of small control signal. Since the insensitivity to amplitude variation allows discriminator to be operated at a high power level in spite of the small $P_{R}$, crystals can be operated in a region of good conversion gain. For $P_{0}=200 \mu$ watts, $R=100$ ohms, $P_{R}=8 \mu$ watts, we expect a differential output of $g_{\nu}=0.04$ volt per radian.

As indicated in Fig. 1, this output is amplified in a differential de amplifier. A cross-cotpled circuit ${ }^{3}$ was used for this purpose; the actual circuit is shown in Fig. 3. The circuit uses all readily availalple techniques to achieve stable, hum-free operation. The heaters of the 12AX7 tubes are fed in series from the negative power supply. A K-band spectrum analyzer, an oscilloscope, and headphones are used to tume the oscillator to the reference signal. Once a klystron is within about $1,000 \mathrm{cps}$ of the reference, it will phase-lock itself automatically. Since the 2 K 50 klystron is microphonic, good sound isolation is essential. The experiment was, therefore, carried out in an anechoic chamber. But, any good acoustic isolation for the klystron should be sufficient.

The reference signal was supplied from a very stable, cavity-tuned, planar triode S-band oscillator. This oscillator, in turn, was locked by an analogous circuit to the tenth harmonic of the output of the M.IT. frequency standard. ${ }^{4}$ The correction signal was applied to the plate of the oscillator triode.

For the S-band oscillator, a klystron also could be
${ }^{2}$ M. W. P. Strandberg, "Microwave Spectroscopy," Methuen land Co., London, Eng.; 1954.
${ }^{3}$ J. N. van Scoyoc and G. F. Warnke, "A d-c amplifier with crosscoupled input," Electronics, vol. 23, pp. 104-107; February, 1950.
${ }^{1}$ C. G. Montgomery, "Technique of Microwave Measurements," M.I.'T. Radiation Lab. Ser., McGraw Hill Book Co., Inc., New York, N. Y., vol. 11, pp. 347-375; 1947.


Fig. 3-Cross-coupled de amplifier.
used; the 707 IJ klystron was phase-locked with this same equipment to the M.1.T. frequency standard.

## Analysis of the Phase-Stabilizing Feedback Loop

The powr spectrum of the output of a klystron, or any conventional oscillator, is not a single sharp line. There are three reasons. First, the klystron puts out a noise band as broad as the pass band of the loaded plate cavity. This noise may be thought of as simple diode noise. ${ }^{2}$ Second, variations of the supply voltages on the klystron electrodes impress a frequency modulation on the carrier. Third, microphonic pickup also causes frequency modulation in the klystron through the relative physical motion of the frequency-determining elements of the escillator.

The frequency-morlulated output is written

$$
\begin{equation*}
S=S_{n} \sin \left\{\omega_{c} t+\int \omega(t) d t\right\} \tag{1}
\end{equation*}
$$

where $\omega(t)$ is the frec uency modulation, and $\zeta(t)=\int \omega(t) d t$ is the phase modulation. The frequency change produced by voltage variations on an oscillator electrode may be written as

$$
\begin{equation*}
\frac{\partial \omega}{\partial m}=\beta \text { radians } \sec ^{-1} / \mathrm{volt} \tag{2a}
\end{equation*}
$$

where $m(t)$ is the random part of electrode voltage. The acoustical pickup might be caused by a variation of the distance of the grids in the gap of the plate cavity or by the vibration of the repeller perpendicular to the tube axis. The effect of this motion may, in general, be expressed as

$$
\begin{equation*}
\frac{\partial \omega}{\partial d}=\epsilon \text { radians } \mathrm{sec}^{-1} / \mathrm{cm} \tag{2b}
\end{equation*}
$$

where $d(t)$ is a characteristic distance in the physical frequency determining circuit. Both $m(t)$ and $d(t)$ are assumed to represent ranclom noise having a normal amplitude probability distribution. Hence, $\omega(t)$ can be represented as ${ }^{5}$

$$
\begin{align*}
\omega(t) & =\sum_{n=1}^{N} c_{n} \cos \left(\Omega_{n} t-Z_{n}\right)  \tag{3a}\\
\overline{\omega^{2}(t)} & =\frac{1}{2} \sum_{n=1}^{N} c_{n}^{2}=\int_{b}^{\top} W^{\top}(\Omega) d \Omega \tag{3b}
\end{align*}
$$

where $W(\Omega)$ is the frequency modulation power in radians $\sec ^{-1}$, and $b$ and $r$, the lower and upper cut-off frequencies, will be discussed later [see (8a) and (12)]. We would like to know the power spectrum of $S$ with the modulation (3a). The Fourier spectrum of a carrier, frequency-modulated by several independent sine waves, has been calculated by Crosby. ${ }^{6}$ He found sidebands displaced by $\Omega_{n}$ from the carrier with amplitude $1 / 2\left(c_{n} / \Omega_{n}\right)$, and cross-modulation bands of higher order in $c_{n} / \Omega_{n}$. The energy of the modulated carrier is concentrated within either twice maximum frequency deviation or twice modulating frequency, whichever is greater.

Thus we have two different cases to consider, $\left(c_{n} / \Omega_{n}\right)>1$ and $\left(c_{n} / \Omega_{n}\right) \ll 1$. In the first case we expect to find the power distributed within a band of width $\left[\overline{\omega^{2}(t)}\right]^{1 / 2}$ around $\omega_{c}$. This assumption was verified in an experiment in which a noise voltage of $10^{5} \mathrm{cps}$ bandwidth and of known rms voltage was applied to the repeller of a klystron whose output was observed in a

[^56]spectrum analyzer. In the second case, the cross modulation was neglected, and we find
\[

$$
\begin{align*}
S= & S_{0}\left\{\sin \omega_{c} l+\sum_{1}^{N} \frac{c_{n}}{2 \Omega_{n}}\left\langle\sin \left[\left(\omega_{c}+\Omega_{n}\right) l-Z_{n}\right]\right.\right. \\
& \left.\left.-\sin \left[\left(\omega_{c}-\Omega_{n}\right) l+Z_{n}\right]\right\rangle\right\}, \tag{4}
\end{align*}
$$
\]

with the power spectrum

$$
\begin{equation*}
p(\Omega) d \Omega=\frac{1}{2} S_{0}^{2} \frac{W(\Omega)}{\Omega^{2}} d \Omega . \tag{5}
\end{equation*}
$$

We now wish to calculate the effect of the stabilizing feedbark loop on the power spectrum. In order to obtain the open loop gain, $\mu\left(\Omega_{n}\right)$, we consider the $n$th term in (4). 'This term is the result of a phase modulation, $\zeta\left(\Omega_{n}\right)$. This modulation, present in the output of the oscillator, gives rise to a voltage output, $G_{P} \cdot \zeta\left(\Omega_{n}\right)$ from the phase discriminator. This output in turn is amplified lyy a factor of $\mu_{D}$ in the de amplifier and converted according to (2a) into frequency modulation. The latter process can also be described as phase modulation, with a gain of ( $\beta / i \Omega$ ) radians/volt. Hence,

$$
\begin{equation*}
\mu(\Omega)=G_{P} \mu_{D} \beta \frac{1}{i \Omega}=\frac{r}{i \Omega} \quad \text { with } r \equiv G_{P} \mu_{D} \beta . \tag{6}
\end{equation*}
$$

If the feedback loop is now closed, the $n$th term will be reduced by a factor $[1 /(1-\mu(\Omega))]$, and the power spectrum of the stabilized oscillator, if $r$ is assumed to be real, is given by

$$
\begin{equation*}
p_{s}(\Omega) d \Omega=\frac{1}{2} S_{0}^{2} \frac{W(\Omega) d \Omega}{r^{2}+\Omega^{2}} \tag{7}
\end{equation*}
$$

where $p_{s}(\Omega)$ is the stabilized noise power in watts/radian $\sec ^{-1}$. With an estimate of $W(\Omega)$ and $r$, the order of magnitude of the residual noise left in the stabilized system may be calculated. From (6) and typical numbers $G_{P}=0.04$ volts $/ \mathrm{radian} ; \mu_{D}=2 ; \beta=10^{7}$ radians $\mathrm{sec}^{-1} / \mathrm{volt}$, $r$ may be computed as $r=8 \cdot 10^{5}$ radians sec ${ }^{-1}$.
'This leads to a stabilization cut-off frequency $\Omega_{c} / 2 \pi$, where $\left|\mu\left(\Omega_{c}\right)\right|=1$, of

$$
\begin{equation*}
\frac{\Omega_{c}}{2 \pi}=\nu_{c}=130 \mathrm{kc} . \tag{8a}
\end{equation*}
$$

We can estimate $W(\Omega)$ from (3b). If we assume it in a first approximation to be independent of $\Omega$ and use an empirical value of $\overline{\Omega^{2}(t)}=10^{8}\left(\text { radian } \mathrm{sec}^{-1}\right)^{2}$ for the unstabilized klystron, and assume that these deviations are the result of noise up to $\Omega_{c}$ we compute

$$
W(\Omega)=100 \text { radians } \sec ^{-1}, \text { for } \Omega<\Omega_{c}
$$

Hence

$$
\begin{equation*}
p_{s}(\Omega)=8 \cdot 10^{-11} P_{c} \text { in the pass band. } \tag{8b}
\end{equation*}
$$

We estimate ${ }^{2}$

$$
\begin{equation*}
p_{\text {diode }} \doteq 10^{-16} P_{c} . \tag{8c}
\end{equation*}
$$

l3y integrating (7) we find
$P_{s t}=100 P_{c} \int_{-}^{r} \frac{d S}{r^{2} \perp \Omega^{2}}=\frac{100 P_{c}}{r} \arctan 1 \doteq 10^{-4} P_{c}$ (80)
where $P_{c}$ is the carricr power in watts; $p_{\text {diode }}$ is the power density resulting from diode noise in the klystron in watts/radian $\mathrm{sec}^{-1}$; and $P_{s t}$ is the total noise output of the stabilized circuit, up to $\nu_{c}$, in watts.

The phase angle $\zeta$ of the stabilized signal still has a (iaussian amplitude prob)ability distribution.

$$
\begin{align*}
P(\zeta) & =\frac{1}{\left(2 \pi \zeta^{2}\right)^{1 / 2}} \exp \left(-\zeta^{t} / 2 \overline{\zeta^{2}}\right) \\
\overline{\zeta^{2}} & =\frac{P_{s t}}{P_{c}}=10^{-4}, \tag{9}
\end{align*}
$$

where $\left(\overline{\zeta^{2}}\right)^{1 / 2}$ is the rms value of the phase variation, computed for the stabilized system.

Following conclusions are suggested by our results:

1. Eqs. (81) and (8d) show that the noise power in the whole spectrum of the stabilized oscillator is far below the carrier. This justifies the assumption of the validity of (4).
2. Eqs. (8b) anci (8c) show that between $\omega_{i}-2 \pi \nu_{c}$ and $\omega_{i}+2 \pi \nu_{c}$ the noise power produced by the klystron diode noise is negligible, compared with the frequency-modulation noise of the stabilized output. Although an ordinary broadband amplitude-modulation detector is insensitive to frequency modulation, so the diode noise is dominant, the high- $Q$ experiments described in the introduction will detect the frequency modulation noise.
3. Eq. (8a) shows that the stabilization loop has a pass band of only 130 kc . The dc amplifier will show no phase shift in this band; therefore, the loop will have a phase shift of $\pi / 2$ and will be absolutely stable.
4. Eq. (9) shows that the rms phase shift is roughly 1 degree and that a shift bigger than $\pi / 4$ is quite impossible. Hence, the phase discriminator operates in the linear region of its discriminator characteristic.
5. The stabilization changes the klystron power spectrum that was originally spread over a finite width (of approximately 100 cps ) into an impulse function and a very low and broad noise band. This result rests on the assumption that the reference signal is monochromatic. In practice, the spectrum width of the stabilized oscillator will be reduced to the width of the reference signal.

## Modifications of the Stabilizing Circuit

However obvious it may seem, it does appear worthwhile to point out that this phase-locking circuitry may also be used to make a divider of particular use in the microwave region where no other kind exists. The operation of the circuit has been discussed from the point of view of locking a microwave oscillator to the harmonics of a lower frequency oscillator. However, the correction signal may also be applied to the lower frequency oscillator to transfer to it the stability of the high-frequency oscillator. This divider type of operation would be quite useful, i.e., for the general utilization of output of molecular microwave oscillator frequency standards. ${ }^{7,8}$

[^57]

Fig. 4-Stabilizing circuit using carrier-suppressed modulation.


Fig. 5-Stabilizing circuit using second reference.

Figs. 4 and 5 show two of the many modifications of the feedback loop. These two forms allow a discussion that is sufficiently general that it can be applied to variations of the basic circuit.

In Fig. 4 the hybrid junction is used to apply carriersuppressed modulation to the reference signal (in crystal $A$ ) and to combine it with part of the oscillator output. Detection of the resultant signal in crystal $B$ is followed, after amplification, by a second detection in a phase detector whose output is the stabilizing signal.

The whole stabilization problem may be transformed to any convenient frequency, of course, by converting the oscillator signal with a stable reference. This may be accomplished in the manner shown in Fig. 5. Here, the oscillator is converted by the reference to a frequency of $f$ cps. The converted signal is amplified and compared with a second reference in a phase detector whose output serves again as the stabilizing signal. The relative stability of the second reference can be worse than that of the first reference by the ratio of their frequencies. Thus the second reference may be obtained from a tunable source so that the frequency of the stable oscillator can be varied even if the first reference is fixed.

In the circuits of both Fig. 4 and Fig. 5 use is made of IF power coming from a detector crystal; in the original circuit de power coming from two detector crystals was used. This is of interest in connection with the problem of detector noise. If a crystal diode (1N26) rectifies a small signal $P_{c}$, a noise power density $p_{k}(\Omega)$ will appear in excess of the themal noise: ${ }^{2}$

$$
\begin{equation*}
p_{h}=\frac{10^{23} k T P_{c}^{2}}{\Omega} \text { watts/radian } \mathrm{sec}^{-1} \tag{10}
\end{equation*}
$$

If an IF carrier is generated in the diode. a similar noise power spectrum is found, distributed around the carrier as it was before around the dc carrier, in accordance with (10).$^{9,10}$ Whether this power results from both frequency modulation and amplitude modulation or from amplitude modulation alone is not, at present, decided. It would seem that the latter is more likely. If this is true we have a means o! minimizing the effect of the crystal ${ }^{*}$ noise on the stabilizing signal. This AM crystal noise

[^58]may be cancelled by using a phase detector that is insensitive to amplitude modulation; for example, one that is adjusted to work with zero output at equilibrium. If the oscillator is stabilized with the de circuit the influence of the crystal noise on the output power spectrum cannot be balanced out because it arises from two independent sources-the two detecting crystals. Since this appears offhand to be a serious fault of the dc stabilization, we conclude with a calculation of influence of crystal noise on stabilized oscillator spectrum.

A noise voltage $s_{N}(\Omega) d \Omega$ introduced at the discriminator (see Fig. 1) is reduced by $1 /(1-\mu(\Omega))$ by feedback, and appears as phase modulation $\mu_{D} \beta s_{N}(\Omega) d \Omega / i \Omega(1-\mu)$ at the output. The noise power

$$
\begin{equation*}
p_{N}(\Omega) d \Omega=R p_{k} P_{c} \frac{\mu_{D}^{2} \beta^{2}}{\Omega^{2}+r^{2}} d \Omega \tag{11}
\end{equation*}
$$

is calculated in a fashion similar to $p_{8}(\Omega)$ in (7). We find that $p_{N}(\Omega)$ exceeds $p_{s}(\Omega)$ only when

$$
\Omega<\frac{2 R 10^{16} k T P_{c}^{2} \mu_{D}{ }^{2} \beta^{2}}{W(\Omega)}=350 \text { radians } \mathrm{sec}^{-1}
$$

if we set $p_{N}(\Omega)=p_{s}(\Omega)$ and assume $P_{c}=10^{-4}$ watts. This means that only in a band of 400 cps around the stabilized signal the noise introduced by the crystal exceeds $p_{s}(\Omega)$. Furthermore, even 1 cps from the carrier, $p_{N}(\Omega)$ $<10^{-7} P_{c}$. The contribution of the total crystal noise to phase modulation noise $P_{N t}$ is $\int_{t}^{T} p_{N}(\Omega) d \Omega$. We get

$$
\begin{equation*}
P_{N t}=\frac{1}{2} 10^{18} k T R P_{c}^{3} \mu_{D}{ }^{2} \beta^{2} \frac{1}{r^{2}} \ln \left[\frac{1}{2}\left(1+\frac{r^{2}}{b^{2}}\right)\right] . \tag{12}
\end{equation*}
$$

For $b=10^{-4}$, the lowest frequency observable within an hour, we obtain $P_{N t}=6 \cdot 10^{-7} P_{c}$. Comparing (12) with (8d) it may be seen that $P_{N t}$ is far less than $P_{s t}$. Eq. (9) indicates that the probability of saturating the discriminator has not increased. This shows that the stability of the output signal is not significantly affected by crystal noise. The circuit of Fig. 4 has therefore no special advantage. Its realization has been tried with some success, but it is difficult to set all the adjust ments properly. The circuit of Fig. 5 makes the oscillator more flexible frequency-wise but it does require a broad band IF amplifier. The original dc circuit is found to be simple and quite effective.

## Correspondence

## Kompfner Dip Conditions*

Kompfner ${ }^{1}$ has described operation of a traveling-wave tube amplifier so that zero output occurs for nonzero input. 'This occurs for a particular value of beam voltage and current; measurement of these quantities enables exact calculation of circuit phase velocity and impedance. Kompfner presents calculations based on the assumptions of a long tube, and zero space charge.

Because of the wide use of Kompfner's technique, it appears justifable to generalize his theory through use of l'ierce's threcwave theory. ${ }^{2}$ This will remove the assumptions that the tube be long and that space charge be negligible, but leave the assumption $C \ll 1$. The notation of Pierce is used throughout.

TABLE I
Conditions for the Kompfner Dip $Q C=0$

| $L$ | $C N$ | $\left(\beta-\beta_{d}\right) l$ | $b$ | $d$ |
| :---: | :--- | :---: | :---: | :---: |
| 0 | .3141 | -3.0040 | -1.522 | 0 |
| 3.201 | .2931 | -2.9116 | -1.581 | .2 |
| 6.017 | .2755 | -2.8369 | -1.639 | .4 |
| 8.527 | .2603 | -2.7721 | -1.695 | .6 |
| 10.79 | .2471 | -2.7169 | -1.750 | .8 |
| 12.85 | .2354 | $-2.66) 7$ | -1.805 | 1.0 |
| $Q C=0.2$ |  |  |  |  |
|  |  |  |  |  |


| $Q C=0.2$ |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $L$ | $C N$ | $\left(\beta-\beta_{\varepsilon}\right) l$ | $b$ | $Q C / C N$ | $I I$ |
| 0 | .33663 | -3.1780 | -1.504 | 0.5947 | 1.890 |
| 3.391 | .3105 | -3.0492 | -1.563 | 0.0441 | 1.745 |
| 6.318 | .2893 | -2.949 | -1.620 | 0.6913 | 1.626 |
| 8.898 | .2716 | -2.8620 | -1.677 | 0.7364 | 1.526 |
| 11.20 | .2565 | -2.7916 | -1.732 | 0.7797 | 1.441 |
| 13.29 | .2434 | -2.7332 | -1.787 | 0.8217 | 1.368 |
| $d=0$ |  |  |  |  |  |
| $Q C / C N$ | $C N$ | $\left(\beta-\beta_{\theta}\right) l$ | $b$ | $Q C$ | $I I$ |
| 0 | .3141 | -3.004 | -1.522 | 0 | 0 |
| 0.5947 | .3363 | -3.178 | -1.504 | .2 | 1.890 |
| 0.7280 | .3434 | -3.239 | -1.501 | .25 | 2.158 |
| 1.2531 | .3990 | -3.843 | -1.333 | .5 | 3.545 |
| 1.3803 | .4347 | -4.438 | -1.625 | .6 | 4.231 |
| 1.6098 | .4659 | -5.368 | -1.834 | .75 | 5.070 |
| 2.0354 | .4913 | -6.396 | -2.072 | 1.00 | 6.174 |
| 2.3697 | .5275 | -7.549 | -2.278 | 1.25 | 7.410 |
| 2.7164 | .5522 | -8.672 | -2.499 | 1.5 | 8.500 |

$L$ is the total circuit loss in $\mathrm{db}, l$ is the physical length of the active part of the circuit and beam, and $H=2 \pi C N \sqrt{4 Q C}$.

Application of a small signal of voltage $V$ to the tube input sets up three waves of incremental propagation constants $\delta_{1}, \delta_{2}$, and $\delta_{3}$ which are the roots of (7.14) of Pierce (which contains a misprint), ${ }^{3}$

$$
\begin{equation*}
\delta^{2}=\frac{1}{(-b+j d+j \delta)}-4 Q C . \tag{1}
\end{equation*}
$$

The total voltage amplitudes of these three waves are given by (9.4) of Pierce ${ }^{4}$ with

* Received by the IRE, February 23. 1955.
${ }^{1} \mathrm{~K}$. Kompiner, "On the operation of the travelingwave tube at low level," Jour. Brit. IRE, vol. 10, pp. 283-289; August-September, 1950 .
${ }^{2}$ J. R. Pierce, "Traveling-IVave Tubes," D. Van Nostrand Co., New York, 1950.

2 lbid., p. 113.
4bid., p. 133.


Fig. 1-Plot of conditions for the Kompfer dip.


Fig. 2-1 lot of conditions for the Kompfner dip. Circuit loss $L=0$.
$i=v=0$, namely

$$
\begin{equation*}
V_{1}=\frac{V \delta_{1}{ }^{2}}{\left(\delta_{1}-\delta_{2}\right)\left(\delta_{1}-\delta_{3}\right)} \tag{2}
\end{equation*}
$$

and cyclical permutations. The total voltage at $z=0$ is equal to the circuit voltage, so $V(0)$ is the actual input circuit voltage; the circuit voltage at $z$ is related to the total voltage at $z$ by (7.17) of Pierce. ${ }^{6}$ With (1) this can be written

$$
\begin{equation*}
\frac{\boldsymbol{V}_{a l}}{\Gamma_{1}}=\frac{\delta_{l}{ }^{2}+4 Q C}{\delta_{1}^{2}} \tag{3}
\end{equation*}
$$

and cyclical permutations. In (3) we have added to Pierce's notation numerical subscripts needed to distinguish the three component waves. Thus, the total circuit voltage at $z$ is

$$
\begin{align*}
& \frac{V_{c 1}+V_{c 2}+V_{c 3}}{V^{\prime}} \exp (j 2 \pi . V) \\
&= \frac{\left.\left(\delta_{1}^{2}+4 Q C\right) \operatorname{ex1}\right)\left(2 \pi C N \delta_{1}\right)}{\left(\delta_{1}-\delta_{2}\right)\left(\delta_{1}-\delta_{3}\right)} \\
&+\frac{\left(\delta_{2}^{2}+4 Q C \cdot \exp \right)\left(2 \pi C . V \delta_{2}\right)}{\left(\delta_{2}-\delta_{3}\right)\left(\delta_{2}-\delta_{1}\right)} \\
&+\frac{\left(\delta_{3}^{2}+4 Q C\right) \exp \left(2 \pi\left(N \delta_{3}\right)\right.}{\left(\delta_{3}-\delta_{3}\right)\left(\delta_{3}-\delta_{2}\right)} . \tag{4}
\end{align*}
$$

Now using I'ierce's (15) of appendix 7, ${ }^{6}$ namely

$$
\begin{equation*}
d=0.0183 / / / C N \tag{5}
\end{equation*}
$$

we can compute the seal roots $C N$ and $b$ of the complex equation

$$
\begin{equation*}
\frac{V_{c 1}+V_{c 2}+V_{c 3}}{V}=0, \tag{6}
\end{equation*}
$$

where $L$ and $Q C$ are parameters. The ront of lowest $C N$ is tabulated in Table 1, and plotted in graphical form in Figs. 1 and 2.
H. R. Jomnson
llughes Res. and Dev. Labs. Hughes Aircraft Co.
Culver City, Calif.

- Ibid. p. 255.


## Correction

J. II. Crysdale, one of the authors of the discussion on "Large Reduction of VHF Transmission Looss and Fading by the I'resence of a Mountain Obstacle in Beyond-Line-of-Sight l'aths," which appeared on pages 627-628 of the May, 1955 issue of the l'roceedings of 1ue IRE, has brought the following correction to the attention of the editors.

The phase of the second term in eq. (1) is

$$
\left(f_{2}+\psi_{1}-\phi_{1}\right)
$$

and not

$$
\left(f_{2}-\psi_{1}-\phi_{1}\right) .
$$

The figures accompanying the discussion have been interchanged.

The multiplying factor in footnote 2 is

$$
\exp \left(-j \xi_{1}\right)
$$

## Frequency Stable LC Oscillators*

I would like to comment on the above paper by J. K. Clapp. ${ }^{1}$ Mr. Clapp admits that the linear analysis by Eison and others indicates that the series-tumed oscillator is no more stable than a high $C$ Colpitts oscillator having the same circuit $Q$ and the same impedances presented to the tube. He then brings up the effect of harmonic components and resorts to linear theory to attempt to prove that the series-tuned circuit is superior.

In this attempt he h.2s made two basic errors. (1) He substitutes an equivalent circuit (Fig. 4) which is a special case pertaining only to resonant operation. Because $R g=-R s$, only three operating conditions are possible: a) resonance where we have either zero voltage or infinite current, b) below resonant frequency where the series impedance is pure capacitive reactance and the current leads the voltage by 90 degrees, c) above resonant frequency where the series impedance is pure inductive reactance and the current lags the voltage by 90 degrees. (2) The generator phase angle $\phi$ represented by (30) has no meaning. The phase relation between the current through the generator and its terminal voltage is entirely determined by the load connected to the terminals. The phase relation between the current through a generator and its internally generated voltage is determined by the sum of the impedances of the generator and the load. Had Mr. Clapp included the load imperlance into his calculation he would have found that $d f / d \phi^{\circ}$ or Fig. 4 was a discontinous function.


## Branch 1 Branch 2



It should be clear from the above that (40) is not valid because it was not based on ralid assumptioni. An analysis of an equivalent circuit substituting a constant current generator for the tube will indicate that the function $d f / d \phi$ is the same for the Colpitts and for the series-tuned circuit for equal $Q$ 's and equal impedances presented to the tube. The calcuation of points on a

* Received by the IRE. August 26. 1954.
${ }^{1}$ Proc. IRE, vol. 42, pp. 1295-1300; August. 1954.

Nyquist diagram which follows shows that the two circuits are entirely equivalent at resonance, near resonance, and at the second harmonic of resonant frequency (Fig. 1).

$$
\begin{aligned}
& Z_{1}=R+j\left(X_{L}-\frac{X_{C}}{2}\right) \\
& Y_{1}=\frac{1}{R+j\left(X_{L}-\frac{X_{C}}{2}\right)} \\
& Y_{2}=\frac{1}{-j \frac{X_{C}}{2}} \\
& Y_{T}=Y_{1}+Y_{2}=\frac{1}{R+j\left(X_{L}-\frac{X_{C}}{2}\right)} \\
& -\frac{1}{\frac{X_{C}}{2}} \\
& Y_{T}=\frac{j \frac{X_{C}}{2}-\left[R+j\left(X_{L}-\frac{X_{C}}{2}\right)\right]}{j \frac{X_{C}}{2}\left[R+j\left(X_{L}-\frac{X_{C}}{2}\right)\right]} \\
& Z_{T}=-\frac{j \frac{X_{C}}{2}\left[R+j\left(X_{L}-\frac{X_{C}}{2}\right)\right]}{R+j\left(X_{L}-\mathbf{X}_{C}\right)} \\
& i=-E i \mathrm{gm} \\
& E_{T}=i_{T} \\
& \text { Eigm } j \frac{X_{C}}{2}\left[R+j\left(X_{L}-\frac{X_{C}}{2}\right)\right] \\
& R+j\left(X_{L}-X_{C}\right) \\
& I_{1}=E_{T} I_{1} \\
& \frac{\operatorname{Eigmj} \frac{X_{C}}{2}\left[K+j\left(X_{L}-\frac{X_{C}}{2}\right)\right]}{R+j\left(X_{L}-X_{C}\right)} \\
& \frac{1}{R+j\left(X_{L}-\frac{X_{C}}{2}\right)} \\
& E i \operatorname{gnc} j \frac{X_{C}}{2} \\
& I_{1}=\frac{2}{R+j\left(X_{L}-\mathbf{X}_{C}\right)} \\
& E_{0}=-I_{1} j \frac{X_{C}}{2}=\frac{E i \operatorname{gm} \frac{X_{C^{2}}}{4}}{R+j\left(X_{L}-X_{C}\right)} \\
& =\frac{E i \mathrm{gm} X_{C}{ }^{2}}{4\left[R+j\left(X_{L_{H}}-X_{C}\right)\right]} \\
& u \beta=\frac{E_{0}}{E i}=\frac{\mathrm{gm} \mathbf{V}_{C}^{2}}{4\left[R+j\left(R_{L}-X_{C}\right)\right]} .
\end{aligned}
$$

For a practical case let $g m=10^{-3}$ at

$$
\begin{aligned}
& \qquad \omega_{0}+X_{C}=20 \Omega, \quad Y_{L}=20 \Omega, \quad R=.1 \Omega \\
& \text { at } \\
& \text { uß } \quad \frac{10^{-3} \times 400}{4 \times .1}=1 \\
& \omega_{0}+\frac{1}{10} \%, \quad X_{C}=19.98 \Omega, \quad Y_{L}=20.02!, \quad R=: 1 \Omega \\
& \qquad \begin{array}{l}
u \beta=\frac{10^{-3} \times 400}{4(.1+j .04)}=\frac{.4}{.4+j .16} \\
=
\end{array}
\end{aligned}
$$

at

$$
\begin{gathered}
2 \omega_{0}, \quad X_{L}=10 \Omega, \quad X_{L}=40 \Omega, \quad R=.1 \Omega 2 \\
u \beta=\frac{10^{-3} \times 100}{4(.1+j 30)}=\frac{.1}{.4+j 120}=\frac{1}{4+j 1200}
\end{gathered}
$$

by a similar development for the series tuned oscillator (Fig. 2)

$$
u \beta=\frac{g n ı X_{C}^{2}}{4\left[R+j\left(X_{L}-X_{C V}-X_{C}\right)\right]}
$$



Fig. 2.
I.et

$$
\mathrm{gm}=10^{-3}
$$

at
$\omega_{0}, X_{L}=500 \mathrm{~s} 2, X_{C V}=400 \Omega 2, X_{C}=100 \mathrm{~s} 2, R=2.5 \Omega$

$$
u \beta=\frac{10^{-3} \times 10^{4}}{4(2.5)}-1
$$

at

$$
\begin{gathered}
\omega_{0}=x_{0}^{2} \%, \quad X_{C}=99.9 \Omega, \quad X_{C V}=399.6 \Omega \\
X_{L}=500.5 \Omega, \quad R=2.5 \Omega \\
u \beta=\frac{10^{-3} \times 10^{4}}{10+j 4}=\frac{100-j 40}{116} \\
=.862-j .344
\end{gathered}
$$

at

$$
\begin{gathered}
2 \omega_{0}, \quad X_{C}=50 \Omega, \quad X_{C V}=200 \Omega \\
X_{L}=1000 \Omega, \quad R=2.5 \Omega \\
u \beta=\frac{10^{-3} \times 2500}{4(2.5+j 750)}=\frac{2.5}{10+j 3000} \\
=\frac{1}{4+j 1200}
\end{gathered}
$$

This should end all contention that the series-tuned circuit has any electrical advantage over any other, so that any choice of oscillator circuits may be made on the basis of practical advantages. On this basis it would seen that a circuit which is useful over a $2: 1$ or $2.5: 1$ frequency ratio would in most cases be preferable to one which is useful over a $1.2: 1$ range.

Actual experiments have borne out the above theoretically derived proof that the series-tuned oscillator is basically no more stable than other configurations.

The opinions herein are those of the writer and are not to be construed as representing the views of the Navy Department of the Naval Service at large.
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## Rebuttal ${ }^{2}$

Commander Bernard questions the equivalent circuit of Fig. 4 of this paper and states that it is applicable as a "special case pertaining only to resonant operation." Frankly, that is exactly what it is intended to represent. Apparently, he loses sight of the significance of the generator resistance, $R g=R_{s}$, as representing a voltage rise, which balances the drop in the circuit-series resistance, $R_{\mathrm{s}}$, at resonance and at the stable current of operation. The voltage is certainly not zero, and the current is certainly not infuite. The only possible mode of operation is at series resonance, hence discussion of operation above or below resonance is misdirected. The entire development was based on circuit as shown, i.e. with load connected.

The paper states, and it is reiterated, that the author "admits that the seriestuned oscillator is no more stable than a high-C Colpitts having the same circuit Q and the same impedances presented to the lube," and a mathematical development is presented to prove the point again. This point is but of minor academic interest.

If the references of the paper are consulted, particularly reference 8 to Prof. Edson's book, it would be appreciated that over a wide range of frequencies up to some tens of megacycles, in a high- $C$ circuit, presenting the same impedances to the tube, it is difficult or impossible to realize the same $Q$. On this practical consideration, the series-tuned circuit will always have higher frequency stability than high- $C$ Colpitts.

A further practical consideration is that in a low-impedance high- $C$ circuit the tuning range is severely restricted, if usually available components of small volume and cost are employed.

As Prof. Edson sums it up: "In the Colpitts oscillator, the reactances resuired for optimum stability are often impracticably small; and an attempt to realize the calculated values is frustrated by poor values of Q, impracticably large variable condensers, and other similar limitations."

One of the main objects of the author's paper was to present circuits having frequency ratios greater than 1.2:1 (the ratio of the clrcuit of the first paper ${ }^{3}$ ) while maintaining the advantages of the series-tuned ozcillator. Since the paper describes circuits of up to 2.5:1 frequency ratio, the objection to the series-tuned oscillator on the basis of frequency ratio cannot be sustained.

Commander Bernard's experiments to prove the series-tuned oscillator no more stable than other oscillators must have been conducted at frequencies high enough to make the advantages of the series-tumed oscillator hegin to deteriorate. It would be interesting to know how equality of the $Q$ 's and equality of impedance levels presented to the tube were established. The experiences of others, as given by references 1 and 4 , indicate that improvements in stability of from 10 to 100 times over conventional circuits are readily obtained.
G. G. Gouriet's and the author's procedure in taking the phase shift of the fundamental, due to harmonic distortion, as
${ }^{2}$ Received by the IRE, January 25, 1955,
${ }^{3}$ J. K. Clapp. «An inductance-capacitance oscillator of unusual frequency stability, ${ }^{n}$ PROC. IRE, vol. 36. pp. 356-358; March, 1948.
equivalent to the nonlinear effect, was based on Llewellyn's work, referred to in reference 9. The method used by Gouriet, and followed by the author, for evaluating the effect of the reactive component, and the conclusion that the effect is reduced by increasing the inductance of the oscillator circuit are not correct. This problem will be considered in a later communication. More importantly, this work pointed out the deleterious effect of disturbances, such as unwanted feedback from amplifiers, which had not been appreciated previously.

The circuit in question is amenable in design alld adjustment and gives satisfactory performance in many fields, ranging from standard frequency crystal oscillators, through many varieties of laboratory and testing oscillators, to master oscillators for transmitters.
J. K. Clapp

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## Characteristic Impedance of AirSpaced Strip Transmission Line*

The authors of a recent paper ${ }^{1}$ describe an improved mathematical technique for determining the characteristic impedance of a TEM wave in an air-spaced transmission line of the form illustrated in Fig. 1. The method leads to results which differ considerably from those previously published, ${ }^{2}$ the latter being based on approximate solutions originated by Maxwell, Palmer and others. The difference is most marked for relatively high values of impedance (e.g. $Z_{0}>80 \mathrm{~s}$ ) .


Fig. 1-Cross section of air-spaced micrustrip line.
As part of a general program on strip-line techniques we have, at these laboratories, carried out a fairly comprehensive analysis of the properties of different strip-line geometrics with the aid of an electrolytic tank. The tank employed was of the double-layer type, ${ }^{3}$ thus ensuring a good approximation to an infinite medium in the transverse plane. Instrumentation difficulties, however, limited the over-all accuracy (in the case of impedance measurements) to an estimated $\pm 3$ per cent.

A few of these results, restricted to the region of interest, are presented in Fig. 2. The thickness of the comfuctors is fixed at

* Received by the IRE, February 28. 1955.
${ }^{1} \mathrm{~K}$. G. Black and T. J. Higgins, "Kigorous Dctermination of the P'arameters of Microstrip Transmission Lines." Symposium on Microwave Strip Circuits, Tufts College, Mass,; October 11-12, 1954.
${ }_{2}$ IF. Assadourian and E, Rimai, "Simplitied theory of microstrip transmission systems, " I'ROC. IRE, vol. 40, pp. 1651-1657; 1952.
"An electrolytic tank for the measurement of the steady-state response, transient response, and allied properties of networks," Jour. IEE, part 1, vol, 96 , p. 163; May, 1949.


Fig. 2-Curves of characteristic impedance for air-spaced microstriv line.
$t=b_{2} / 32$ throughout, but repetition of some of the measurements for $t=b_{2} / 64$ and $\ell=b_{2} / 96$ showed that within this range, at any rate, the impedance is largely independent of strip thickness. (The change in conductor loss is, however, appreciable.) It will be noted that they are in excellent agreement with the numerical results calculated by Black and Higgins, ${ }^{1}$ and hence confirm the inadequacy of the earlier formulations. The agreement would seem to be more than a coincidence and I therefore feel that the authors deserve considerable praise for their hard-won achievement.

The sandwich (or tri--plate) line has also been studied and good agreement has likewise been obtained with pultished results." In both cases considerable data has been collected in regard to conductor losses, capacitative coupling between adjacent strip conductors (considerably lower in the case of the tri-plate line), and the effect on attenuation of changes to the shape of the conductor edge (found to be negligibly small). It is hoped to iaclude much of this data in a fortheoming pablication.

In conclusion it should be stressed that Assadourian and Rimai never claimed anything other than an approximate solution. It has always been clear that a rigorous solution to the air-spaced line would le of rather academic interest only, since practical microstrip lines usually involve a dielectric supporting slab. In the latter case the boundary conditions require (as a minimum) that there exist a longitudinal component of the electric vector. Sciegienny and Schetzen, whose work is mentioned briefly, ${ }^{6}$ made some progress towards a rigorous solution, and, for an idealized type of slab supported line, identified a dominant mode of the Ell type. However, in discussions with them in 1953 I was led to understand that attempts to obtain an explicit solution had not been successful, and that numerical computations of wavelength, results of which are quoted, ${ }^{\text {b }}$ had proved very tedious.
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-S. B. Cohn, "Characteristic impedance of the shieldedastrip transmission linon TFANS. IRE, vol. MTT-2, pp. 52-57; July, 1954.
\&J. Sciegienny and M. Schetzen, "Theoretical Analysis of a Strip Transmisginn System " Paretical Quart. Prog. Rep. Res. Lab. Elec. MIT. April 15 . 1953.

## Reflection Coefficients of Irregular Terrain at $10 \mathrm{Cm}^{*}$

A recent paper on the above subject by K. Bullington ${ }^{1}$ reminds the present authors of some relevant datia accumulated while they were members of the Research Laboratories of Sperry Gyroscope Co. at Garden City, New York.

The work of Bullington was concerned with gross reflection of radio waves encountered in point-to-point relay transmission. Because the separation between points of transmission and reception is large, the reflection takes place from a large area of the carth, usually being quite inhomogeneous. In order to understand this and related problems better, we chose to study' reflections from sample regions, carefully selected to be homogeneous in character. It is believed that the behavior of reflections from large regious can best be understond by studying homogeneous samples of the various types of surfaces that are likely to be encountered. Our experiments were performed at a wavelength of 10 cr and inclucled the study of salt and fresh water, dry sand and soil, moist sand, dry soil with vegetation, and ice. The bulk of the data was accumulated in 1943, on various sites on Long Island, New York.

The method of measuring the reflection coefficient of a surface employed in this investigation is illustrated in Fig. 1. A suitable radio transmitter and a receiver were supported on portable towers above the surface to be studied as shown. The direct and reflected waves arrive at the receiver after traveling the paths indicated. The intensity of the field at the receiver depends upon the vector sum of the two waves. The relative phase of the direct and reflected waves may be changed by changing the height of the receiver. By observing the character of the resulting interference pattern as a function of height of the receiver, the strength of reflected waves can be determined in the usual way, making it possible to determine the magnitude of the refection coefficient.

(a)

(b)

Fig. 1 (a) Arraugement of anparatus in making measurements of reflection =vefficient. (b) Idealized urements of teflection suef
geometry of the experiment.

* Keceived by the IRE. January 17. 1955. ${ }^{1} \mathrm{~K}$. Bullington. "Reflection cuefficient of irregular terrain." Proc. IRE, vol. 42, pp. 1258-1262; August. 1954.

The symbols used in the following discussion are as follows: $d$ is the horizontal separation between the transmitter and the receiver, $r_{1}$ is the length of path of the direct wave, $r_{2}$ is the length of the path of the reflected wave, $h_{t}$ is the height of the transmitter above plain earth, $h_{r}$ is the height of the receiver, and $\psi$ a is the angle of inciclence of the reflected wave.

Using the above notation and assuming that $d \gg \lambda, h_{t}=h_{r} \gg \lambda$, and $r_{1}=d$, the reflection coefficient $R$ can be shown to be

$$
R=\frac{1}{\cos \psi_{2}} \frac{1+\frac{E_{\mathrm{anx}}}{E_{\mathrm{min}}}}{1-\frac{E_{\mathrm{mux}}}{E_{\mathrm{miz}}}}
$$

where $E_{m a x}$ and $E_{\text {min }}$ are the maximum and minimum values observed at the receiver as it is moved vertically through the interference pattern. In practice, one may not always be able to satisfy the approximations used in arriving at the above formula, and more elaborate calculations become essential. Moreover, the transmitting and receiving antennas are partially directional, thus changing the relative strength of the direct and reflected waves. All of these details were accounted for in our work by the expedient of using specially prepared nomographs. The reflection coefficients for the various types of earth are presented in Figs. 2 to 13. An attempt was made in each case (0) correlate these data with theory. Whenever possible, values of dielectric constant and conductivity were obtained from existing sources and appropriate curves were computed using the well-known methods. ${ }^{2}$


Fig. 2-Reflection coefficient for salt water. Merrick Canal. Surface: Sea water covered with estimated $t$ inch ripples. Points experimental; theoretical curves drawn for $\varepsilon=69, \sigma=6.5 \times 10^{-11}$ emu.


Fig. 3-Reflection coefficient for salt water. Short Beach. Surface: Tidal flat covered with 18 inches of retical curves drawn for $\epsilon=69, \sigma=6.5 \times 10^{-11}$ emu
${ }^{2}$ C. R. Burroughs. "Radio propagation over plain earth-field strength curves," Bell Sys. Tech. Jour vol. 16, pp. 45-75; January. 1937.
F. E. Terman, "Radio Engineers Handbook, Mchraw- Hill Book Co., Inr New York; pp. 700-707. eral Communication Commission, Document 47475.


Fig. 4-Reflection coefficient for fresh water, Kenyon Farm. Surface: Fresh water pond, smooth. Theoretical curves plotted for $\epsilon=80$.


Fig. 5-Reflection coefficient for dry sand. Oak Beach. Surface: Small dry sand hillocks, some vegetation. Theoretical curves drawn for $e=4$.


Fig. 6-Keflection coefficient for harrowed field. New York State Agricultural Institute. Surface: Harrowed field, clay-sand soil, lumps 1 to 2 inches in diameter. Theoretical curves drawn for $=4$.


Fig. 7-Reflection coefficient of a tidal flat at low tide, Short Beach. Surface: Tidal flat with some organic material. Theoretical curves drawn for $t=10$.


Fig. 8-Reflection coefficient for moist sand. Short Beach. Surface: Moist sand, some algae, very smooth. Theoretical curves drawn for $e=15$.


Fig. 9-Reflection coefficient of dry soil. Hicksville Airport. Surface: Rolling field, dry soil, grass 4 inclies long.


Fig. 10-Reflection coefficient of a grass covered field, Garden City, New York. Surface: Slightly rolling, grass 4 to 18 inches high, dry.


Fig. 11-Reflection coefficient of an agricultural field, New York State Agricultural Institute. Surface: Beet field covered with weeds.


Fig. 12-Reflection coefficient of brush covered terrain, New York State Agricultural Institute. Surface: Growth of pine trees 3 to 10 feet tall, bushes, weeds, and gravelly soil.


Fig. 13-Keflection coefficients of ice. Lake Tiorati. New York. Surface: Smooth ice, 22 inches thick Theoretical curves drawn for $=4$.

Too accurate an agreement cannot be expected as the dielectric constant and conductivity are not directly applicable to the heterogeneous material encountered in our experiments.

Our results show, as would be expected from optical reasoning, that earth, which is rough when compared to $\lambda / 2$, does not produce specular reflection and the reflection coefficient has no meaning. For example, plowed fields produce almost complete scattering. Earth covered with vertical grass may produce specular reflection for horizontal polarization, and seattering for vertical polarization. In the case of smooth and isotropic surfaces, the consentimal theory of reflection coefficient applies, provided that the dielectric constant and conductivity of the surface are known. By fitting curves to the experimental data presented above, values of dielectric constant and conductivity have been determined and appear to be in fair agreement with known or estimated values of these constants as obtained by other means.

We wish to acknowledge the help of Messrs. W. Frost, J. Singer, and M. Dickerson for help in accumulating the above information and to V. R. Learned for numerous suggestions. We also appreciate the encouragement and help of Dr. W. TT. Cooke, and the permission of Sperry Gyroscope Co. to publish this information.
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## Note on Helix Propagation*

It has often been stated that all space harmonic components of a helix mode have the same group velocity. ${ }^{1}$ By a tacit assumption that the familiar relation among power flow, energy storage, and group velocity is valid for the separate space harmonies, it has been implied that all space harmonics carry power in the same direction. This writer thinks that this implication is erroneous because the above relation is not a valid one for the separate space harmonics. The usual proof ${ }^{2}$ fails because these fields separately carry complex power into the cylindrical surface in which the helix lies. Helical-line cut-off propertie: can be clarified if one is not bound by the concept to which this writer takes exception.

Brief evidence in support of the writer's position is shown by examination of power flow conditions as the high frequency cutoff point of the delay-line mode (Sensiper's $\beta_{0}$ mode) is a pproached. Here the impedance (total axial power flow divided by suluare of current) drops rapidly and is finally zero at cutoff, where the field is a complete standing wave. It is proposed that the decrease in

[^59]total power flow is not the result of all powers in space harmonic fields decreasing, but rather by a subtraction of the power in the -1 space harmonic held from the aggregate power in the remaining fields. The -1 lield carries power in the backward direction, which is in the same direction as its phase velocity. As cutoff is approached, this field spreads radially as its phase velocity approaches the velocity of light; it becomes increasingly dominant and finally cancels the forward power to produce cutoff. This reasoning applies to both the shielded and unshielded helical lines and tends to support earlier evidence ${ }^{3.4}$ that curoff is not simply the result of a radiation condition, which is peculiar to the unshielded helix only.

> I. Stark

Res. and Dev. Labs.
Hughes Aircraft Company
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${ }^{3}$ I. Stark, "The lower morles: of a concentric line having a helical inner conductor "Jour. A ppl. Phys., vol. 25.pp. 1155-1162; September, 1954.

4 J. R. Pierce and P. K Tien. "Coupling of modes in helixes," Proc. IKE, vol. 42, pp. 1389-1396; September, 1954.

## Design Considerations of Junction Transistors at Higher Frequencies*

The lowest frequency for which the Tnet work obtained in this papert remains correct is governed by the validity of the approximation of the frequency variable $s_{p}$ :

$$
\begin{equation*}
s_{p}=\frac{W_{0}}{\left(I I_{p} \tau_{p}\right)^{1 / 2}}\left(1+j \omega \tau_{p}\right)^{1 / 2} \approx \frac{W_{0}}{\sqrt{\bar{D}_{p}}}(j \omega)^{1 / 2} . \tag{1}
\end{equation*}
$$

By using the exact expression above for $s_{p}$, oue can extend the applicability of the $T$-network down to zero frequency. Retaining the previous expressions for $z_{A}, z_{B}$, and $z_{c}$, in terms of $z_{a 0}$ and $z_{50}$, and using for the collector capacitor impedauce on a $2 b$-ohm impedance level

$$
\begin{equation*}
Z_{2}=\frac{2 b}{C_{c} s}=\frac{2}{K s_{p}^{2}-A}, \tag{2}
\end{equation*}
$$

where $s=j \omega, A=C_{c} / b \tau_{n}$, and $k=D_{p} C_{c} / b V^{r}{ }_{0}{ }^{2}$, one now obtains for the $\uparrow$-urm impedances $z_{A}, z_{B}$, and $z_{C}$,

$$
\begin{align*}
& z_{A}=\frac{2\left(K-\frac{A}{s_{p}^{2}}\right)+\frac{\tanh s_{p} / 2}{s_{p} / 2}}{\left(K s_{p}-\frac{A}{s_{p}}\right) \operatorname{coth} s_{p}+1},  \tag{3}\\
& \frac{\tanh s_{p} / 2}{s_{p} / 2}  \tag{4}\\
& z_{B B}=\frac{\left.K^{\prime} s_{p}-\frac{A}{s_{p}}\right) \operatorname{coth} s_{p}+1}{},
\end{align*}
$$

and

$$
z_{C}=\frac{\frac{2}{s_{p} \sinh s_{p}}}{\left(K s_{p}-\frac{A}{s_{p}}\right) \operatorname{coth} s_{p}+1}
$$

* Received by the IRE, March 16, 1955.

11. Statz, E. A. Guillemin, and K. A. Pucel Proc. IRE, vol. 42, pp. 162 m -1628; November 1954

The poles of these functions are obtained as the solutions of the equation

$$
\begin{equation*}
k s_{p}-\frac{A}{s_{p}}+\tanh s_{p}=0 ; \tag{6}
\end{equation*}
$$

or, with $s_{p}=j u$, from the equation

$$
\begin{equation*}
\tan u=-K u-\frac{A}{u} . \tag{i}
\end{equation*}
$$

For the typical values $C_{c}=100 \mu \mu$, $b=2 \times 10^{-5}, \tau_{p}=10^{-4}$, as given in the article, $A \approx 0.05$; thus, one may ignore the term $A / u$ in comparison with $K u$ and obtain the same solutions for $u$ :

$$
\begin{equation*}
u_{\nu}=\frac{\nu \pi}{2} \text { for } \nu=1,3,5, \cdots \tag{8}
\end{equation*}
$$

I'sing the same delinition of the frequency variable $p$,

$$
\begin{equation*}
p=\frac{4}{\pi^{2}} \frac{W_{0}^{2}}{D_{p}} s \approx \frac{s}{\omega_{\alpha}}, \tag{9}
\end{equation*}
$$

and introducing the definition

$$
\begin{equation*}
d=\frac{4}{\pi^{2}} \frac{\Pi_{0}^{\prime}{ }^{2}}{D_{p} \tau_{p}} \approx \frac{1}{\omega_{\alpha} \tau_{p}}, \tag{10}
\end{equation*}
$$

one obtains

$$
\begin{equation*}
s_{p}=\frac{\pi}{2} \sqrt{p+l}=j u . \tag{11}
\end{equation*}
$$

Hence the poles can be expressed in terms of $p$ as

$$
\begin{equation*}
p_{\nu}=-\nu^{2}-d \text { for } \nu=1,3,5, \cdots \text {. } \tag{12}
\end{equation*}
$$

However, since $d$ is of the order of $10^{-4}$, then

$$
\begin{equation*}
p_{p} \approx-\nu^{2} \text { for } \nu=1,3,5, \cdots \tag{13}
\end{equation*}
$$

which coincide with the previous results.
The residues in these poles, which may be obtained following the pattern outlined in the article, are:

$$
\begin{align*}
& \text { residues of } z_{A} \approx \frac{16}{\pi^{2}},  \tag{14}\\
& \text { residues of } z_{B} \approx \frac{32(-1)^{(\nu-1) / 2}}{K \pi^{3} \nu}, \tag{15}
\end{align*}
$$

$$
\begin{equation*}
\text { residues of } z_{C} \approx \frac{32(-1)^{(v+1) / 2}}{K \pi^{3} v} \tag{16}
\end{equation*}
$$

since $A / K \ll 1$ and

$$
\left|\frac{\tan \nu \pi / 4}{\nu \pi / 4}\right| \ll 2 K \text { for } \nu=1,3,5, \cdots
$$

These residues are the same as those obtained previously.

However, $z_{c}$ no longer has a pole at $s_{p}=0$ ( $n \mathrm{r} p=0$ ). For as $s_{p} \rightarrow 0$,

$$
\begin{aligned}
z_{C} & =\frac{2}{\left(K s_{p}^{2}-A\right) \cosh s_{p}+s_{p} \sinh s_{p}} \\
& \rightarrow \frac{2}{\left(K s_{p}^{2}-A\right)+s_{p}^{2}} .
\end{aligned}
$$

Thus, $z c$ is finite at the origin, but has a pole at

$$
\begin{equation*}
s_{p}^{2} \approx \frac{A}{K+1}, \tag{17}
\end{equation*}
$$

or, since

$$
s_{y}=\frac{\pi}{2} \sqrt{p+d} \quad \text { and } \quad A=\frac{\pi^{2}}{4} K d,
$$

this pole is also given by $p=-d / K$. For this pole,

$$
z_{C} \rightarrow \frac{8}{K \pi^{2}} \frac{1}{p+d / K},
$$

and its residue there is $8 / K \pi^{2}$, which is the same as that obtained in the article for the pole of $z_{C}$ at the origin.

This residue in conjunction with the new pole position will appear in the equivalent circuit as a resistor shunted across the series capacitor in the $z_{c}$-arm of the $T$-network, Fig. 9 of the article. The net result is that the new $T$-network will permit a base current to flow even at zero frequency-which is more in accordance with physical facts. This is the major change effected by the use of the exact expression for $s_{p}$. It can be shown that $z_{A}$ and $z_{B}$ also have simple poles at $p=-d / K$ with residues $-d / K$ and $+d / K$, respectively.

The infinite partial fraction expansions for $z_{A}, z_{B}$, and $z_{C}$ corresponding to those obtained in the article are:

$$
\begin{align*}
z_{A}= & \frac{16}{\pi^{2}}\left[\frac{-\pi^{2}}{16} \frac{d / K}{p+d / K}+\frac{1}{p+1}+\frac{1}{p+9}\right. \\
& \left.+\frac{1}{p+25}+\cdots\right]  \tag{18}\\
z_{B}= & \frac{32}{K \pi^{3}}\left[\frac{\pi^{3}}{32} \frac{d}{p+d / K}+\frac{1}{p+1}-\frac{1 / 3}{p+9}\right. \\
& \left.+\frac{1 / 5}{p+25}-\cdots\right]  \tag{19}\\
z_{C}= & \frac{32}{K \pi^{3}}\left[\frac{\pi / 4}{p+d / K}-\frac{1}{p+1}+\frac{1 / 3}{p+9}\right. \\
& \left.-\frac{1 / 5}{p+25}+\cdots\right] . \tag{20}
\end{align*}
$$

Note that for $d=0$ (which corresponds to approximating $s_{p}$ by $W_{0}(j \omega)^{1 / 2} / \sqrt{D_{p}}$ ), these expressions coincide with those obtained in the article. I'sing the exact expressions for $z_{A}, z_{B}$, and $z_{C}$,

$$
\begin{aligned}
& z_{A}(0) \approx 1 ; \quad z_{B}(0) \approx 1 \\
& z_{C}(0) \approx 8 / \pi^{2} d=2 D_{p} \tau_{\mathcal{p}} / W_{0}{ }^{2}
\end{aligned}
$$

for $0<d \lll 1$ and $K \gg 1$.
If the infinite expansions are terminated after a finite number of terms, the approximations will be improved if the approximate net effect of the abandoned terms is taken into account in the range of $p$-values considered.

For example, if terms up to and including the pole at $p=-25$ are retained, for the range $0 \leqq p \leqq j 6$, the terms dropped are essentially resistive and nearly equal to their respective values at $p=0$. Consequently, one may represent their combined effect in each series by a constant evaluated at the origin. These constants are the same as those obtained in the article; thus, the modified circuit will revert to that of the article when $d=0$.

The normalized equivalent circuit for the T-network corresponding to Fig. 9 of the article is shown here in Fig. 10. The portion enclosed by the dotted lines extends the applicability of the circuit obtained in the
article to radian frequencies below $1 / \tau_{p}$. This circuit may be denormalized by use of Table 1 of the article. ${ }^{1}$


Fig. 1-Explicit form of the T-network applicable to low frequencies (four terms in the expansion).

In conclusion it is informative to observe that this equivalent circuit describes the transistor correctly at zero frequency. For this purpose the current through the collector (which is considered to be shorted) is compared with the current through the base. From the circuit the ratio of base to collector current is simply equal to $z_{B} / z_{C}$ $\approx\left(\pi^{3} K / 32\right) /(\pi K / 4 d)=W_{0}^{2} / 2 D_{p} \tau_{p}$ since $K$ $\gg 1$ and $0<d \ll 1$.

1t is also very simple to calculate the ratio directly from the physical properties of the transistor. The emitter current is approximately equal to the collector current $I_{c} \approx I_{e}=A q D_{p} \operatorname{grad} p \approx A q D_{p} p_{o} / W_{0}$, where $A$ is the emitter area, $p$ is the concentration of holes, and $p_{0}$ is the concentration of holes at the emitter side of the base. The base current, which is equal to the total recombination current, is given by $\left(q / \tau_{p}\right) \int p d v$ $\left.\approx\left(A q p_{\mathrm{e}} W_{0}\right) / 2 \tau_{p}\right)$. The ratio of the emitter current to the base current is $W_{0}{ }^{2} /\left(2 D_{p} \tau_{p}\right)$, which agrees with the above result. The fact that the emitter efficiency is not equal to unity was not considered, but it is included in the complete circuit of Fig. 3 of the article.

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## "yrneh"*

The writer is indebted to Professor True Mclean for having called to his attention the true origin of the term "yrneh." 1 It appears that the term was originated in 1910 by the famed Professor Vladimir Karapetoff ${ }^{2}$ of Cornell University.

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* Received by the 1RE. April 8. 1955.
${ }^{1} \mathrm{H}$. Stockman. "On reciprocal inductance," 1'roc. IRE, vol. 43, p. 341; March, 1955 .
${ }^{2}$ 'V. Karapetoff, "The Magnetic Circuit." McGrawHill Book Co., New York. N. Y., p. 10; 1911.


## On Reciprocal Inductance*

I wish to concur with Mr. Baghdady ${ }^{1}$ in his suggestion of the term "inertiance" for reciprocal inductance. A further suggestion is the use of the term "erny" for the units of "inertiance" is a sort of tribute to Ernst Guillemin of M.I.T., who has done so much to popularize the use of reciprocal inductance in his courses on "Guillementary Circuit Analaysis."

## 11. T. McAleer

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* Received by the IRE, December 30. 1954.
${ }^{1}$ Proc. 1.R.E., D. 1807; December, 1954.


## Optimum Patterns for Endfire Arrays*

In a recent paper, Dullamel has described a method of synthesizing an equal-minor-lobe, or Tchebycheff, directivity pattern for an endfire linear array. This method is based on a generalization of the synthesis procedure suggested by l-I. J. Riblet ${ }^{2}$ in extending the work of C. L. Dolph ${ }^{8}$ for the broadside array to include the case of arrays having an element spacing less than a half wavelength. Dolph's original method yields the optimum pattern ouly for element spacings greater than a half wavelength, whereas Riblet's method permits an optimm pattern to be specified for any element spacing. However, Riblet's method is applicable as such to arrays having an odd mmber of elements only, whereas Dolph's method may be used for even or odd numbers of elements.

The purpose of this communication is to point out that the method originally described by Dolph may be applied directly to the case of the endfire array. An optimum pattern may be obtained for the endfire array for any element spacing, in contrast with the case of the broadside array. (It

[^60]should be noted, however, that element spacings greater than a half wavelength normally are not used in an endfire array in order to avoid extra major lobes in the pattern.4) There are two advantages of using Dolph's method. First, the same method may be applied directly to arrays having even or odd numbers of elements. Second, the resulting equations for determining the relative currents for the elements (excitation coefficients) are of somewhat simpler form than the corresponding equations derived by Dulfamel from Riblet's method.


Fig. 1-Graphical const ruction of optimum directivity pattern $S_{0}(\psi)$ froin fourth-degree Thebycheff nolynomial $T_{4}(z)$ for five-clement endfire array with $d / \lambda=\frac{1}{}$.

Application of Dolph's method to the endfire array is illustrated in Fig. 1 for the five-element array used by Dullamel (cf. his Fig. 5). Notation used here is the same as that in DulHanel's paper. 'The space factor $S_{5}(\psi)$ expressed as a function of the auxiliary variable $\psi$ is determined from the fourth-degree Tchel)ycheff polynomial $T_{4}(z)$ by means of a linear transformation $z=A$ $\cos (\psi / 2)$. (DulHamel determines his space factor from the second-degree polynomial $T_{s}(z)$ with $z=a \cos \psi+b$.) In either case, the physical directivity pattern, expressed as a function of $\phi$, is described by the space factor $S_{5}(\psi)$ over a limited (visible) range of $\psi$. As shown in Fig. 1, the unknown constants $A$ and $\alpha$ are determined by the two conditions that for $\phi=0,|S(\psi)|=R$, or $z=z_{0}$, while for $\phi=\pi|S(\psi)|=1$ or $z=-1$.

Further details of this method cannot be

- This was noted by Dullamel, loc. cit., p. 655.
included here because of space limitations. However, a more complete description, including numerical examples for the sevenelement array used by DuHamel and for a four-elcment array, plus a description of an alternative method of overdesigning a supergain antenna, has been submitted for publication to the Transactions of the IRE, l'rofessional Group on Antennas and Propagation. ${ }^{5}$ It also might be noted that the bidirectional array considered by DuHamel may be treated by a similar adaptation of Dolph's method. ${ }^{6}$
R. L. I'ritchard 21.53-A Daisy lane Schenectady, N. Y.

巨 R. L. Pritchard, "Discussion on optimuni patterns for endfire arrays," Trans. IRE, vol. AI'3. Dp, 40-43; January, 1955.
*R. L. Pritchard, "Optimum directivity patterns for linear point arrays," Jour. A cous. Soc. Amer., vol. 25 , pp. 890-891; September, 1953.

## The Unit for Frequency*

The Proceedings uses cycle, kc , and me as units of frequency, whereas the Radiation Laboratory Series, for example, uses cps, $\mathrm{kc} / \mathrm{sec}$, and $\mathrm{Mc} / \mathrm{sec}$. The conflict between the desire for convenience and the desire to retain the usual meaning of the word "cycle" is perhaps epitomized by the following sentence (Proc. IRE, vol. 42, p. 1372 ; Sept., 1954, top of the page). "The light is chopped by a 90 -cycle-per-second sector disc, and the ac photoresponse is measured by a 90 -cycle amplifier-detector."

A happy solntion would be more widespread use of the term "hertz," meaning cycle-per-second. Thus the units of frequency would be hertz (or hz), khz, and Mhz. A rate of sweeping frequency could have the unit Mhz $/ \mathrm{sec}$, ins.tead of $\mathrm{mc} / \mathrm{sec}$ or $\mathrm{Mc} / \mathrm{sec}^{2}$. Noise power density could have the unit watts/Mhz, instead of watts/me or watts/(Mc/sec).

The term hertz is listed in "Electronics Dictionary" by Cook and Markns, McGrawHill look Company, 1945; "German Military Dictionary," TM 30-506; "Antennas" by Kraus, McGraw-Hill Book Company, 1950; and "Science," December 24, 195.t. l'. W. Crist
Airborne Instruments Lab. Mineola, N. Y.

[^61]C. A. Adridge was born in Canandaigua, N. Y., on May 2, 1922. He was employed from 1940 to 1944 by the Consolidated Machine Tool Corp.

C. A. Alidridge and in cooperative programs studied mechanical engineering from 1940 to 1942 at the Rochester Institute of Technology, and electrical engineering from $19+2$ to 1944 at the Iniver. sity of Rochester. From 1944 to 1946 he was with the Navy at the Naval Research Jaboratories and in the South Pacific.

Mr. Mdridge received his B.S. degree in physics in Junc, 1950 from Syracuse University. He completed a graduate year in physics and then joined the Electronics Laboratory in June, 1951.

Since then, he has been engaged in development work in the fields of colored television studio equipment, electromechanical simulators, and transistor circuitry, and is continuing his graduate work at Syracuse University.
13. I'. Bogert was born on September 26, 1923, in Waltham, Mass. He received the degree of 13.S. in physics in February, 1944. the M.S. degree in

B. Р. Boglent mathematics in 1946 and a I'h.I. in mathematics in 1948 from Massachusetts Institute of 'Technology.

During 1944 and 1945, he was a staff member of the M.J.T. Radiation Jaboratory. In 1948, he joined the technical staff of the Bell Telephone Laboratories, Murray Ilill, and has been primarily concerned with research in physical acoustics, and in narrow band speech transmission.

He is a member of Sigma Xi, and a Fellow of the Acoustical Society of America.
\&
IV'. F. Chow (II'53-SM'5.3) was born in Shanghai, China, on June 7, 1923. After receiving the $B . S$. degree in electrical engineering from Ta Tung

IV. F. Chow University in 1945, he joined the Chapei Power Co. In 1948 he came to the I'nited States receiving his M.S. degree in electrical engrineering in 1949 and his Ph.D. in electrical engineering in 1952, beth from the Eniversity of Minnesota, while he was serving there as a teaching assistant.

Dr. Chow joined the General Electric Co. in 1952. He is engaged in the research and development of transistor circuitry.

Ir. Chow is a member of Eta Kappa Nu and Sigma Xi.
B. F. C. Cooper ( $\mathbf{M}^{\prime}+77$ ) was born in England in 1917. He received the degrees of Bachelor of Science in 1939 and Bachelor of Engineering in 1941

B. F. C. Cooper from the I'niversity of Sydney. He has been a member of the research staff of the Division of Radiophysics, Commonwealth Scientific and Inclustrial Research Organization since 1940.

During World War II he worked on many developmental aspects of ground and airborne radar, and in the immediate post-war period he was responsible for the design of an airborne distance measuring equipment. In 1946-47 he spent a period with the National Research Council of Canada, where he developed an airborne ground-profile recorder.

Since returning to Australia he has developed instrumentation for rain-physics research, and also a magnetic drum storage system for the C.S.I.R.O. MK I Computer. At present he leads a group working on transistor electronics.

Martin I'eter was born July 12, 1928, in Switzerland. He received the diploma with distinction in physios and mathematics from the Eidgenössische Technische Hoch-

M. Peter schule in Zurich in 1952.

Mr. Peter was awarded the Kern prize for his thesis on colloidal ferroelectrics.

Mr. Peter subsequently came to the [nited States and entered the Massachusetts Institute of Technology, where he is a candidate for the Ph.I). degree in the Department of Ihysics. While studying for his degree, he is also a research assistant in the Rescarch Laboratory of Electronics.

Mr. Peter is a member of the American Physical Society and is an associate of Sigma Xi .

## $\therefore$

For a photograph and biography of M. IV. I. Strandberg, see page 756 of the June, 1955 issue of the Procerdings of the IRE.
R. WV. Grow (S'48-A'52) was born in Lymndyl, Utah, on October 31, 1925. He received the B.S. and M.S. degrees in clectrical engineering from the

R. IV. Grow University of Utah in 1948 and 1949 , respectively, and the Ph.D. degree in electrical engineering from Stanford University in 1955. From 1952 to 1953 he was an RCA Fellow in Electronics, under the Na tional Kesearch Council.

From 1949 to 1951, Dr. Grow was employed as an electronic scientist at the I'. S. Naval Research Laboratory, where he worked successively in the fields of radar countermeasures and nuclear physics, participating in the 1951 Atomic Weapons 'rests. In 1951 he became associated with the Electronics Research Laboratory at Stanford C"uiversity, where he has been engaged in microwave tube research specializing in traveling-wave tubes and backward-wave oscillators. At present he is a Research Associate at the Applied Electronics Laboratory.

Dr. Grow is a member of Signa Ki, Tan Beta Pi, Phi Kappa Phi, and Phi Eta Sigma.

## $\%$

E. Keonjian (M'50-SM'52) was born in Tiflis, Russia. He received his B.S. and M.S. degrees from the Leningrad Institute of Electrical Engineer-

E. KEONJIAN ing in 1932.

- Ifer graduation, Mr. Keonjian joined the Leningrad Central Radio-laboratory and in 1935 was transferred to the Leningrad Kesearch Institute of Electronics as a senior engineer. In 1940, he taughtelectrical communication at Leningrad Institute of Electrical Engineering.

In 1947, Mr. Keonjian cane to this country and joined Westinghouse Electric Corporation. In 1949, he was appointed as a lecturer in electrical communication at the City College of New York. Since 1951, Mr. Keonjian has been a member of the engineering staff of General Electric, where he is engaged in the development work concerned with transistor circuitry,

Mr. Keonjian is an author of numerous works on electronics, published here and abroad. He is a member of the Research Society of America, holds a professional engineeering license in the State of New York, and is a co-author of the book, "Principles of Transistor Circuits."
R. C. Knechtli was born in Geneva, Switzerland, on August 14, 1927. He received his M.S. degree in electrical engineer-
ing in 1950, from the Swiss Federal Institute of Terhnology (E.'Т.H.). He was employed by Brown Boveri and Company (Baden, Switzer-

R. C. Knechtli land), from 1951 to 1953 where he engaged in research on microwave circuits and microwave tubes. From 1951 to 1952, he was a research assistant at the Massachusetts Institute of Technology, where he also worked on nicrowave tubes.

Since 1953, he has been it rescarch engineer with RCA Iaboratories, at Princeton, New Jersey.
J. D. Kraus (A'32-M'43-SM'43-F'54) was born at Smo . Trbor, Mich., on June 28, 1910. He attended the I'niversity of Michigan receiving the l3.S.

J. D. Ǩraus degree in 19.30 , M.S. in 1931, and I'h.D. degree in physics in 1933. I Hring the next few years he was active on industrial noise reduction problems and in nuclear researeh. From 19,38 to 1940 I r. Kraus was an antenma consultant. From 1940 to 1943 , he was physicist and division head at the U. S. Naval Ordatace Laboratory, Washington, I). C. and between 194.3 to 1946 was research associate and group leader at N.W.R.C.'s Radio Rescarch Iaboratory at Harvard Cniversity.

In 1946 I $)$. Kraus joined the staff of the ()hioState [ niversity where he is now Professor of Electrical Engineering and Director of the Katlo ()bservatory. Professor Kraus is the author of books on antemmas and electromagnetic theory.

He is a member of the American Astrommical Society and the American I'hysical Society.
J. (i. Linvill (A'49) was born on August 8, 1919, in Polo, Mo. He received an A.B. degree from William Jewell College in 1941 and continued his
J. G. Linvilil
 studies at Massachusetts Institute of Technology, where he was awarded the S.I3. in 194.3, S. M. in 1945, and Sc.I). in 1949, all in electrical engineering. While at M.I.T. he was a member of the faculty, serving as assistant professor in electrical engineering from 1949 to 1951.
At the same time, he was a consultant to Sylvania E:lectrical I'roducts.

In 1951, Mr. Linvill joined Bell Telephone Laboratories, where he worked on ac-
tive network problens involving applications of transistors as the active element. Since March of this year, he has been Associate Professor of electrical engineering at Stanford University.

He is a member of the American Institute of Electrical Engineers, Sigma Xi, and Eta Kappa Nu.

For a photograph and biography of J. R. Macdonald, see pages 1571-1572 of the October, 1954 issue of the Procemincos of the IRE.
S. Matt ( $\Lambda^{\prime} 53$ ) was born in Cleveland, Ohio, on September 3, 1923. He received the Bachelor of Science degree in electrical engineering from Ohio

S. Matt I'niversity in 1944. Following graduation, he served with the U. S. Army Air Force. He returned to Ohio University in 1946 to teach in the Department of Electrical Engineering.

In 1947, he entered California Institute of Technology and obtained his M.S. degree in electrical engineering in 1948. IIe remained there two years as a research assistant in the High Voltage Laboratory. He then entered Ohio State I Iniversity and received his I'h.l). degree in 1953.

Dr. Matt was a research assistant and then an instructor while at the Ohio State University. At present, he is at the General Vlectric Advanced Electronics Center at Ithaca, N. Y.

Dr. Natt is a member of Sigma Ni amd I' Mit Epsilon.
A. R. Moore was born in New York, N. Y., on Jamary 14, 1923. He received the B.S. degree in chemistry in 1942 from the ['olytechnic Institute

A. R. Moore of Brooklyn. He worked on phototube and thyratron development at RCA Victor in Harrison, N. J. and I.ancaster, Pa., from 1942 to 1945. In 1945 he entered Cornell Itiversity, and received the Ph.1). in physics in 1949, specializing in physics of solids During his last two years at Cornell he was an RC. Fellow. He joined the RC.I Laboratories IVivision at I'rinceton, N. J. in 1949, where he has worked on semiconductor physics.
R. F. Rutz (A'51) was born in Alton, Ill., on February 9, 1919. He received the B.A. degree in 1941 fromShuritdeff College, Alton,
III., and the M.S degree in physics in 1947 from the State I niversity of Iowa.

Mr. Rutz joinerl the Sandia Laboratory of the Los Ala-

R. F. Rutz mos Scientific Research Laboratory, subsequently Sandia Corporation, at Albuquerque, N. M. in 1948, where he worked in the Electronics Research Department. In 1951 he joined the Research Laloratory of the International Business Machines Corporation at Poughkeepsic, N. Y., where he is currently working on transistor research.
lle is a member of Sigma Xi , and the American Physical Scciety.
A. P. Stern (.1'51i was born on July 20 , 1225, in Budapest, Hungary. He studied at the L'niversities of Budapest and Lausame and at the Swiss Fed-

A. P. Stern eral Institute of Technology in Zurich, where he acquired a Master's degree in electrical engineering in 1948.

From 1948 to 1951, Mr. Stern was engaged in research and devilopment work in the field of gaseons discharges in Switzerland. In 1950, he became Instructor for illumination engineering and photometry at the Swiss Federal Institute of 'Technology'.

Mr. Stern came to the United States in 1951 and joined the staf of the General Electric Company's Blectronics Laboratory in Syracuse, N. Y. . It the present time, Mr. Stern is supervisor of solid state circuit development in the Electronics Laboratory.

Mr. Stern is a member of the Scientific Research Society of America.
K. F. Stripp was Jorn on Jamary 26, 1920, in Union (ity, N. I. He obtained a BS. degree in chenistry from the Polytectnic Institate of

k. F. Stripl Brooklyn in 1950. He was teaching assistant in the chemistry department at California Institute of Technology from 1950 to 1951 and obtained his lh.I). in physical chemistry from Yale in 1953. He spent three summers at R(.) Laboratories and had been assocrated with them on a full-time basis since July 1, 1953.

Dr. Stripp was a member of the American Chemical Socicty, AAAS, American Institute of Chemists, Sigma Ki, and Phi Lambda Upsilon. IIe died suddenly on July 10, 1954.
J. J. Suran ( A '52) was barn in New York, N. Y., on January 11, 1926. After having served for three years with the U. S. Army during World War II,

J. J. Suran he received the B.S.E.E. degree from Columbia University in 1949 and continued graduate studies there and at the Illinois Institute of Technology.

From 1949 to 1953, Mr. Suran was employed in control systems de:ign and development by J.W. Meaker and Co., and in FM communications research and development by Motorola, Inc. Since 1952, he has been a member of the Electronics Laboratory of the General Electric Company.

Mr. Suran is a co-author of the book, "Principles of Transistor Circuits." Ife has a professional engineering license in New York State and is a nember of VIEE.
A. W. Warner (M'52) was born in Sewickley, Pa. in 1915. He received the B.A. degree, with a major in physics, from the Dniversity of Delaware in 1940 and the
M.S. degree in physics from the U'niversity of Maryland in 1942.

In the same year Mr. Warner was a

A. W. Warner member of the faculty of Lehigh University, leaving in July to join the Western Electric Company, where he worked on the development of crystal-unit test equipment. In 1943 Mr. Warner became a member of the technical staff of Bell Telephone Laboratories, and is engaged in the design of high-frequency plated crystal units.

For a photograph and biography of D.A. Watkins, see page 106 of the January, 1955 issue of the Procerdings of the IRE.
H. A. Wheeler (A'27-M'28-F'35) was born in St. Paul, Minn., on May 10, 1903. He received the I3.S. degree in physics from George Washington Cniversity in 1925. From 1925 to 1928 he studied in the physics department of The Johns Hopkins Cniver-
sity, and lectured there during 1926 and 1927. He was employed as a laboratory assistant in the radio section of the National Bureau of Standards

H. A. Wheeler in 1921, leaving in 1923 to assist Professor Hazeltine and later to join the Hazeltine Corporation in 1924. He was in charge of their Bayside laboratory from 1930 to 1937, where he later became vice-president and chief consulting engineer.
He has specialized in the design of radio receivers (including FM and TV), the theory of communication networks, radar (including IFF during World War II), antennas, and microwave equipment.

In 1946 Mr . Wheeler opened his own consulting office in Great Neck, N. Y. He is now also president of Wheeler Laboratories, Inc. From 1950 to date, he has been serving part-time as consultant to the Office of Secretary of Defense in the fields of guided missiles and electronics.

Mr. Wheeler is a Fellow of the American Institute of Electrical Engineers, an Associate of the Institution of Electrical Engineers, and a member of Sigma Xi. He received the Morris Liebmann Memorial l'rize in 1940.

## IRE News and Radio Notes

## Second Symposium on Vacuum Trechnology Invites I'apers

The Committee on liacuum Techniques, Incorporated invites the submission of papers at the Second Symposium on Vacuum Technology to be held at Mellon Institute in Pittsburgh, October 13-15. Those interested should write to Rudy Fiohler, Committee on Vacuum Technique:s, Inc., I3ox 1282, Boston 9, Massachusetts.

The program will deal with equipment, instrumentation, developments in vacuum technology, standards, nomenclature, methords and techniques, and vacuum systems applications and processes

## M.I.T. Gives Summer Course on Noise in Electron I)evices

Ilans for a two-week Special Summer I'rogram on Noise in Electron Devices at the Massachusetts Institute of Technology have been announced by Ernest H. Huntress, Director of the M.I.T. Summer Session. The program will be held from July 18 through July 29 and is planned primarily for those who are or plan to become research workers in the field.
L. D. Smullin, of the Research Laboratory of Electronics at M.I.T., and H. A. Haus, Assistant Professor in the Department of Electrical Engineering, will direct the program. Other members of the M.I.T. staff who will lecture include I'. Elias, Assistant Professor of Electrical Engineering; Y. W. Lee, Associate Professor of Electrical

Engineering; and L. Tisza, Associate Professor of Physics.

Guest lecturers will include D. O. North and I. W. Peter, both of the David Sarnoff Research Center; C. F. Quate and T. E. Talpey, of the Bell Telephone Laboratories; and A. van der Ziel, I'rofessor of Electrical Engineering at the ['niversity of Minnesota.

Full details and application blanks for this Special Summer Program may be obtained from the Summer Session office, Room 7-103, Massachusetts Institute of Technology, Cambridge 39, Massachusetts.

## Instrumentation Conference to Meet in Atlanta

The Professional Group on Instrumentation and the Atlanta Section will sponsor an Instrumentation Conference and Exhibit at the Biltmore Hotel in Atlanta, November 28 through November 30.

The theme of the conference is "Data Handling." Prospective authors of papers are invited to submit abstracts of 200 words or less on data gathering, processing, utilization, and processing systems not later than September 1. Titles and abstracts should be addressed to B. J. Dasher, School of Electrical Engineering, Georgia Institute of Technology, Atlanta, Georgia.

## Audio Engineering Society Will. Meet October 12 to 15 in N. Y.

"Practicality" will be the theme of the

1955 Convention of the Audio Engineering Society Twenty-five, which is scheduled for October 12 to 15 . Sessions will be held in New York at the Hotel New Yorker and will be concurrent with the Audio Fair.

According to Richard II. Ranger, president of Rangertone, Incorporated and program chairman for the event, theronvention will include panel discussions on transistors, amplifier design and tape recording. Their purpose, he said, will be to bring out the "right and easy way" to handle each type of equipment. The agenda will also include theoretical and scientific papers. The annual banquet of the society will be held on the evening of October 12.

Col. Ranger, who is Executive Vice-I'resident of the society this year, is being aided by Effingham Kettleman of RCA.

## Final Call for <br> PGED Papizks

The PGED, which will holds its First Annual Technical Meeting in Washington, D. C., October 24 and 25 , is now making its final call for papers.

The meeting will include three parallel sessions: Solid State Devices, Microwave Tubes, and Non-Microwave Tubes.

## Nominations for 1956 Officers

At its May 14, 1955 meeting, the IRE Board of Directors received the recommendations of the Nominations Committee and the reports of the Regional Committees for officers and directors for 1956. They are:
President, 1956-A. V. Loughren
Vice-President, 1956-Herre Rinia
Director-at-Large, 1956-1958 (two to be elected) L. V. Berkner, E. W. Herold, T. A. Iunter, J. R. Whinnery

Regional Directors, 1956-1957 (one to be elected in each Region)
Region 1-C. R. Burrows, H. F. Dart, L. B. Grew

Region 3-J. G. Brainerd, L. R. Quarles
Region 5-J. J. Gershon, R. E. Moe
Region 7-1.. E. French, C. F. Wolcott
According to Article VI, Section 1, of the IRE Constitution, nominations by petition for any of the above offices may be made by letter to the Board of Directors, giving name of proposed candidate and office for which it is desired he be nominated. For acceptance a letter of petition must reach the executive office before noon on August 12, 1955, and shall be signed by at least 100 voting members qualified to vote for the office of the candidate nominated.

## Vehicular Communications Paper Deadline Announced For the First of August

The Professional Group on Vehicular Communications will hold its Sixth Anmual Mecting September 26 and 27 at the Multnomah Hotel, Portland, Oregon.

Deadline for papers is August 1; title of paper, abstract, full name and address should be submitted to Newton Monk, Bell Telephone Labs., 463 West Street, New York 14, N. Y.

## Art in Electronics ('ompetition Announced for WESCON Silow

The West Coast Electronic Manufacturers Association is sponsoring a competition in art as applied to electronics during the 1955 IVESCON Show which will be held in San Francisco, California, August 24-26.

Employees of electronic manufacturers and their families are eligible to compete for cash and vacation trips in a contest planned to augment the association's annual scholarship program. Entries will be exhibited at WESCON and auctioned to the highest bidders. Proceeds from the sale will go to the WCEMA Scholarship Fund.

Purpose of the contest, according to Paul M. Cook and Mrs. Jan Smith, co-chairmen for the event, is to stimulate an awareness of the visually attractive materials, components and designs of electronics and their possible use as objects of art.

Trip prizes will be awarded in each of four categories: photography, painting and sculpture, decorative accessories, and jewelry. All entries must be created from component parts commonly used in electronics.

A prize of $\$ 250.00$ will go to the creator of the object bringing the largest price in the auction.

Complete details and entry blanks for the competition are available from Berk Baker, Eitel-McCullough, Incorporated, 798 San Mateo Avenue, San Bruno, California. They may also be obtained at the WCEMA office, 339 South Robertson Boulevard, Beverly Hills, California.

## Benjamin Bauer and Kenneth Goff Honored by PG on Audio

Under a plan approved last year by the Administrative Committee of the Professional Group on Audio, awards were presented during Audio Sessions at the National Convention to Kenneth E. Goff and IB. B. Bauer. Mr. Bauer, of Shure Brothers, Incorporated, received $\$ 200$ "in recognition of many excellent audio papers appearing in IRE publications


Kenneth W. Goff over a period of years." Mr. Goff, who is associated with the Acoustics Laboratory at M.I.T., received $\$ 100$ as the author, under 30 years old, "of an especially meritorious paper dealing with a subject related to audio appearing in any IRE publication." The name of Mr. Goff's paper was The Development of a Variable Time Delay and appeared in the 1953 Convention Rec-

B. B. Bauer ord of the Institute of Radio Engineers.

Mr. Goff graduated in 1950 from West Virginia I'niversity with the B.S. degree in electrical engincering. ITpon graduation, he joined the M.I.T. Acoustics Laboratory and entered the M.I.T. Graduate School. In accepting the audio award, Mr. Goff said that the "tremendous challenge of the unsolved problems in the field of Audio, together with the opportunities for organization and individual recognition made possible by the PGA, combine to present a very attractive picture to those of us who are just beginning our work in electrical engineering."

Mr. Bauer received the E.E. degree in 1937 from the University of Cincinnati. He has also received the Industrial Electrical Engineering degree from Pratt Institute and has attended the I'niversity of Chicago and Illinois Institute of Technology. Mr. Bauer suggested, in accepting his award, that the money be used for work among PGA Student Members. "I suggest," he said, "the creation of a fund, which could be augmented from time to time by PGA proper, and used for annual Student Awards, for meritorious papers on subjects connected with Audio Technology."

## Calendar of Coming Events

SRI and Nat. Ind. Conf. Board Symposium on Electronics in Automatic Production, Sheraton-Palace, San Francisco, Calif., Aug. 22-23
URSI Symposium on Solar Eclipses and the Ionosphere, Royal Society, Burlington House, London, England, Aug. 22-24
IRE-West Coast Electronic Manufacturers' Association WESCON, Civic Auditorium, San Prancisco, California, Aug. 24-26
Emporium Section Sirteenth Annual Summer Seminar, Emporium, Pa., August 26-28
IRE-ISA Tenth Annual Instrument Conference, Shrine Auditorium, Los Angeles, Calif., Sept. 12-16
Association for Computing Machinery, Annual Meeting, Moore School of Electrical Engineering, U. of I'a., Sept. 14-16
IRE Professional Group on Nuclear Science-Second Annual Meeting, Oak Ridge National Labs., Oak Ridge, Tenn., Sept. 14-17
IRE Cedar Rapids Section Symposium on Automation, Cedar Rapids, Ia., Sept. 17
RETMA Automation Symposium, U. of Pennsylvania, Philadelphia, Pa., Sept. 26-27
PG on Vehicular Communications Sirth Annual Meeting, Multnomah Hotel, Portland, Ore., Sept. 26-27
IMSA Annual Convention, Hotel Seneca, Rochester, N. Y., Sept. 26-29
International Analogy Computation Meeting, Société Belge des Ingenieurs des Télécommunications et d'Electronique, Brussels, Belgiam, Sept. 27-Oct. 1.
IRE-AIEE Conference on Industrial Electronics, Rackham Memorial Building, Detroit, Michigan, Sept. 28-29
National Electronics Conference, Hotel Sherman, Chicago, Ill., October 3-5 Audio Engineering Society Convention, Hotel New Yorker, New York City, Oct. 12-15
Second Symposium on Vacuum 'Technology, Mellon Inst., I'ittsburgh, Pa., Oct. 13-15
IRE-RETMA Radio Fall Meeting, Hotel Syracuse, Syracuse, N. Y., Oct. 17-19
Eighth Annual Gaseous Electronics Conference, General Electric Res. Lab., Schenectady, N.Y.,()ct.20-22
PG on Electron Devices Annual Technical Meeting, Shoreham Hotel, Washington, D. C., Oct. 24-25
IRE East Coast Conference on Aeronautical and Navigational Electronics, Lord Baltimore Hotel, Baltimore, Md., Oct. 31-Nov. 1

Symposium on Applied Solar Energy, Westward Ho Hotel, Phoenix, Ariz., Nov: 1-5
IRE-AIEE-ACM Eastern Joint Computer Conference, Hotel Statler, Boston, Nov. 7-9
IRE-AIEE-ISA Electrical Techniques in Medicine and Biology, Shoreham Hotel, Washington, D. C., Nov. 14-16
PGI and Atlanta Section Data Processing Symposium, Hotel Biltmore, Atlanta, Ga., Nov. 28-30

# LONG ISLAND SECTION HONORS FOUR IRE FELLOW RECIPIENTS AT SPECIAL AWARD CEREMONIES 

Citations Presented to Loughlin, Learned, Gaffney, and Dunning



Fellow citations are preserted members by John Dyer as Pres. Fyder wat hes. Left to righ are B. B. Loughlin, V. R. Learned, F. J. Gaffney, D. M. Dunning. J. N. Dyer, and John Ryder.

A special award meeting was held by the Long Island Section on March 20 to honor four Section members who were made Fellows of the IRE. John Dyer, Regional Director, presented the award citations and President John Ryder made the principal address. After the award presentations, a cocktail party was held in the main ballroom of the Garden City Hotel.

The four Long Island Section members who received the awards were Orville M. Dunning, Francis J. Gaffney, Vincent R. Learned, and Bernard B. Loughlin. Mr. Dunning, a member of the Board of Directors and Vice-President in Charge of Engineering at Hazeltine Ccarporation, received the Fellow award " . . . for his contributions to the field of sound recording and his effective organization of engineering effort." Vice-Iresident for Engeneering at Marion Electrical Instrument Company, Mr. Gaffney was made a Fellow" . . .for his contributions to the field of electrical Measurements." ". . . for his contributions to research and development of microwave electron tubes," Dr. Learned, of Sperry Gyroscope, received the Fellow award. A Consulting engineer with Hazeltine Corporation, Mr. Loughlin was made a Fellow ". . for his contributions to color television, frequency modulation, and superregeneration."

## John R. Pierce Elected to National Academy of Sciences

John R. Pierce, Editor of IRE and Director of Electronics Research at Bell Telephone Laboratories, was elected to the National Academy of Sciences at its 92nd annual meeting held in Washington recently.

The National Academy of Sciences, a private, non-profit organization, serves as an adviser to the Federal Government in scientific matters and acts in the furtherance of science for the general welfare. The membership of the academy numbers approximately 500 in the physical and biological fields.

In addition to Dr. Pierce, IRE members who have been elected to the academy include M. J. Kelly, J. B. Fisk, W. Shockley, L. V. Berkner and F. E. Terman.

## Program for Applied Solar Energy Symposiun Announced

A preliminary program has been announced for the World Symposium on Applied Solar Energy to be held in Phoenix, November 1 through 5. Sponsors of the meeting are the Association for Applied Solar Energy, Stanford Research Institute, and the University of Arizona.

Among the papers scheduled are: The Sun's Energy, Farrington Daniels, University of Wisconsin; Survey of the Domestic Uses of Solar Energy, H. C. Hottel, Massachusetts Institute of Technology; Space Cooling With Solar Energy, George O. G. Lijf, Denver, Colorado; Food and Fuel from Solar Energy, F. A. Brooks, University of California; Chlorella for Animal Food, Jack Meyers, University of Texas; Engineering for Algae Culture, A. W. Fisher, Jr., Arthur D. Little, Inc., Cambridge, Mass.; and Solar Energy Utilization by Higher Plants, Paul C. Mangelsdoff, Harvard University. Maria Telkes, New York University, Solar Stills; R. C. Jordan, University of Minnesata, Mechanioal Energy from Solar Energy; Paul Erlandson, Southwest Research Institute, Direct Conversion of Solar Energy; L. J. Heidt, Massachusetts Institute of Techol ogy, Converting Solar to Chemical Energy; G. L. Pearson, Bell Telephone Laboratories, Photovoltaic P-N Couples. J. E. Hobson, director of Stanford Research, will lead a discussion on "The Economics of Solar Energy" at the outset of the meeting

Contributions from abroad will be made by: Felix Trombe, Laboratoire de l'Energie Solaire, Paris. High-Temperature Furnaces; R. N. Morse, Commonwealth Scientific and Industrial Research. Organization, Australia, Solar Water Heaters; and Hiroshi Tamiya, Tokugawa Institute for Biological Research, Tokyo, Chlorella for Food.

General chairman for the symposium is Lewis W. Douglas, of the Southern Arizona Bank and Trist Company, Vice-Chairman and program director is Merritt L. Kastens, Assistant Director of Stanford Research Institute. Headquarters for the symposium during its planning phase are located in Suite 204, Mayer-Heard Building, Phoenix.

## PGNS Will Meet in September

The Second National Annual Meeting of the PG on Nuclear Science will be held in Oak Kidge, Tenn., September $1+16$, with the Oak Ridge Chapter as host.

Persons wishing to gresent papers are invited to submit, before July 15, 200-word abstracts to H. E. Banta, Papers Committee, Oak Ridge National Laboratory, Box P, Oak Kitge, 'renn. Since pasers will be accepted in 15,30 , and 45 minute categories, specification of the time required is requested.

## Professional Group News

## Eight New Chapters Approved

At its meeting of April 3, the IRE Executive Committee officially approved the following chapters: Los Angeles Chapter, PG on Automatic Control ; Los Angeles Chapter, PG on Reliability and Quality Control; Chicago Chapter, PG on Communications Systems; Pittsburgh Chapter, PG on Electron:c Computers; Atlanta Chapter, PG on Instrumentation.

On May 3 three chapters were approved by the Executive Committee. They were: San Francisco Chapter, PG on Audio; Twin Cities Chapter, PG on Automatic Control; Long Island Chapter, PG on Instrumentation.

At an earlier meeting the Northwest Florida Subsection was made a full Section and the For: Muachuca Subsection was established in the Phoenix Section.

## Technical Committee Notes

The Antennas and Waveguides Committee met at IRE Headquarters on April 20 with P. H. Sinith presiding. The committee approved the Proposed Standard on Antennas and Waveguides: Definitions for Waveguide Components for submission to the Standards Committee, Subcommittee 2.4 on Waveguide and Waveguide Component Measurements presented draits of "Methods
of Measurement" comprising discussions on Measurements of Phase Shift, Measurement of Power Handling Capacity, and Measurement of $Q$. A good deal of discussion followed, with the suggestion that this subject be re-examined by the subcommittee.

The Antennas and Waveguides Committee met at IRE Headquarters on March 9 with P. H. Smith presiding. The committee reviewed the Proposed Standards on Antennas and Waveguides, Definitions for Waveguide Components. W. E. Waller reported on the Proposed Standards: Waveguide and Waveguide Component Measurements.
D. E. Maxwell presided at the Audio Techniques Committee meeting at IRE Headquarters on April 21. L. D. Runkle, Chairman of Subcommittee 3.1 on Definitions, reported that after two meetings the subcommittee has adopted and tentatively approved 27 definitions. The committee reviewed the Proposed Standards on Methods of Measurement of Gain, Loss, Amplification, Attenuation and Frequency Response, which they expect to complete at their next meeting.
1). E. Maxwell presided at the Audio Techniques Committee meeting at IRE Headquarters on March 24. L. D. Runkle, Chairman of Definitions Subcommittee 3.1, reported that his subcommittee is reviewing a list of proposed definitions. R. C. Moody, Chairman of the West Coast Subcommittee 3.3, reported that his subcommittee is working on a Proposed Standard on the Measurement of Intermodulation Distortion, which they hope to finish by the end of the year. The committee reviewed the proposed Standards on Methods of Measurement of Gain, Loss, Amplification, Attenuation and Frequency Response.

The Electronic Computers Committee met at IRE Headquarters on March 24 with Robert Serrell presiding. After discussion of the work of the Definitions Subcommittees, Mr. Brown made and Dr. Haynes seconded the following motion, which was unanimously approved: "The Electronic Computers Conmittee requests that its Subcommittees 8.4 and 8.5 make plans to distribute glossaries of new terms being considered by them. The glossaries should be
sent with the PGEC Transactions and include all desirable explanations of the terms."

The Facsimile Committee met at IRE Headquarters on April 22 with Chairman H. Burkhard presiding. The committee discussed the Facsimile Test Chart which they are preparing. The following definitions were approved: electrostatic recording and magnetic recording.

The Facsimile Committee met at the Times Annex on March 18 and H. Burkhard presided. The committee reviewed their tentative definitions of terms. The following terms were proposed as possibly requiring definitions: xerographic recording, electrostatic recording, magnetic recording, ferrographic recording.

The Information Theory and Modulation Systems Committee met at IRE Headquarters on March 9 with J. G. Kreer, Jr. presiding. Dale Pollack was appointed chairman of the Modulations Systems Subcommittee and Peter Elias was appointed chairman of the Information Theory Subcommittee. The following terms were referred to the subcommittee on Information Theory for study: binary digit, message coding versus symbol coding, systematic versus unsystematic coding, and corrector and characteristic as used in coding.
H. R. Mimno presided at the Navigation Aids Committee meeting at IRE Headquarters on April 22. The committee commenced its exploratory examination of the proposed measurements standard on the VHF Omnirange. The discussion covered the Introduction, Description and Characteristics to be Measured, together with the introductory paragraphs on Specialized VOR Test Equipment. Mr. Moskowitz recorded certain editorial changes arrived at by general agreement, and noted additional suggestions to be considered by an editorial group.

The Piezoelectric Crystals Committee met at IRE Headquarters on March 21. W. P. Mason presided. W. L. Bond discussed nomenclature used at the recent meeting of the committee of the International Crystallographic Congress. Three proposals on methods of determining the piezoelectric, elastic, and dielectric constants of
crystals and the parameters of piezoelectric vibrators were discussed.

Erust Weber presided at the meeting of the Standards Committee at the Jade Room of the Waldorf-Astoria Hotel on March 24. Dr. Weber explained IRE standardization procedure to new members, and A. G. Jensen described the functions of the Standards Coordinator. There was a review of the past year's work and an announcement of future plans. Mr. Jensen introduced the new members of the Standards Committee. E. A. Laport expressed the appreciation of the group to the recent chairmen, Dr. Brainerd, Mr. Jensen, and Dr. Weber, for assuming the responsibility of leadership of the committee. He stated that the committee had been fortunate in having capable chairmen to lead the group.

The Standards Committee met at IR1: Headquarters on April 7 with Chairman E. Weber presiding. A report of the work of the Ad Hoc Committee on Spurious Radiation was submitted by R. F. Shea, with a recommendation that this committee be dissolved. This motion was unanimously approved. The formation of a Nuclear Techniques Technical Committee with G. A. Morton as chairman was unanimously approved. No action was taken on the Proposed Standards on Pulses: Methods of Measurement of Pulse Quantities or the Proposed Definitions on Induction and Dielectric Heating since there were no representatives of the sponsoring committees present.

The Video Techniques Committee met at IRE Headquarters on March 31 with W. J. Poch presiding. Mr. Jones reported on the recent activities of the Subcommittee on Video Transmission. The problent of specifying the exact conditions for taking measurements of differential gain and differential phase were discussed at length. Dr. Athey reported that he was making progress in collecting a list of standard terms in the field of kinescope recording. The following definitions were approved by the committee: flyback, linearity control, camera tule, contrast ratio, geometric distortion, de restorer, dc transmissior, return interval. return trace, raster, retrace interval, brightness control, nominal line width, progressive scanning.

## Sixth PGUE Group Administrative Committee Meeting March 22 During IRE Convention



PGUE officers are, left to right: Morris Kenny, retiring Secretary; Julia Herrick, Vice-Chairman; Amor Lane, Chairman of Membership Committee and retiring Group Chairman; Oskar Mattiat, Editor of PGUE Transactions; Morton Fagen, Chairman; Frank Massa, Chairman of Nominations Committee and retiring Ad- Study and Review Comnittee; and Donald Berlincourt, Associate Editor

# 1955 Western Electronic Show and Convention 

## Tentative Program <br> San Francisco; California August 24-26

The 1955 Western Electric Show and Convention will meet in San Francisco on August 24, 25 and 26. This year there will be 570 exhibits representing more than 600 producers. Co-sponsored by the West Coast Manufacturers' Association and the San Francisco and Los Angeles Section of the Institute of Radio Ergineers, representing the Seventh Region, WESCON will be attended by more than 20,000 visitors.

A I'nited Airlines "Airlift" has been arranged to transport part of the visitors who will attend. The Airlift was arranged by Noel E. Porter of the Hewlett Packard Company and WESCON Chairman, and Mal Mobley, Jr., Business Manager. In addition to the many ("nited schedules, special Manliner flights will be arranged exclusively for WESCON visitors and exhibitors from major cities to San Francisco. United Airlines has arranged to handle reservations on its own line or any other scheduled airline. Confirmation wilf be made directly by (Tnited or through any local airline office or travel agent. Show and Convention officials; urged that reservations be made as soon as possible in order to be assured of the best flights and schedules.

The Technical Program will consist of twenty-four sessions and over 100 papers. This specialized program has been closely integrated by the program committee which is made up of IRE; members, coordinating both Professional Groups and Section activities.

## Wednesday Morning

## Solid State Devices

Transistors Today, J. A. Morton
Large Signal Semi-Conductor Devices, John Saby
Iligh-Frequency Power Gain of Junction Transistors, R. L. Pritchard
Recent Developments in Germanium Alloy Junctions, C. W. Mueller
A Now Iigh-Ambient Transistor, R. R. Rutherford and J. J. Bowe

## Information Theory

Limiting Frequency-Modulation Spectra, N. Blachman

The Definition of a General Metric of Information, N. Abramson
An Analysis of Optimum Sequential Detectors, J. J. Bussgang and D. Middleton
Analysis of Automatic Bias Control for Threshold Detectors, E. Ackerlind
Generating a Gaussian Sample, S. Stein and J. E. Storer

Proof of the Sampling Thearem for Stationary Processes, A. Rosenbloom and J. Heilfron

## Reliability and Quality Contral <br> Engineering and Testing for Reliability, H. G. Romig

Parts Versus Systems: The Reliability Dilemma, David A. Hill
An Effective Reliability Program Based Upon "A Triad for Design Reliability," F. E. Dreste
A Basic Study of the Effects of Operating and Environmental Factors on Electron Tube Reliability, W. S. Bowie
Surface Contamination of Dielectric Materials, Saul Chaikin

## Propagation

An Explanation of Fading in Microwate Relay Systems, H. Maynuski
Some Notes on Propagation ozer a Spherical Earth, S. J. Fricker
Radio I'ower Received via Tropospheric Scattering, A. Waterman
Atmospheric Attenuation of Microzeave Radiation, G. R. Marner
Theory of Deviative Absorption in the F? Layer and its Relation to Temperature, R. Gallet

Symposiunt on Industrial Electronics and Nuclear Engineering

## Wednesday Afternoon

## Broadcast and TV Receivers

A Thin Cathode Ray Tube, William R. Aiken
Beam Focusing and Deflection in the Aiken Tube, R. Madey
Radiation Measurements at VIIF and UHF, A. B. Glenn

An Experimental Automobile Receiver Employing Transistors, 1.. A. Freedman, T. O. Stanley and 1). I). Holmes

High-Efficiency, Unipotential Post Focus, Tri-Color P'icture Tube, Wilfrid F. Niklas

## Circuit Theory I - Transistors and Blocking Oscillators

Advantages of Direct Coupled Transistor Amplifiers, Richard Hurley
Junction Transistor Blocking Oscillators, J. G. Linville

The Design of Blocking Oscillators as Fast Pulse Regenerators, F. K. Bowers
Slability of 1 Iulti-Mode Oscillating Systems, R. W. De Grasse
(Additional paper to be announced.)
Electronic Instrumentation in AircraftJoint Symposium of the Professional Group on Aeronautical and Navigational Electronics and the Institute of deronautical Sciences
Experiments with Radio Controlled, Dynamically Similar Models, E. G. Stout
Rote of Eilectronics in Engineering Flight Testing, W. L. Howland
Instrumentation for Rocket Engine Testing, R. F. Compertz
(Additional papers to be announced.)

## Antennas I

Recent Developments in Microwave Antcnnas, L. C. Van Atta

Printed Surface Wave Antennas, H. W. Cooper
Cireularly-Polarized Slot Radiators, A. J. Simmons

Radiation from Ferrite-Loaded Slot Radiators,
D. J. Angelakos and M. Korman

A Large Aperture Differential Polarization Antenna for Radio Astronomy Use, V. H. Goerke and O. D. Remmler

## Instrumentation

Beamplexer-High Speed Channel Mulliplexing Unil, 11. Moss and S. Ḱuchinsky
A Stable Diode Chopper Circuit, H. Patton
A Completely Automatic Impedance I'lotter, J. R. \inding

A Broadband Microwave Frequency Meter, P'. H. Vartanian and J. L., Melchor
An Expanded Sale Froquency Meter, Duane Marshatl
Measurement of Time Varying Frequencies, Martin Craham

## Thursday Morning

## Electronic Component Parts

Design and Properties of High l'oltuge Glass Capacitors, G. I'. Smith
Characteristics of Modular Electronic Components, IV. G. James
Simple Electronic Transformer Design, R. I.ee
Mrasurement of P'arameters Controlling l'ulse Front Response of Transformers, I. R. Gillette, K. Oshima and R. M. Rowe
Development of MIL-T-27-A: Transformers and Reactors, İ. M. Wiler
International Research in Electronics and Allied Fields. Symposiunt I-The Role of the IRE and URSI

## High Power Tubes

M-Type Backwerd Wrave Oscillators, J. IHull Considerations of Various Struclures for Ifigh Average Powers in the UHF Region, I). Preist

Design Informalion on Large Signal Travel-ing-I'ave 1 mplifiers, J. E. Rowe
A New Beam Power Tuhe for UIIF Service, W. B. Bemmet

An Ion Trapped IIigh V'oltage Pentode, R. E. Hellers

## Automatic Control

Non-Linear Compensation of an Aircraft Instrument Servo-mechanism, I). Lebell
The Stabilization of Von-Linear Servomechanisms Encountered in Antenna Instrumentation, J. Bacon
Synthesis of a Non-Linear Control System, I. Flugge-Lotz and C. F. Taylor

Theory of Non-Lincar Fecdback Systems Having a Multiple Number of First-Order Operating Points, J. A. Narul
Noise th Non-Linear Servos, G. O. Young and C. J. Savant

Telemetry and Remote Controi. Wow and Flutter Compensation in FM Tclemetry, W. H. Chester
Aliasing Errors in. Sampled Data Dystems, A. J. Mallinckrodt

Air-to-Ground Propagation over Descrt Terrain at Telemetering Frequencies, G. I.. McCone
ulse Width Data Mulliplexing of an FM／FM Subcarrier，A．S．Westuest The Use of A C Excited Gauges in a PDM／PM Telemetering System，W．F．Carmody

## Thursday Afternoon

## Microwave Theory

Periodic Structures for Traveling－Wave Tubes， M．Chodorow
Conversion of Maxwell＇s Equations into Gen－ eralized Telegraphist＇s Equations，S． 1. Schelkunoff
On the Expansion of Fields in Lossless．Micro－ wave Junctions，＇T＇．「「eichmann
Conformal Mapping of Rounded Polygons by a Wave－Filter Analogue，H．A．Wheeler

## Broadcast Transmission Systems

The Perfect Television System，O．H．Schade The Subjective Sharpness of Sinulated Color TV Pictures，H．F．Huntsman
The Conversion of a Standard TV Mobile Unit for Greater Flexibility and Operating Convenience，H．F．Huntsman
High Speed Duplication of Magnetic Tape Recordings，J．M．Leslie
Color TV Magnetic Tape Recording System， H．F．Olson

## Computers I－Digital Computer Applications and Design Techniques

A Punched Card Method of Evaluating Sys－ tems of Boolean Functions with Special Reference to Analysis of Relay Circuits， IV．R．Abbott
The Elecom 50－A New Type of Computer， Evelyn Berezin and Jhyllis Hersh
Logical Design of the Remington Rand High Speed Printer with Emphasis on the Check－ ing and Ediling Features，M．Jacoby
Theory，Principles and Applications of Sta－ tistical Computers，H．Blasbalg and W．O＇llare
A Glow Transfer Shifting Register Utilizing $R-F$ Gas Discharge，D．C．Engelbart
Ferroelectric Hysteresis in Barium Titanate Single Crystals，H．H．Wieder

## Engineering Management

Snall Engineering Company Organization－ a Philosophy and Method，T．W．Jarmie
Is the Yardstick for Estimating Individual Engineering and Scientific Potential Reli－ able？A．H．Schooley
Management in Production Engineering， C．Blahna
Market Development－The Neglected Com－ panion of Product Development，A．D． Ehrenfried
Cross Functional Engineering Managemetn， C．M．Ryerson

Aeronautical and Navigational，

## Electronics

An Improved Simultaneous Phase Compari－ son Guidance Radar，H．H．Sommer
Antenna Design Considerations for Heli． copters，J．B．Chown
High Voltage Impulse Generation for Meas－ urement of Receiver Susceptibility to Inter－ ference Encountered in Aircraft，A．Newan and J．R．Stahmam
Experimental Results of Conductive Cooling Tests on Airborne Equipment，R．L．Ber－ ner

## Thursday Evening

．Medical Electronics Panel Discussion

## Friday Morning

## Computers II－Analogue

 Computer Components and ApplicationsAutomatic Data Accumulation System for Wind Tunnels，John Wedel
Data Recorder for Evaluation of a Fire Con－ trol System，J．T．Ator and L．P．Retzinger， Jr．
Transistors in Current Analog Computing， B．P．Kerfoot
The Use of Electronic Analog Computers in the Solution of Certain Radur Noise Prob－ lems，J．A．Aseltine
Precision Electronic Switching with Feedbuck Amplifiers，C．M．Edwards

## Circuit Theory II－ Synthesis Problems

New Methods of Transformerless Driving－ Point Impedance Synthesis，Stanley Ifurst General Synthes is of Quarter－Wave Impedance Transformers with Given Insertion Loss Function，Henry J．Riblet
The Approximation Problem in the Synthesis of $R-C$ Networks，K．L． Su and B．J． Dasher
A Precise Method of Designing High－and－
Low－Pass R－C Fillers with Active Elements， M．McWhorter
Signal Flow Graphs for Random Signals， W．H．Huggins

## Medical Electronics

Recent Developments in Color－Translating Ullura－Violet Microscopy，R．B．Holt
Some Theoretical and Practical Aspects of Microscanning，W．E．Tollers，et al．
The Electrocardiophone－A New Surgical Tool，A．J．Morris and J．I．Swanson
Instrumentation for Spectral Phonocardi－ ography，George N．Webb

## Electron Tubes

A UHF Traveling－Wave Amplifier Tube Em－ ploying an Electrostatically Focused Hol－ low Beam，C．B．Crumly

Design of Solenoids for Traveling－Wave Tubes，J．E．Etter A．W．Friend and W．Watson
Light Weight Solenoids of Aluminum Fnil， W．G．Worcester and A．L．Weitzmam
The Serrodyne－A Single Sideband Synchro－ dyne，R．C．Cumming
Recenl Dark Trace Tube Developments， S．Nozick
Recent Developments in the Use of Dispenser Cathodes in Lowe and Medium Power Mag－ netrons，R．S．Briggs

## Microwave 「Techniques

Waveguides for Long Distance Communica－ tions，A．C．Beck
Recent Advances in Microwave Filter Tech． niques，Seymour Cohn
Geometrical Methods far the Analysis of Two－ Part Networks，G．A．Deschanips
Some Applications and Characteristics of Ferrite at Wavelenyths of 9.87 and 1.9 cms ， Clyde Stewart
Measurement and Control of Microwave Fire－ quencies by Lower Radio Frequencies，R．C． Mackey et al．

## Friday Afternoon

## Antennas II

Radiation Characteristics zoith Power Gains for Slots on a Sphere，Y．Mushiake and R．E．Webster
Rudiation Patterns of Asymmetrically Fed Prolate Spheroidal Antennas，H．A．Myers
Phase Properties of Antennas for the Dovap Millile Tracking System，T．Morita and C．W．Stecle
Rotationally Symmetric Dialectric Microwave Lenses with Two－Dimensional Wide Angle Scanning Characteristics，A．Mayer and E．Wantuch

## Radio Relay Systems I）esign

Design of FM Radio Relay Equipment for Multi－Channel Operation，J．W．Halina
Factors Affecting the Spacing of Radio Ter－ minal in an UHF Link，J．H．Gerks
Radio Communication with Secondary Power， H．E．Hollmann
Single Sideband Mulliplexing as it Applies to Microwave Relays，T．L．Leming

## International Researcii in Electronics and Allied Fields <br> II．The International Geo－ physical Year Program

The International Geophy－ical Year， 1957 1958，R．J．Slutz
Absorption Measurements During the Inter． national Geophysical Year，Gordon Little
Vertical Incidence Ionasphere Sounding Measurements during I．G．Y．，J．M．Watts Back－Scattering Measurements During
I．G．Y．，A．M．Peterson


## Professional Groups

Aeronautical \& Navigational ElectronicsChairman, Edgar A. Post, Navigational Aides, United Air Lines, Operations Base, Stapleton Field, Denver 7, Colo.
Antennas \& Propagation-Chairman, Delmer C. Ports, Jansky \& Bailey, 1339 Wisconsin Ave., N.W., Washington 7, D. C.
Audio-Chairman, IV. E. Kock, Bell Tel. Labs., Murray IIill, N J.
Automatic Control-Choirman, Robert B. Wilcox, Raytheon Manufacturing Co., 148 California St., Newton 58, Mass.
Broadcast \& Television Receivers-Chairman, W. P. Boothroyd, Philco Corp., I'hiladelphia 34, Pa.
Broadcast Transmission Systems-Chairman, O. IV. B. Reed, Jr., Jansky \& Bailey, 1735 DeSales St., N.W., Washington, D. C.

Circuit Theory-Chairman, J. Carlin, Microwave Res. Inst., Polytechnic Inst. of Brooklyn, 5.5 Johnson St., Brooklyn 1, N. Y.

Communications Systems-Chairman, A. C.

Peterson, Jr., Bell Labs., 463 West St., New York 14, N. Y.
Component Parts-Chairman, Floyd A. Paul, Reliability Bendix Development Lab., 116 W. Olive Avenue, Burbank, Calif.
Electron Devices-Chairman, George Espersen, Microwave Tube Section, Philips l.abs., Irvingtoti-on-Hudson, N. Y.

Electronic Computers-Chairman, J. H. Felker, Bell I.abs., Whippany, N. J.
Engineering Management-Chairman, C. J. Breitwieser, Lear, Inc., 3171 S. Bundy Drive, I.os Angeles 34, Calif.
Industrial Electronics-Chairman, George P. Bosomworth, Engrg. Lab., Firestone Tire \& Rubber Co., Akron 17, Ohio
Information Theory-Chairman, Louis A. DeRosa, Federal Telecommunications Lab., Inc., 500 Washington Avenue, Nutley, N. J.
Instrumentation-Chairman, Robert L. Sinck, Consolidated Engrg. Corp., 300 N. Sierra Madre Villa, Pasadena, Calif.

Medical Electronics-Chairman, Dr. Julia F. Herrick, Inst. of Experimental Medicine, Mayo Found., Rochester, Minn.
Microwave Theory and Techniques-Chairman, A. C. Beck, Bell Labs., 463 West St., New York 14, N. Y.
Nuclear Science-Chairman, Dr. Donald II. Loughridge, Dean of Engineering, Northwestern Tech. Inst., Evanston, Ill.
Reliability and Quality Control-Chairman, Leon Bass, Jet Engine Dept., General Elec. Co., Cincinnati 15, Ohio
Production Techniques-Chairman, R. R. Batcher, 240-02-42nd Ave., Donglaston, L. I., N. Y.

Telemetry and Remote Control-Chairman, C. II. IIoeppner, Stavid Engineering, Plainfield, N. J.
Ultrasonics Engineering-Chairman, M. 1). Fagen, Bell Labs., Whippany, N. J.
Electron Devices-Chairman, J. S. Saby, Electronics Laboratory, G. E. Co., Syracuse, N. Y.

## Sections*

Akron (4)-H. L. Flowers, 202919 St., Cuyahoga Falls, Ohio; H. F. Lanier, 49 W: Lowell St., Akron, Ohio
Alberta (8)-Officers to be elected.
Albuquerque-Los Alamos (7)-T. F. Marker, 313340 Street, Sandia Base, Albuquerque, N. Mex.; T. G. Banhs, Jr., 1124 Monroe Street, S.E., Albuquerque, N. Mex.
Atlanta (3)-D. L. Finn, School of Elec. Engr'g., Georgia Inst. of Tech., Atlanta, Ga.; P. C. Toole, 605 Morningside Dr., Marietta, Ga.
Baltimore (3)-C. D. Pierson, Jr., Broadview Apts. 1126, 1:6 West University Pkwy., Baltimore 10, Md.; M. I. Jacob, 1505 'rredegar Ave., Catonsville 28, Md
Bay of Quinte (8)-J. C. R. Punchard, Elec. Div., Northern Elec. Co. Ltd., Sydney St., Belleville, Ont., Canada; M. J. Waller, R.R. 1, Foxboro, Ont., Canada
Beaumont-Port Arthur (6)-I.. C. Stockard, 1390 Lucas I)r., Beaumont, Texas; John Petkorsek, Jr., 4390 El Paso Ave., Beaumont, Texas
Binghamton (1)-O. T. Ling, 100 Henry Street, Binghamton, N. Y.; Arthur Hamburgen, 102 S. Nanticoke Ave., Endicott, N. Y.

Boston (1)-A. J. Pote, M.I.T., Lincoln Lab., Room C-249-A, Box 73, Lexington 73, Mass.; T. P. Cheathanı, Jr., Hosmer St., Marlborough, Mass.
Buenos Aires-J. M. Rubio, Ayacucho 1147, Buenos Aires, Argentina; J. L. Blon, Transradio Internactional, San Martin 379, Buenos Aires, Argentina
Buffalo-Niagara (1)-D. P. Welch, 859 Highland Ave., Buffalo 23, N. Y.; William S. Holmes, 1861 Ellicot Rd., West Falls, N. Y.
Cedar Rapids (5)-Ernest Pappenfus, 1101 30 Street Dr., S.E., Cedar Rapids, Iowa;

[^62]E. L. Martin, 111923 St., S.E., Cedar Rapids, Iowa
Central Fiorida (3)-1 Ians Scharla-Nielsen, Radiation Inc., P.O. Drawer ' $Q$ ', Melbourne, Fla.; G. F. Anderson, Radiation Inc., P'O. Box 'Q', Melbourne, Fla.
Chicago (5)-J. J. Gershon, De Vry Tech. Inst., 4141 Belmont Ave., Chicago 41, Ill.; J. S. Brown, 9829 S. Hoyne ive., Chicago 43, IIl.
Cincinnati (4)-R. A. Maher, 6133 Sunridge Drive, Cincinnati 24, Ohio; W. S. Alberts, 6533 Elwynne Dr., Silverton, Cincinnati 36 , Ohio
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Connecticut Valley (1)-II. E. Rohloff, The Southern New England Tel. Co., 227 Church St., New Haven, Conn.; B. R. Kamens, 45 Brooklyn Circle, New Maven 15, Conn.
Dallas-Fort Worth (6)-J. A. Green, Box 7224, Dallas 9, Texas; G. K. Teal, Texas Instruments Inc., 6000 Lemmon Ave., Dallas 9, Texas
Dayton (5)-A. B. Henderson, 801 Hathaway Rd., Dayton 9, Ohio; N. A. Nelson, 310 Lewiston Rd., Dayton 9, Ohio
Denver (6)-R. E. Swanson, 1777 Kipling St., Denver 15, Colo. ; S. M. Bedford, Jr., Mountain States Tel. \& Tel., Room 802, Denver, Colo.
Des Moines-Ames (5)-A. A. Read, 511 Northwestern Ave., Anes, Iowa; W. L. Hughes, E. E. Dept., Iowa State College, Ames, Iowa
Detroit (4)-N. D. Saigeon, 1544 Grant, Lincoln Park 25, Mich.; A. L. Coates, 1022 E. Sixth St., Royal Oak, Mich.
Elmira-Corning (1)-J. L. Sheldon, 179

Dodge Ave., Corning, N. Y.; J. I'. Hocker, Corning Glass Works, Corning, $\mathcal{X}$. Y.
El Paso (6)-J. F. Stuart, Box 991, El l'aso, Texas; W. T. McGill, 7509 Mazatlan Rd., El Paso, 'lexas
Emporium (4)-L:. II. Boden, R.D. 1, Emporium, Pia.; I1. S. Hench, Jr., R.I). 2, Emporium, Pa.
Evansville-Owensboro (5)-A. P. Haase, 2230 St. James Ct., Owensboro, Ky.; D. D. Mickey, Jr., Eng'g. Dept., General Electric Co., Owensboro, Ky.
Fort Wayne (5)-J. J. Iffland, 2603 Merivale St., Kirkwood Park, Ft. Wayne 8, Inc.; T. L. Slater, 1916 Eileen Dr., Waynedale, Ind.
Hamilton (8)-G. M. Cox, $15 \pm$ Victoria St., S., Kitchener, Ont., Canada; A. I.. Fromanger, Box 507, Ancaster, Ont., Canada
Hawaii (7)-G. W. Clark, Box 193, Lanikai, Oahu, T. I.; J. R. Sanders, c/o Matson Navigation Co., Box 899, Honolulu, 'Г. H. Houston (6)-L. W. Erath, 2831 Post Oak Rd., Houston, Texas; J. M. Bricaud, Schlumberger Well Surveying Corp., P.O. Box 2175, Houston, Texas

Huntsville (3)-D. E. French, 1403 E. Clinton St., Huntsville, Ala.; T. L. Greenwood, 1709 LaGrande St., Iluntsville, Ala.
Indianapolis (5)-J. T. Watson, 407 N. Penn. 508, Indianapolis 4, Ind.; M. J. Arvin, 4329 Fletcher Ave., Indianapolis 3, Ind.
Inyokern (7)-G. D. Warr, 213-A Wasp Rd., China Lake, Calif; B. B. Jackson, 54-B Rowe St., China Lake, Calif.
Israel-Franz Ollendorf, Box 910, Hebrew Inst. of Tech., Haifa, Israel ; J. 1f. Hallserstein, P.O.B. 1, Kiriath Mozkin, Israel
Ithaca (1)-Ben Warriner, General Electric Co., Advanced Electronics Ctr., Cornell University, Ithaca, N. Y.; R. L. Wooley, 110 Cascadilla St., Ithaca, N. Y.
(Cont'd on next page)

Fections ront＇d）
ansas City（6）—Kenneth V．Newton，Ben－ dix Aviation Corp．，Box 1159，Ǩansas City 41，Mo．；Mrs．（i．L．Curtis，Radio Industries，Inc．， 1307 Central Ave．，Kan－ sas（ity 2．K゙an．
Little Rock（6）—J．E．II $\mathfrak{y}$ lie， 2701 N．Pierce， Little Rock，\rk．：Jim Spilman，A．R．\＆T＇．
 Rork，．Irk．
London（8）－C．F．Macl）onald， 328 St． James St．，London，Ont．，Canada；J．D．B． Moore． 27 MeClary Ase．，London，Ont．， Cimada
Long Island（2）－ 11 ．F．Batikey，Ilazeltinc Corp．，58－25 Little Neck l＇kway，Little Neck，I．．I．，N．Y．；I＇．（r．Ilansel，Addison Lane，Creenvale，L．I．，N．「．
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Toronto（8）－A．P．H．Barclay， 2 Pine Ridge l）r．，Toronto 13，Ont．，Canada； II．W．Jackson， 352 Laird Ir．，Toronto 17，Ont．，Canada
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## Subsections

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Wholey， 342 Verano Dir．，Los Altos，Calif．
Pasadena（7）－Officers to be elected．
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The following issues of "Transactions" have recently been published, and are now available from the Institute of Radio Engineers, Inc., 1 East 79th Street, New York 21, N. Y. at the following prices. The contents of each issue and, where available, abstracts of technical papers are given below.

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## AERONAUTICAL AND NAVIGATIONAL ELECTRONICS

Vol. ANE-2, No. 1, March, 1955

## University of Dayton Honored

Chairman's Report
Trajectory Precision Requirements for Automatic Landing-J. L. Ryerson

Requisite to landing an aircraft antomatically is the perception of its altitude. This reports attempts to resolve the theoretical accuracy of an aircraft's altitude by standard statistical equations for the propagation of error.

Available aircraft trajectory data have been employed to determine the bandwidth of aircraft trajectories during approach and "landing." Such data have also been incorporated to establish the bandwidth of the noise superimrosed by apparent radar target "wander."

A method of removing the noise by optimum filtering techniques, is discussed. A means of retrieving the high-frequency components of the aircraft trajectory information by inertial equipment is proposed. It is further proposed that similar smoothing techniques be utilized to remove low-frequency drift terms which appear as noise in the inertial equipment. The two sources of information are subsequently recombined to obtain broadband trajectory coverage having a greatly improved signal-tonoise ratio.

Airborne UHF Communication Equipment -G. H. Scheer

After more than two years' experience, it has been determined that airborne UHF communications is satisfactory. Many unforeseen types of interference have been found, some of them serious. Antenna patterns on aircraft are not ideal, but are usable. The newest subminiature airborne transceiver has undergone unique overational engineering tests. Results show that its design is highly desirable from the aspects of performance, installation and main-
tenance. Until the ultimate design of an equipment requiring no maintenance, industry should improve quality of components to increase reliability and reduce repairs.

Flight Director Design Trends-G. Iddiags and E. Martino

The flexibility of the fligltt director to include practically any all-weather mission is accomplished by the coupling of the basic navigation data to the central flight director computer with special signal-shaping circuits designed on a modular hasis. This paper describes the nature of these mission-coupling devices and an extension of their use to improve the basic flight panel information.

Cooling Requirement Charts for Electronic Equipment-L.J. Lyons

Efficient installation of airborne electronic equipment requires adequate data to design a heat removal system for the airplane. For many years aircraft generators have been supplied with rating charts to define their cooling requirements: electronic equipment now requires similar treatment. The form of the chart depends on the nature of the equipment, but consists, essentially, of a notograph showing required cooling airflow and bressure drop vs air temperature. When pertinent, it should also include such factors as life, output, surrounding wall temperature, etc. Such charts make it convenient to design and compare performance of ram and fan driven cooling systems, and to compare different equipment designs in terms of their flow and pressure drop requirements. Examples of several different types of charts, their preparation and their use, are presented.

The Accuracy of the VHF Omni-Range System of Aircraft Navigation: A Statistical Study -W. G. Anderson

This paper describes a statistical treatment of the errors encountered in the VHF Visual Omni-Range Navigation System. The first part of the paper consists of a rudimentary discussion of statistical theory which serves to acquaint the reader with the methods used in the
second portion of the report
Each of ten different errors are described in terms of a normal distribution function with the parameters $\bar{x}$ and $s$, the mean and standard deviation respectively. The ten errors are summed up and are compared to the error distribution which was obtained from error data gathered by the pilots of various scheduled airlines. The close agreement justifies the method used.

Of the ten errors, only five are significant. and the V(OR ground station error is shown to be greater than the sum of the other nine errors.

Effect of Internal Fluctuations and Scanning on Clutter Attenuation in MTI Radar-G. S. Grisetti, M. M. Santa and G. M. Kirkpatrick

The approximate effect of internal fluctuations and scanning on clutter attenuation in MII radar is derived and the results presented on two charts. It is proposed that derivative antenna patterns be used to increase the clutter attenuation when scanning. The theoretical improvement from the use of derivatives is appreciable, as shown on the second chart.

## AUDIO

Vol. Al'-3, No. 3, May-June, 1955
Complexity and Unreliability in Electronic Equipments-G. H. Scheer

Connecting Piezoelectric Pickups to Mag-netic-Pickup Amplifiers-B. B. Bauer

IRE-PGA Election and Convention Summary.

Bereskin New Editor for IRE TRANSACTIONS on AUDIO

Formation of Syracuse IRE-PGA Chapter -W. W. Dean

Dayton IRE Section Organizes Audio Chap-ter-A. B. Henderson

PGA Chapter NEWS
Perceptibility of Flutter in Speech and Music-F. A. Comerci

The perceptibility of flutter at various rates, in recordings of speech and music, was investigated in relation to the development of a flutter meter which will give a direct indication of the effect upon programs as juclged by listeners.

Sound Measurements at Very High Levels - Arnold Peterson

The behavior of a number of microphones at high sound levels is described. Some of the problems encountered in making measurements at high sound levels are discussed.

Electronic Organ Tone RadiationD. W. Martin

The principles of design for electronic organ tone chambers are outlined. The differences between the design goals for loudspeaker enclosures for organs and for other purposes are explained in fundamental terms. The construction of new organ tone cabinets for indirect radiation is described in detail. A few organ installation examples are given.

PGA Institutional Listings

## COMPONENT PARTS

$$
\text { PGCP-3, March, } 1955
$$

## Report from the Chairman

The Effective Leakage Resistance of Several Types of Capacitors-R. W. Tucker and S. D. Breskend

A rate-of-charge method for measuring the effective leakage resistance of good quality, ligh valued capacitors is described. This method yields results rapidly and directly. The effective leakage resistance of various tybes of capacitors as a function of time of applied
voltage at different tenıperatures was measured. A method for calculating the change in capacitance with time of applied voltage is given. A polytetrafluoroethylene capacitor had the best direct-current properties of any of the types tested.

Transformer Design Chart-Reuben Lee and N. E. Mullinix

This naper describes a transformer design chart by which the design of two winding, sixty cycle, low voltage transformers can be made without most of the time-consuming design procedures. Its use is intended for transformer design engineers; therefore, it provides only the winding information that varies for each design. A specific series of cores and operating conditions are assumed. The turns, wire size and approximate winding resistance of both the primary and secondary can be determined from voltage requirements and secondary volt ampere rating. Equations upon which the chart is based are given together with an example problem. A set of rules for applying the chart to other transformer designs is also given.

Problems Encountered and Procedures for Obtaining Short-Term Life Ratings on Resis-tors-W. T. Sackett, Jr.

The purpose of this paper is to give an abbreviated review of Battelle Institute's activities in the electronic-component field, and to give a more detailed discussion of the particular phase of those activities having to do with the development of procedures for obtaining short-term life ratings on components.

Subminiature Transformers and Their Application to Junction-Transistor CircuitsE. F. Dunkin and D. L. Johnson

The technical limitations of subminiature audio or control frequency transformers or inductors are discussed and some features of their design described. Statistical methods are used to control quality and specification limits as the physical sizes are too small to allow for individual adjustment. A miniature toroidalshell construction has given results comparing favorably with laminated assemblies below a certain size and in this form, optically-finished lapped joints are employed in the magnetic circuit. The use of subminiature core assemblies as transductors has also been investigated.

## ELECTRONIC COMPUTERS

## Vol. EC-4. No. 2, June, 1955

PGEC Student Activities and Education in Computers-H. H. Goode

A Survey of Electronic Analog Computer Installations-L. B. Wade and A. W. Wortham

A survey has been made of real-time electronic analog computer (differential analyzer) installations. This survey was conducted so that a directory of the installations could be compiled and so that various data regarding the installations could be made available for analysis. The survey was conducted by a mail questionnaire. Information was obtained regarding size of installation, size of staff, weekly usage of the equipment, age of installation, and availability to outside organizations from 96 installations having a total of 8,320 computer amplifiers. The results of the survey have been analyzed and are presented in this paper, together with the directory.

A Digital Computer for Use in an Operational Flight Trainer-W. H. Dunn, C. Eldert, and P. V. Levonian

The requirements for a digital computer for use in an operational flight trainer are presented with emphasis being placed on the realtime aspects of the problem. The general purpose digital computer is shown to be inadequate for this purpose and a special purpose digital computer is described which meets the requirements.

A Diode Multiplexer for Analog VoltagesH. J. Gray, Jr., M. Rubinoff and J. Tompkins A diode multiplexer switch is clescribed for time-sharing 64 analog voltages in a digital computer application. Apart from its relative simplicity and economy, the multiplexer characteristics of microsecond switching speeds, maximum settling time of 133 microseconds for a 10 -volt operating range, and accuracies of better than 1 per cent full scale are confirmed both by theoretical equations and by experimental results.

Some Notes on Logical Binary Counters-

## R. M. Brown

The properties of binary counters which utilize non-transient storage elements for the count information are presented. The four possible sets of logical connections between the two storage elements necessary for each stage are described. The binary numbers represented in the storage elements are shown to be the actual count in one set of elements and in the other set a Gray code representation of twice the actual count. Examples of bi-directional counters are given.

A Variable Binary Scaler-D. B. Murray
The binary elements of a counter or "scaler" may be interconnected in many ways. This paper discusses a class of interconnections in which some elements are "forward-counting" and some are "reverse-counting." By changing the interconnections any arbitrary integral scaling ratio (up to the counter capacity) may be obtained.

Time-Delay Circuits-W. E. Thomson
Contributors
News: S. B. Disson
Reviews: Reviews of Current LiteratureH. D. Husky, ed.

## ENGINEERING MANAGEMENT

## PGEM-3, March, 1955

More Engineering Per Dollar-Burgess Dempster

Some Factors Related to Management of an Applied Research Project-Harley Iams

Among the factors important to the management of an applied research project are the arrangement of working areas, the provision of labor-saving supporting equipment, and the supplying of adequate shops for making experimental parts. But even more important is the management of the people, including their organization into effective teams, seeing that engineers and scientists are retained for their best contribution to national defense, and upgrading the abilities of the staff.

A Practical Approach Toward Integration of Project and Group Theories in Establishing an Engineering Organization-C. F. Horne

Engineering Management and the Changing World-Maurice Nelles

Are Engineers People?-A. M. Zarem
Biographical Notes on the Authors

## INDUSTRIAL ELECTRONICS

## PGIE-2, MARCh, 1955

Mutual Problems in Industrial Electronics and Communications-E. W. Allen

Numerical Control of Machine ToolsLeroy U. C. Kelling

Industry needs more flexible methods of programming machine cycles to achieve automatic operation of machine tools in limitedquantity production. Part of this need is met by numerical control systems which command the machine in accordance with prepared numerical instructions read from a storage medium. Such numerical instruction can be stored in binary or decimal numerical form on punched telegraph tape, punched tabulating cards, magnetic tape, and many other data
storage mediums. Numerical control systems are well adapted to control of machine tools such as lathes, turret punch oresses and boring, drilling and milling machimes. The numerical control system for a Wiedemann turret punch press is a typical example illustrating the problems of joining controls and machines into a smoothly working combination.

Electromechanically Stabilized DC Amplifier for Use in Transducing and Telemetering of Milli-Voltage and Micro-Ampere SignalsHubert A. Riester, Jr.

Industrial Applications of X-Ray Tech-niques-T. H. Rogers

The Application of Radioactivity to Measurement and Control-Norman E. Walters

A Magnetic Thickness Gage for Rubber and Plastic Applications-Albert M. Dexter

Advantages of Electronic Process ControlC. E. Mathewson

Contrary to measurement requirements, the modern concept of closed loop, or feedback, control has created an urgent requirement for high sjeed response. The pure dead times and variable exponential lags characteristic of bneumatic communication systems have limited control performance to an appreciable extent. Electrical communication is the obvious solution and equipment designed for its utilization will be described herewith.

Electronically Produced and Controlled Illumination-Harold E. Edgerton

After a brief summary of the theory of clectronically produced pulses of light, energy storage systems, and energy converting devices, a review is made of commonly used triggering and controlling elements. A series of practical devices utilizing the previous:y described methods are discussed, showing typical stroboscopes, high-speed single-flash photographic lights, high-speed motion-picture lights, and flash sequence equipment.

Automatic Detection of "Green-Rot" in Shell Eggs-K. H. Norris

Power Oscillators for Dielectric HeatingT. L. Wilson

Survey of Today's Use of Power Rectifiers in Industry-L. W. Morton

Electronic Considerations in the Theory and Design of Electric Spark Machine ToolsE. M. Williams and J. B. Woodford, Jr.

## INFORMATION THEORY

Vol. IT-1, No. 1, MARCH, 1955
An Analysis of the Detection of Repeated Signals in Noise by Binary Integration-J. V. Harrington

An analysis of the detection of repetitive signals in noise by binary integration techniques is made. An expression for the effective signal-to-noise ratio of the quantized video is obtained and is shown to apply to any halfwave second detector. A comparison of analog and digital integration is made, and it is further shown that digital integration is, at most, 1.9 db poorer due to the quantization loss. However, the loss due to nonideal analog integration can make the two types equivalent. The optimum settings for quantizer and counter thresholds are derived, and expressions for the finaldetection and false-alarm probabilities are determined. Lastly, the results are modified to include the effect of nonuniform amplitudes in the set of signals being quantized and integrated.

An Expansionfor Some Second-Order Probability Distributions and its Applications to Noise Problems-J. F. Barrett and D. G. Lampard

In this paper it is shown that, in general, second-order probability distributions may be expanded in a certain double series involving orthogonal polynomials asscociated with the corresponding first-order probability distributions. Attention is restricted to those second-
order probability distributions which lead to a "diagonal" form for this expansion.

When such distributions are joint probability distributions for samples taken from a pair of time series, some interesting results can be demonstrated. For example, it is shown that if one of the time series undergoes an amplitude distortion in a time-varying "instantaneous" nonlinear device, the covariance function after distortion is simply proportional to that before distortion.

Some simple results concerning conditional expectations are given and an extension of a theorem, due to Doob, on stationary Markov process is presented.

The relation between the "diagonal" expansion used in this paper and the Mercer expansion of the kernel of a certain linear homogeneous integral equation, is pointed out and in conclusion explicit expansions are given for three specific examples.

Predictive Coding-Peter Elias
Predictive coding is a procedure for transmitting messages which are sequences of magnitudes. In this coding method, the transmitter and the receiver store past message terms, and from them estimate the value of the next message term. The transmitter transmits, not the message term, but the difference between it and its predicted value. At the receiver this error term is added to the receiver prediction to reproduce the message term. This procedure is defined and messages, prediction, entropy and ideal coding are discussed to provide a basis for Part II, which will give the mathematical criterion for the best predictor for use in the predictive coding of particular messages, will give examples of such messages, and will show that the error term which is transmitted in predictive coding may always be coded efficiently.

The Linear, Input-Controlled, Variable-Pass Network-B. E. Keiser

This paper describes the study and development of a linear, variable-pass network system which is controlled by the Fano short-time autocorrelation function of the input. Given an input function, the message, whose short-time power spectrum varies in an unpredictable manner with time, and to which there has been added a different function, the disturbance, whose short-time power spectrum is either time-invariant or varies in a completely known manner, a linear, input-controlled, variablepass network can be specified which minimizes the mean-square error between the message input and the total output, taking network delay into account. Methods for mathematical computation of the mean-square error have been devised.

The linear, input-controlled, variable-pass network has been found to have a lower meansquare error than that attainable with an optimum-mean-square, linear, fixed, selective network, for certain types of input messages.

Spectral Density Functions in Pulse Time Modulation-H. Kaufman and E. H. King

Spectral power density functions corresponding to various types of pulse shapes, probability distribution functions arising in the study of pulse time modulation problems are computed. The results are presented in tabular form. The following cases are considered: PAM and PPM, for arbitrary pulse shape, PDM, for rectangular, Gaussian, and error-function pulse shapes, and SEM, for rectangular pulse shape.

A Note on the Sampling Theorem-L. J. Fogel

The human operator often perceives rate as well as amplitude information in sampling various displayed continuous parameters. It is therefore necessary to extend the Sampling Theorem to allow the analysis of certain manmachine relations. The result is stated and the required mathematics included in the appendix. Certain distinct problem areas where this extension can be fruitfully employed are indicated.

Statistical Calculation of Word Entropies for Four Western Languages-G. A. Barnard

Using a modified version of Shannon's method, comparative figures for the word-letter entropies of printed English, French, German, and Spanish are obtained and the method described.

Papers presented at WESCON in Los Angeles, August 25-27, 1954

On the Modulation Levels in a Frequency Multiplexed Communication System by Statistical Methods-R. L. Brock and R. C. McCarty.

This paper presents a mathematical analysis with experimental verification of the distribution of the instantaneous voltage of a complex signal resulting from the combination at random of a small number $n$ of sinusoidal oscillations. The resulting calculated distributions are plotted in the form of a set of probability curves for comparison with curves obtained by experiment. Further laboratory measurements in which the individual sinusoidal oscillators are frequency-modulated in a manner suitable for communicating information in a binary form, yield substantially no change in the amplitude distribution as determined for the unmodulated oscillators. Consequently, the results of the mathematical analysis may be applied in the determination of $M$, the degree to which each subcarrier may amplitude modulate a final carrier in an fm-am frequency multiplexed system. $M$ may be determined for any desired degree of overmodulation in excess of one per cent and for as many subcarriers as are required in the system. Modulation levels determined according to approximate methods and the method described here are tabulated and compared.

On the Response of a Certain Class of Systems to Random Inputs-Jack Heilfron

This paper deals with the connection between vector Markoff processes and the response of a lumped constant parameter linear system composed of a finite number of elements. It was known that if a Gaussian process which is one component of a vector Markoff process passes through such a system, the result is also Gaussian and may be considered as one component of a higher dimensional vector Markoff process. We show that the term Gaussian may be excluded in the above statement. The practical importance of this result is that if one can determine the initial and transition probabilities of this vector Markoff process, one can also determine the complete statistical properties of the output of the system. This further implies that the determination of the properties of the output for the class of not necessarily Gaussian inputs mentioned above is not as difficult as might be expected from the results for just the first probability distribution for non-Gaussian inputs.

Noise in Driven Systems-J, M. Richardson

It is known that a direct relation exists between the noise in a system in equilibrium and transient drift toward equilibrium. It seems that a similar relation should exist for a system in a nonequilibrium stationary state. It is now necessary to distinguish between two types of transients; those produced by selecting those systems satisfying certain initial and those produced by actual physical perturbation. It is is shown that a simple relation exists between noise and the transients produced by selection and that no relation exists in the case of transients produced by perturbation. In the equilibrium case it is shown that the two types of transients, though still logically and operationally distinct, can be described by the same impedance operator.

Design and Performance of Phase-Lock Circuits Capable of Near-Optimum Performance Over a Wide Range of Input Signal and Noise Levels-R. Jaffe and E. Rechtin

## MICROWAVE THEORY AND TECHNIQUES

Vol. MTT-3, No. 3, Apleil, 1955
Editorial Comment by Saad
Frontispiece of W. W. Mumford
Editorial-W. W. Mumford
Advances in Microwave Theory and Tech-niques-D. D. King

Planar Transmission Lines-David Park
This paper derives formulas for the transmission properties-characteristic impedance and attenuation-in the principal noode of a transmission line consisting of one or two long strips of metal foil embedderl in a dielectric material between two long metal strips considerably wider than the central ones. The width and spacing of the central strips is arbitrary, and it is also necessary to take account of their thickness in computing the attenuation. A graphical method is given for evaluating the characteristic impedance in general, and analytic approximations are given for a number of special cases. Finally the question of the leakage of power from between the outer strips is considered briefly.

Measurement of Time-Quadrature Components of Microwave Signals-J. H. Richmond

A phase-sensitive coherent detector used for microwave laboratory measurements is described. The receiver measures the real $(|E| \cos \alpha)$ and imaginary $(|E| \sin \alpha)$ components of a signal $E$ with equipment which is less elaborate than that required for measuring the amplitude $|E|$ and phase $\alpha$. Furthermore, many calculations are more convenient if $E$ is presented in rectangular rather than polar form.

Measurements made with the receiver on known fields in waveguides are included to demonstrate its accuracy. The receiver has a sensitivity of -125 dbw at $9,375 \mathrm{mc}$.

Optimum Design of Stepped TransmissionLine Transformers-S. D. Cohn

This paper describes the optimuin stepped-transmission-line transformer structure for matching two unequal characteristic impedances. For any specified bandwidth, the steps are designed to yield a Tchebycheff-type (or equal-ripple) reflection-coefficient response. Over this band, the maximum vswr is less than that obtainable with any other stepped-transformer having the same number of steps. The design method and the technique for eliminating discontinuity-capacitance effects are given. The measured results for a coaxial and a waveguide model are presented and found to verify the method.

The Use of Scattering Matrixes in Microwave Circuits-E. W. Matthews

Difficulties arising from the use of the impedance concept in microwave circuitry have led to the introduction of the scattering representation for work at these frequencies. This paper presents a development of the scattering approach in terms of fundamental transmissionline phenomena. The physical meaning of the quantities involved is brought out wherever possible and the relationships among the various elements of the scattering matrix are given. Several examples of the application of scattering techniques to analysis of the properties of microwave junctions are presented, and methods for measuring scattering parameters of such junctions are outlined.

Some Applications and Characteristics of Ferrite at Wavelengths of 0.87 Cm and 1.9 Cms-Clyde Stewart

This paper describes the use of ferrites in waveguide to produce Faraday rotations at 0.87 cms and 1.9 cms wavelengths. The Dicke radiometric receiver is briefly reviewed and its improvement by the use of ferrite waveguide components is described. Experimental equipment for securing data on the behavior of
ferrites is discussed. Details are given for the construction of a midirectional waveguide transmission line for 0.87 cms wavelength.

Probes for Microwave Near-Field Meas-urements-J. H. Richmond and T. E. Tice

To be satisfactory for microwatve near-field measurements, a probe must have desirable polarization characteristics, must have an aperture small enough to indicate the field at a point, must deliver sufficient signal voltage to permit accurate measurement, and yet inust not seriously distort the fields. The rlesign of a probe may be simplified if the fields to be measured are known to be almost linearly polarized or to consist only of a traveling wave. Comparison of measurements made with various probes has led to the development of a small open-ended waveguide probe which is simple to construct and has given excellent results.

Measurement Techniques for Multimode Waveguides-A. C. Beck

This paper surveys some of the technipues that have been worked out for multimode waveguide measurements. Equipment has been developed for measuring one mode at a time by taking advantage of the differences between the modes. Illustrations of its use are given.

> Reports of International Organizations- L. G. Cumming and W. W. Mumford

Addenda to "Bibliography on Directional Couplers"-R. G. Medhurst and R. F. Schwartz

## MICROWAVE THEORY AND TECHNIQUES

## Vol. MTT'-3, No. 2, Marcil, 1955

Microwave Printed Circuit-An Historical Survey-R. M. Barrett

The microwave lirinted circuit, as described in this pager, is an extension of the well-known technique which is of such importance in the lower frequency regions, where lumped element circuits are practical. This new circuit possesses all of the virtues of other frinted circuits, such as light weight, cheapmess, ease of manufacture, miniaturization, etc., along with the ability to be used at frequencies as higln as $10,000 \mathrm{mc}$. The basis of the new technique is the planar or "flat strip" conxial transmission system which was developed during Wiond War II but which has remained unpublished and relatively unknown in the postwar period; and for which an adequate theoretical analysis had mot been available.

Microwave Strip Circuit Research at Tufts College-A. D. Frost and C. R. Mingins

The research work on microwave strip circuits which has been in progress since November, 1952 is describecl. Experimental investigations have included measurement of the characteristic impedance of various lines, the design of transitions from coaxial connectors to strip)lime, the preparation and adjustment of matched resistive terminations, and most recently the evaluation of the effects of simple line discontinuities such as bends or steps. Theoretical studies on the calculation of characteristic impedance and line loss have also been carried out.

Characteristics and Some Applications of Stripline Components-W. E. Fromm

Basic characteristics of Stripline in various frequency bands from 1000 to $16,000 \mathrm{mc}$ are summarized. Various components such as transitions to coaxial line, attenuators, hybrid rings, directional couplers, and filters are shown. Some applications of these components in practical high performance microwave circuits and equipment in the frequency range of $2500-10,000 \mathrm{mc}$ are also described.

Photoetched Microwave Transmission Lines-Norman R. Wild

Microwave transmission line and components of unusual light weight and compact con-
struction can be made employing photoetcling techniques to produce strip type transmission line. This report will be a general description of work done at Sauders Associates, Inc., to develop techniques for the design and manufacture of photoetched microwave transmission lines. Discussion will include measurements of attenuation and radiation leakage on parallel plate strip lines, as well as shielded type Triplate lines, the problem of mode purity and its relation to electrical parameters, various schemes of making transitions from standard waveguide to photoetched strip line. The basic design and ierformance of various components, as well as items of test equipment, such as slotted lines, matched loads, fixed attenuators, variable attenuators, directional couplers, crystal holders, phase shifters, hybrid rings, coax to Iri-plate transistors, etc., will also be treated. In addition, data will be presented showing impedance and susceptance values of simple discontinuities and impedance matching transformers. A simple technique for constructing gyrators and resonators will be presented, and the design and fabrication of an S-band signal generator employing photoctched microwave Tri-plate line will be shown, illustrating that practical inicrowave systems can be constructed far more economically than would be possible utilizing conventional waveguide techniques.

Characteristics and Applications of Microstrip for Microwave Wiring-M. Arditi

The experimental results of the transmission propertirs of Microstrip are compared with the values to be expected from a first order theory based on the assumption of a T.E.M. mode of propagation. The characteristics of various Microstrip components are given. These components include: waveguide or coaxial transducers, hybrids, directional couplers, crystal mixers, attenuators, filters, ferrite modulators, gas discharge modulator tubes and a wideband noise source. The design considerations stress the wide-band properties of Microstrip similar to those found in coaxial lines.

The methods of measurement used in Microstrip are outlined and they show the simplicity of the experimental set-up required for the application of Deschamps' method for determining the principal characteristics.

The applications of Microstrip to the design of complete systems such as microwave re"civers are discussed and examples in " S " band, " C " band and " X " band are given.

Miniature Strip Transmission Line for Microwave Applications-E. N. Torgow and J. W. E. Griesmann

The construction of a strip line whose physical size is kept as small as possible consistent with reasonable electrical performance is presented. This line is fabricated by relatively simple techniques and can be shaped to fit line components into relatively confined spaces. The line has good power handling capacity and moderately low attentuation. Various components have been developed in this line, including a broadband $\frac{3}{8}$-inch coaxial line to strip line adapter, a broadband matched load, attenuators, and high and low pass filters.

Strip Type Components for 2000 Mc Receiver Head End-K. E. Zubulin

Recent experimental work has evolved some components using air-spaced strip type transtnission line that have been used successfully in connection with a variable attenuator, cavity and crystal mixer. Bandwidth, VSIVR, and NF measurements are comparable with a commercial receiver head-end presently in use. The asymmetric air-spaced strip-above-ground transmission line used results in a simple configuration for coupling the line to the cavity. It also facilitates the application of a variable attenuator using a ferrite slab of high attenuation per unit length with good VSWR properties.

Properties of Dielectric Image LinesD. D. King

The properties of a dielectric rod on an
image surface are revieved, and experimental results on straight sections, various bends, and a twist are presented. Techniques for measuring insertion parameters and field distributions are clescribed.

Practical Dielectric-Filled Waveguide-

## T. N. Anderson

This paper describes the develonment of a laminated teflon-filled dielectric waveguide using techniques similar to what has been done in the flexible coaxial line. This paper describes the development of dielectric waveguide giving the theoretical design of the teflon-filled dielectric waveguide from both a mechanical and electrical point of view.

The emphasis on this dielectric waveguide development has been to arrive at a practical waveguide construction whict would be suitable for a radar systems application. This dielectric waveguide is intended to provide a miniaturized waveguide circuit which will have essentially the same peak power handling capabilities as standard waveguide which would be suitable for use up to $200^{\circ} \mathrm{C}$.

The fabrication technique is described along with a description of the measurement procedure for determining the characteristics of this dielectric waveguide including match, attenuation and high power breakdown.

The design of special transitions from airfilled waveguide to dielectric-filled waveguide are described also.

This work was pertormed under contract number AF33 (600) 26763 for Wright Air Development Center and is intended to eventually yield a series of dielectric-filled waveguides, both rigid and flexible.

Measurement of Attenuation and Phase Velocity of Various Laminate Materials at L-Band-M. E. Rigenbach and H. W. Cooper

Measured data are plotted for the characteristic impedance, velocity of propagation, and attenuation of dielectric sheet supported strip transmission lines for four dielectric materials: Teflon bonded glass cloth, epoxy bonded glass cloth, polyester bonded glass mat, and XXXP paper base phenolic. At 1000 megacycles, the teflon material is excellent and the epoxy and polyester materials satisfactory for low $Q$ applications, such as microwave transmission lines.

The equivalent physical length of a dielectric sheet supported strip transmission line right angle is reported.

The Input and Muttal Impedance of Dipole Strips Between Parallel Planes-W.H. Hayt, Jr.

A center-fed filamentary dipole is parallel to and between two parallel, infinite, perfectly conducting planes and carries a sinusoidal current. The longitudinal electric field intensity corresponding to such a current distribution is then obtained by an application of the image principle to the field of a single center-fed filanentary dipole in free space.

This longitudinal electric field is then used directly to obtain the input impedance and mutual impedance between filamentary dipoles of resonant lengths by the induced emf method. The impedances appear as an infinite series of integrals which are approximated by simple expressions having errors of less than onequarter of one per cent. Curves are obtained giving the input impedance of dipoles having various resonant lengths and locations between the guard planes, and for several separations of the guard planes, the latter value being maintained less than one-half wavelength to avoid any propagating modes, Mutual impedance is shown as a function of dipole separations as well.

The results are then extended to dipoles having a cross-section which is a circle or a zero-thickness strip. Curves are obtained for the input impedance of strips between parallel planes for several plane separations and several dipole widths, as a function of dipole length.

Problems in Strip Transmission LinesS. B. Cohn

A review is given of characteristic-impedance formulas for shielded-strip transmission lines. From these formulas, a set of approximate relationships for the attenuation and $Q$ of a dielectric-filled shielded-strip transmission line is derived. The method makes the standard assumption that the current distribution is that of a lossless line and the surface resistivity that of an infinite-plane conductor. Although this method applies accurately to most other types of lines, in this case, an error of the order of $10 \%$ is believed to occur due to the failure of the assumptions at the corners of the strip. However, the error is in a direction that makes the computed values conservative, and the accuracy should be sufficient for most practical purposes. The derivation of a correction term is now being attempted.

In addition to the discussion of attenuation, attention is given in this paper to the design considerations involved in a shielded-strip-line impedance meter, and to some preliminary data obtained with this device. Also, the future topics for investigation under this research and development program are mentioned.

Equivalent Circuits for Discontinuities in Balanced Strip Transmission Line-A. A. Oliner

Theoretical formulas are derived for the equivalent circuit parameters of a variety of discontinuities in balanced strip transmission line. These formulas are simple in form and are obtained by employing a small aperture procedure or a Babinet equivalence procedure in conjunction with an approximate model of the line. The results for a number of discontinuities are presented and comparison is made with the available measured data.

A Universal Approximate Formula for Characteristic Impedance of Strip Transmission Lines with Rectangular Inner ConductorsR. L. Pease and C. R. Mingins

An explicit expression is developed for the characteristic impedance of a microwave strip transmission line with rectangular inner conductor of arbitrary dimensions. The expression is exact for zero thickness and arbitrary width, exact for zero width and arbitrary thickness, and quite accurate (within $3 \%$ for the extreme case of a square inner conductor of dimensions about 0.01 of plate separation, but in most cases of practical interest, within $0.1 \%$ ) throughout the entire range of thickness and width.

Stripline Radiators-E. G. Fubini
Progress on the use of strip conductors as microwave antennas indicates that the technique is flexible and economical. Broadside curtains can be fabricated with sufficient accuracy. Several types of balanced Stripline feed have been considered, and twists have been successfully built. A variety of baluns have been evaluated and used to feed colinear Franklin arrays through binary splits.

Slot Array Employing Photoetched TriPlate Transmission Lines-D. J. Sommers

Microwave printed circuit techniques are readily adapted to the construction of compact antennas ideal for flush mounting on high speed aircraft. This paper describes the development of a two-dimensional X-band array consisting of 16 slots fed by photoetched Tri-plate transmission line. The design of a unity coupled series slot and the resulting mode purity problems are discussed. Several power divider configurations are illustrated and data on the performance of some of these devices is presented. The construction of a 4 slot E-plane, a 4 slot H-plane and the combination $4 \times 4 \mathrm{E}-\mathrm{H}$ plane array utilizing these power dividers is shown. Radiation patterns of each of these arrays were measured and a comparison of the individual and combination array patterns is made.

Bandpass Filters Using Stripline Tech-niques-D. R. White and E. H. Bradley

Strip lines provide a convenient transmis-
sion medium for the realization of microwave filters. Since bandpass filters designed in waveguide and coaxial lines would be large at ultrahigh frequencies, strip lines afford a practical means of realizing filters which are simply fabricated, are readily reproduced, and, in most cases, represent an appreciable savings in size and weight. Of the different types of strip transmission lines currently in use, the so-called "sandwich" structure has been employed at Melpar for two reasons: (1) very broad-band coax-to-strip-line transitions are easily realized; and (2) the electromagnetic field is essentially confined between the two ground planes, thus reducing problems in packaging.

Using design techniques developed for direct-coupled cavity-type waveguide and coaxial filters, experimental strip-line filters having ten per cent bandwidths in the u-h-f spectrum have been developed. These units have less than 1 db mid-band insertion loss and provide a rejection of greater than 40 db at frequencies twelve per cent from the center frequency. The design techniques discussed in this paper are general and, therefore, are not restricted to the realization of the above filter characteristics. Some limitations pertaining to the realizability of the cavity parameters in different dielectric media and the existance of spurious responses are discussed.

Resonator and Preselector in StriplineJ. F. Moore and Max Michelson

One of Raytheon's commercial applications involves microwave circuitry in balanced strip line, with a $\frac{3}{8}$-inch spacing between ground planes. Though the less critical parts of the unit are etched in copper-clad Teflon-Fibreglas, two of the components are of higher $Q$ than can reliably be obtained in the presence of plastic. They are: (1) A resonator, for use as a frequency stabilizing reference element in an AFC circuit; and (2) A four-stage maximally flat preselector filter. These units are made of separate strips of metal, and do not depend on the plastic sheet for their support.

The unit is intended to operate over a $4 \frac{1}{2} \%$ band centered on 6725 mc , and was designed to avoid expensive parts and assemblies.

The design requirements were for a resonator with an unloaded $Q$ of 1980, and with no more than $\pm 0.45 \mathrm{mc}$ frequency variation over operational extremes of temperature and humidity; and for a preselector with less than 3.5 db . loss and a 30 mc pass band. Both units satisfy the overall electrical and mechanical design without requiring special high-cost structures. In fact, the microwave head is, in balanced strip line, about half as expensive as in conventional plombing. In addition, the present cost will be further reduced as larger quantities are considered.

Broad-Band Microstrip Crystal Mixer with Integral DC Return-Eric Carlson

A light and compact microwave mixer using microstrip has been designed for use in airborne equipment. The mixer features a low-input voltage standing wave ratio over a frequency range of one octave. The local oscillator is coupled to the input transmission line by a microstrip directional coupler having an integral $\mathrm{d}-\mathrm{c}$ return. Minimum coupling to the $\mathrm{d}-\mathrm{c}$ return is obtained by placing it in a region of low field intensity.

## RELIABILITY AND QUALITY CONTROL

## PGRQC-5, April, 1955

Statistical Design-A Means to Better Products of Lower Cost-R. C. Miles

Present evidence suggests that electronic reliability problems are being attacked with too little emphasis on the basic problem, which is one of equipment rather than merely component reliability. Among the reasons for this
situation are: lack of enforcement of equipment reliability requirements, resulting in part from difficulty of enforcement; vague or unrealistic statement of equipment reliability requirements; lack of an adequate quantitative basis for predicting the reliability of a proposed equipment design.

The popular concept of "guaranteed" reliability is basically a fallacy, since reliability cannot be positively guaranteed in any useful sense. If the guarantee concept were valid, an equipment using only "guaranteed" components should have a life at least equal to the shortest guaranteed component life; that such is not the case in practice proves the fallacy of reliability "guarantees."

A more practical concept of reliability involves the formulation of a statistical definition. For example, reliability may be defined as "the probability that a component or equipment will operate satisfactorily under given circumstances," time of operation being included as one of the "circumstances."

Although such a definition makes it possible to relate equipment reliability to the reliabilities of the individual components, presently available component data is not suitable for the purpose. In particular, most component reliability data is deficient as regards the variation of reliability with operating time. Even such component data as does exist applies to operating times at which the component reliability has become intolerably poor in terms of the requirements of equipment of even moderate complexity.

It appears further that the majority of existing component improvement programs will do little to improve the situation. Such programs seem to have been conceived on an unrealistic basis, concentrating on improving reliability at operating times near the end of the component useful life, rather than maintaining very high values of component reliability for as long as possible in the interests of improved equipment reliability.

Substantial progress toward more reliable electronic equipment requires a combination of:

1. A realistic quantitative basis for equipment reliability requirements.
2. Education of equipment and component designers and users as to the basic nature of the reliability problem.
3. Better data on the reliability of existing components, in addition to development of in!proved components.
4. Closer cooperation between component and equipment engineers, in order that each may acquire a better understanding of the other's needs and problems.

Contributing Factors to Component Parts Reliability and Extended Service-J. A. Goetz

Reliable electronic equipment performance depends fundamentally upon a sound application of engineering data and service information. Essential elements of a program tailored to this need by the equipment mantfacturer include:

1. Maintenance of a realistic field evaluation program on component parts and assemblies thereof;
2. A coordinated engineering liaison program between consumer and vendor of component parts;
3. A sound source qualification and parts improvement program;
4. Development of adequate specifications and application data covering extended life applications of component parts.

These elements are discussed as they apply to the current manufacture of electronic accounting and data processing machines by IBM.

Acceptance Sampling of Reliable TubesB. P. Goldsmith

The traditional method of checking acceptability of a lot of tubes for a particular elec-
trical characteristic has been to test a large sample and count the number of tubes beyond the minimum or maximum limits-inspection by attributes. This reduces the chances of accepting a lot with a high percentage of defectives, but gives no assurance that the lot is centered close enough to bogey or that the spread of the distribution is tight enough.

The increasing complexity of circuitry and the high standards of performance required in many types of equipment have increased the need for such assurance. It is gained most simply from inspection by variables.

Simply adding variables inspection criteria to the minimum and maxinum limits already on the TSS will not do an efficient job of separating good lots from bad. By proper coordination of the two types of inspection, attributes and variables, a high degree of discrimination can be achieved with a modest amount of testing.

Examples are based on recent TSS for type JAN 5744WA.

Cathode Interface Impedance Desimplified -H. B. Frost

Cathode interface impedance has usually been treated as if it could be represented at any given time by a parallel R-C combination. In actuality, however, the inpedance can be represented accurately only by an R-C network containing four elements. Moreover, the interface undergoes a reversible change of state with a relaxation time near one second as the cathode current is changed. These characteristics merit important consideration when specifications concerning cathode interface impedance are written.

Modern methods of preventative maintenance frequently allow replacement of those tubes with cathode interface impedance before they can cause the failure of large electronic systems such as digital computers. However, serious cases of cathode interface impedance may cause the tube population of such a system to have a short average life, perhaps less than 10,000 hours. Under such conditions, catastrophic failures-primarily an affliction of young
tubes-very likely will be more prevalent than would be the case if the average life were greater. Any increased level of catastrophic failures will cause a reduction system reliability which may be attributed indirectly to cathode interface impedance.

The Definition of Terms of Interest in the Study of Reliability-C. R. Knight, E. R. Jervis, and G. R Herd

The aim of this paper is to propose certain concents and definitions as aids in studies of the reliability of various products. "Reliability" and other terms commonly used in such studies are so defined that they can be measured and expressed quantitatively; and the theoretical relationship of components to the system is discussed. Reliability is studied in terms of discrete variables and continuous variables and their combined effects, with consideration of the interdependence of components. The concept of dependence is developed to facilitate measurement of the effectiveness with which components are incorporated into a system. The paper advances an alternate definition of "satisfactory performance" to the generally accepted one based on current specification practices. The new definition takes into account user acceptance or rejection of the product. Weighting functions are proposed to give mathematical expression to user opinion versus equipment performance characteristics.

## ULTRASONICS ENGINEERING

## PGUE-2, May, 1955

Composite Piezoelectric Resonator-W. G. Cady

Various types of composite resonators and their uses are summarized. The general equations are given for the transducer for plane ultrasonic waves, consisting of a crystal assembly with back and front plates. Applications are made to several simple cases, and expressions are tabulated for the amplitude of vibration for various combinations of half- and quarter-
wave components. Considerat on is given to the effect of the acoustic load on the frequency for maximal amplitude. Theoretical formulas are compared with experimental results for rods of aluminum and fused quartz excited in vibration by piezoelectric crystals.

Ultrasonic Cleaning of Miniature Devices -Q. C. KcKenna

Ultrasonic cleaning gives industry a new method of obtaining cleaning results previously unattained. By irradiating liquid cleaners with appropriately arranged transaucers, large volumes of intricate parts can be cleaned. Barium titanate ceramic transducers offer many advantages as sound generating elements. They can be operated at low voltages compared with quartz and can be cast in shapes which give high ultrasonic intensities. Through focusing, the ultrasonic cleaning process usually results in a more economical methed, saving time, labor, and space.

Power Measurements in UltrasonicsO. E. Mattiat

A Temperature Invariant Solid Ultrasonic Delay Line-Edwin Voznak and R. W. Mebs

A study was made of various metals and alloys in an effort to obtain a solid ultrasonic delay line that would be thermally stable with respect to time delay. Experimental data showing the effect of temperature or the propagation of ultrasonic waves are presented. An isoelastic alloy possessing a temperature coefficient of delay time of not more than 8 $\mathrm{ppm} /{ }^{\circ} \mathrm{C}$ over a temperature range of -50 to $200^{\circ} \mathrm{C}$ is described. This characteristic is superior to that of quartz or mercury by an order of magnitude.

Some Applications of the Linear Piezoelectric Equations of State-Rudolf Bechmann Notes on the Uses of Ultrasonics for the Finishing of Cathode Ray Tube Guns and Gun Components-W. F. Niklas

Biographical Notes on the Authors
Letter to the Editor
Recent Books on Ultrasonies
Cross Inder IRE-PGUE Transactions 1-3

## Books

## Transistor Audio Amplifiers by Richard F. Shea

Published (1955) by John Wiley and Sons. Inc.: 440 Fourth Ave., New York 16, N. Y. 207 pages $+x$ xiii pages +5 page index +5 page bibliography. Ilpages lutrat. $9 \frac{\text { page }}{\times 1} \times 6 . \$ 6.50$.

The stated object of this book is to provide the practical fundamentals of transistor applications and to show how these fundamentals may be used in the construction of audio amplifiers. In the reviewer's opinion the author has fulfilled his purpose competently. The book is a useful, lucid compilation of junction transistor circuit fundamentals, typical data on commercial units, design formulas, and practical audio circuitry.

Intended primarily for the experimenter and the designer of practical circuitry, the hook does not treat transistor physics or technology and uses a minimum of mathematics throughout. The emphasis is definitely on presentation of facts and figures rather than detailed whys and wherefores. One-third of the text is devoted to transistor characteristics and parameters, including the relationships among the various
equivalent parameters in current usage, and citing a good many commercial transistor specifications. One of the eight chapters gives a full treatment of the three basic amplifier configurations, complete with expressions for input and output resistances, and various gain figures under matched as well as unmatched conditions. Numerous graphs are included, which show at a glance how gains and input and output resistances vary with source and load resistances, supply voltages, currents, and temperature. The remaining chapters are devoted to coupled stages, with a detailed comparison of the various possible combinations of the three basic structures; preamplifiers, including, for example, consideration of transistor noise figure; Class A and B power amplifiers, with considerable number of transfer-characteristic graphs comparing departure from linearity under various conditions; and, finally, a few examples of hearing-aid and phonograph amplifier circuits, including the design steps leading to these circuits.

Transistor Audiap Amplifiers covers a fairly wide range of material in relatively few
and small-size pages; the treatment is therefore necessarily brief, and a number of points are brushed over rather lightly. Aside from this, and aside from a few inconsequential errors, Transistor Audio Amplifiers can well be recommended to those interested in becoming familiar with this increasingly important subject.

Ernest R. Kretzmer
Bell Telephone Labs., Inc.
Murray Hill, N. J.

## Sonics by T. F. Hueter and R. H. Bolt

Published (1955) by John Wiley and Sons, Inc,; 440 Fourth Ave., New York 16, N. Y. 440 pages $+x i$ pages +15 page index. Illustrated, $1 \times 6 . \$ 10.00$.

The announced purpose of this book is to provide a treatment of sonics-defined as the technology of sound as applied to problems of measurement, control, and processing-to serve the needs of the physicist as well as the practical design engineer. In this purpose the book has succeeded admirably. An adequate discussion is given of the fundamental acoustical principles, the properties of trans-ducers-principally piezoelectric and mag-
netrostrictive-and their applications in such practical applications as drilling, cleaning, sonic processing of metals, liquids and gases, and ultrasonic inspection of materials. The point is made that there is no frequency division into audible and inaudible sound for many of the processes. Thus, drilling of brittle material by an ultrasonic drill is similar in principle to the drilling of oil wells by much larger units working at subaudible frequencies.

Of interest to the physicist and chemist is a chapter on the principles of sonic testing and analysis. In this chapter there is a discussion of the various techniques for measuring the elastic properties of solids, i.e., both the dissipative and elastic moduli. A discussion of methods for measuring the viscoelastic properties of normal and polymer liquids is included. The appendix presents a short discussion of the significance of such measurements in the interpretation of the structure of liquids and solids

The coverage of techniques is complete and includes foreign techniques as well as domestic. For engineers and physicists desiring to acquaint themselves with the various techniques, methods for constructing apparatus and what can be done with them, Sonics is highly recommended. The book does not completely cover the communication field since no mention is made of the use of wave transmission in delay lines or mechanical wave filters and their applications in the communication systems. In the interpretive field, only basic principles are covered. In the opinion of the reviewer, this is the most complete book on techniques and applications of sonic processes that has yet appeared, and it should be in the library of all engineers and physicists dealing with these processes.
W. P. Mason

Bell Telephone Laboratories Murray Hill, N. J.

## Handbook of Microwave Measurements: Two Vols., Edited by Moe Wind and Harold Rapaport

Published (1954) by Polytechnic Inst. of Brooklyn, 55 Johnson St., Brooklyn 1, N. Y. Volume I, 20 Sec tions; Volume II, Illustrations; 4 Appendices. $8 \$ \times 11$ \$12.00.

The material in these two volumes has been assembled for the guidance of technical personnel in the field of microwave measurements. The editors, recognizing that much of the information in this field is widely dispersed, have endeavored "to present a unified collated handbook of microwave measurement methods" in order that many sound methods may not be overlooked and remain unutilized. The two volumes, comprising upward of one thousand pages, are the work of twenty-five contributing authors. The material has been divided into twenty sections, each devoted to a particular characteristic quantity such as power, attenuation, impedance, etc. A unique feature is that all of the text is contained in Volume One and all of the illustrations are in Volume Two.

The material was originally prepared for the Signal Corps Engineering Laboratories, evidently as an instruction manual for students without previous experience in the field. For this reason highly detailed step-by-step procedural instructions are given.

The amount of detail is perhaps greater than the average reader would wish, but it does serve to acquaint him with many aspects of the diverse measurement methods available. Although many of the methods are evaluated as regards accuracy, the distinction between fundamental methods and those relying on secondary standards is not emphasized. In particular, it was noted that the method of measuring conversion loss by directly measuring the input and output power is not given.

Each section begins with a theoretical analysis, quite detailed and complete, of the subject under discussion. Particularly valuable are the sections on propagation constant, impedance and dielectric constant. A great amount of detailed information is contained in these two volumes and those working in microwave measurements should find much which is useful to them.
C. F. Edwards

Bell Telephone Laboratories. Inc. Holmdel. New Jersey

## Television Interference, Third Edition,

 Edited by Philip S. RandPublished (1953) by Remington Rand. Inc., 315 Fourth Ave., New York, N. Y. 104 pages. $\$ 25$.

This book, like the two preceding editions, consists of reprints of technical articles on the subject of interference suffered by television receivers. The present volume contains 31 such articles and a list of recommended reading.

In view of the recent action of the FCC in proposing rules for the control of spurious radiations, this subject has become a matter of urgent interest to the designers of all electronic equipment. Most of this book is directed to the constructors and operators of amateur transmitters, but the subject of controlling spurious radiations in television receivers is also covered.

In a compilation of this type it cannot be expected that all of the material will be on a uniformly high technical level. However, taken with the preceding two editions, this book brings together a worth while collection of previously published material. Mr. Rand and the Remington Rand Corporation deserve the thanks of the industry for making this material available at a most modest price.

Donald G. Fink
Philco Corporation
Philadelphia, Pennsylvania

## The Oscilloscope at Work by A. Haas and R. W. Hallows

Published (1954) by Iliffe and Sons, Ltd., Stamford St., London S.E., England. 167 pages +4 page index. 319 figures. $8 \frac{1}{2} \times 5 \mathrm{~g}$. 15 s. 0 d .

The book deals primarily with measurements of electrical circuit characteristics by the use of the oscillograph as the indicating instrument. The electrical circuits for which the oscilloscope's use is described include basic electrical circuits, audio frequency amplifiers, rf amplifiers, oscillators, rectifiers, modulators, phase shifting and wave shaping circuits, and certain limited television receiver measurements. In describing the oscilloscope's use, many waveforms are employed and can be a valuable aid to the reader as a general guide in the types and limitations of measurements that can be
made in these fields. One complete chapter is devoted to oscilloscope operating troubles. Although somewhat limited, it is an excellent guide to the nore important defects an oscillograph might have and their effect on application of an oscillograph as an instrument tool.

It is the reviewer's opinion that the title of the book The Oscilloscope at Work is somewhat misleading. The actual fields in which the oscilloscope works are many, and yet this book covers only one small sector of them. The circuits in the first chapter are extremely simple compared to the modern cathode-ray oscillograph.

In summary, let me say that the electrical engineer who has not used an oscillograph may find this a valuable aid in learning some of the fundamentals of the oscillograph in his measurement work.
IV. G. Fockler

Allen B. DuMont Labs.

Operations Research for Management by Joseph F. McCloskey and Florence N. Trefethen

Published (1954) by The Johns Hopkins Press, Baltimore 18, Md. 350 pages +5 page index +xoxiv pages. $9 t \times 6 \%$. $\$ 7.50$.

This is a comprehensive collection of articles prepared by individuals who are experts in their respective areas to this new science. Careful reading should do much to explain Operations Research and how it may be used as a tool of management. The majority of the articles were presented at a seminar held by Johns Hopkins University in the spring of 1952. The volume is divided into three parts preceded by a well-written introduction by Dr. Ellis Johnson. This introduction paves the way for what is to follow, making it more understandable.

Part I covers the history of operations research and the concepts of it as a profession and a science. This section starts with World War II and follows through its evolution to the present time. It points out the similarity of systems evaluations, operations evaluation, operational analysis, and operations research.

The second part, dealing with methodology, describes some of the mathematical and statistical techniques employed as well as some of the basic philosophy underlying the use of these techniques. This part of the text is particularly difficult reading for the layman. Many of the words used are not normally encountered nor will they be understood by the average member of management. The authors, however, have done their utmost to define their unusual or complex vocabulary and symbolic logic in terms that are generally understood.

Part III contains a carefully selected number of case histories. The variety of cases is such that most members of management will find general ideas indicating how operations research might well be applied in their own organization. The article by Dr. Horace C. Levinson, "Experiences in Commercial Operations Research," is worth reading.

One should read the introduction and Part I first. Part III should be read next. Part II should be read last. While the great majority of managers may not thoroughly
understand the complexity of the methods described in Part II, it should give them a very definite idea of the value of these methods. It is unfortunate that a glossary of the technical terms used in this volume was not included. The bibliography is complete and well prepared.

Tom C. Rives General Electric Co.

Syracuse, N. Y

Advances in Electronics and Electron Physics: Volume Six Edited by L. Marton

Published (1955) by Academic Press Incorporated, 23 East 23 St., N. Y., N. Y. 518 pages + xi pages +19 page index. Illus, $9 \frac{1}{2} \times 6 . \$ 11.80$

This book is a collection of eight comprehensive reviews prepared by outstanding authorities from the United States, England and The Netherlands. Under the able editorship of L. Marton of the National Bureau of Standards, Volume Six carries on the fine tradition of this series. Reader interest will be particularly strong among physicists, chemical physicists, and radio engineers whose curiosity extends beyond the mundane problems of the radio-TV spectrum into areas where advance work is laying the foundation for knowledge from which many of the electronic inventions of tomorrow will spring.

The contributing authors are Elihu Abrahams, Rudolf G. E. Hutter, Henry F. Ivey, and W. M. Webster of the I'SA M. E. Haine and A. B. Pippard of England; and J. Smit, J. V'an Den Handel and H. P. J. Wijn of The Netherlands. The book is made up of chapters on: Metallic Conduction of IIigh Frequencies and Low Temperatures, Relaxation Processes in Ferromagnetism Physical Properties of Ferrites, Space Charge Limited Currents, A Comparison of Analogous Semiconductors and Gaseous Electronics Devices, The Electron Microscope, Traveling. Wave Tubes, and Paramagnetism.
lach chapter presents the reader with an erudite cross-section of contemporary research in the subject. A comprehensive list of references follows each writing. The authors offer more to interest the mature scientist engineer than the beginner. But whether the volume is used for instruction or for reference material for research underway, the reader will gain from those parts which pertain to his field of interest or specialization.

The Editorial Board assisting Dr. Marton in bringing this excellent compilation to press consists of Allibone, Casimir, DeVore, Dow, Nier, Nottingham, Piore, Ponte, Rose and Smith.

This is a fine book and would be an excellent addition to the library of anyone interested in advanced work in the areas reported upon.

Harold A. Zahl
Signal Corps Engineering Laboratories Fort Monmouth, New Jersey

Laplace Transforms for Electrical Engineers by B. J. Starkey

Published (1954) by Iliffe and Sons Ltd., Dorset House, Samford St., London, S.E. 1, England. 276 pages +3 page index. Illustrated. $8 \frac{1}{1} \times 5 \frac{3}{8} .308$

Perhaps the most efficient way of evaluating this book is to compare it to Transients in Linear Systems by Gardner and Barnes, since the latter is well-known to virtually
every electrical engineer. The word similar refers only to the expressed intent of the books, each being an introduction to the use of Laplace transform methods in solution of practical problens, Of the two, Starkey begins with more elementary considerations but quickly, in 81 pages, covers most of the material to be found in Gardner and Barnes work, although with fewer examples; Starkey's examples are all concerned with electric circuits while the other work also treats mechanical and acoustical problems.

The bulk of Starkey's book is concerned with topics which do not appear at all in Transients in Linear Systems: i.e., complex variable theory and the evaluation of inverse Laplace transforms by contour integration. These topics are developed from the start with sufficiently clear and detailed explanations so that a person, not previously familiar with Cauchy's theorem, will find everything he needs in order to understand mathematical methods commonly used to find inverse transforms.

This book can be recommended very highly to the serious student who wishes to obtain more than just a smattering of the Laplace transform method, who wants to understand it from a mathematical standpoint sufficiently well so that he is freed from dependence on tables of transforms, a desirable objective that becomes more and more necessary in advanced work.

On the negative side, however, one word of warning is necessary. Starkey depends strongly on intuition and his mathematics has more vigor than rigor. Occasional liberties are taken with regard to convergence of integrals, interchange of order of integration, etc., which a mathenatician would find quite hair-raising. For example, the integral in equation (9.6) does not exist unless one makes a qualification that does not appear in the text until three paragraphs later. Of course, to a practical man these mathematical questions will be regarded as mere nuisances; nevertheless they exist, and a prospective student should be cautioned that even in the most practical problems it will sometimes be necessary to use higher standards of mathematical rigor, not just for artistic reasons but in order to get the right answer.

This criticism is to be regarded as a very mild one; it is undoubtedly good pedagogy to defer considerations of rigor until after the student has a preliminary view of the field.
E. T. Jaynes

Stanford University
Stanford. Calif.

## Electromagnetics by John D. Kraus

Published (1953) by the McGraw-Hill Book Company, Inc., 330 West 42 nd St., New York $36, \mathrm{~N}$. Y 555 pages +10 page index +7 page bibliography +30 page appendix + xiii pages. 379 figures. $6 \times 9 \mathrm{a}$. $\$ 9.00$.

This excellent text on electromagnetic theory is distinguished by its clarity and logicality. Mathematical material is fully developed, few steps being omitted, or, where they are omitted their justification being clearly described in a concise manner. All notation is completely defined. The order of presentation of theoretical developments is generally that which is most satisfactory for purposes of learning; in most cases the
dependent variable is first formulated in an initial equation in terins of independent variables of broad significance, which, in turn are then determined in more specific terms. 'The book is copiously provided with illustrations and with worked ex.mmples, which, in themselves, form part of the text, and serve not only to apply the relationships obtained but also to develop them further. As the author says in the preface, simple special cases are usually considered first, and then with these as a background, the corresponding general cases are evoived. Vector concepts and operations are demonstrated ats a part of the text material, wherever they are needed.

These qualities of lucidity, though obviously indispensable in any textbook, are actually found in so few that Dr. Kraus's book may be fairly said to be one of the best available.

Emphasis is on electromagnetic field theory, but, though this point of view is stressed, circuit theory is not neglected, and its relation to field theory is pointed out.

The first seven chapters of the text are intended for use in an introductory onesemester field-theory course at about the third or fourth-year college level, while the last seven chapters are written for a somewhat more advanced course of the same length at the senior or first year graduate level.

Subjects treated include the static electric field, the steady electric current, the static magnetic field, charged particles in electric and magnetic fields, time-changing electric and magnetic fields, Maxwell's equations, plane waves in dielectric media, plane waves in conducting media, transmission lines, wave guides, antennas, and boundary value problems.

This book can be highly recommended, not only as a teaching text-book, but as a reference book for engineers and physicists.
I). B. Harris

Stanford University
Stanford, California

## Introductory Circuit Theory by Ernst A. Guillemin

Published (1953) by John Wiley \& Sons, Inc., 440 Fourth Ave., New York, N. Y. 545 pages +4 page index $+x x y$ pages. 199 figures. $6 \times 9 . \$ 8.50$.

Ernst A. Guillemin is Professor of Electrical Communication at the Massachusel.ts Institute of Technology, Cambridge, Mass.

We have here a text destined to have substantial influence on electrical engineering education and practice. Its importance is large, and it marks one of those milestones which give the steps-as contrasted with the inchlines-by which major progress is measured. We propose to discuss the work of Dr. Guillemin under four categories: its value in pointing the way toward a radical course (and ultimately curriculurn) change; emphasis on recently recognized techniques which contribute, at a cost, much toward simplification and condensation; new material for which the text would be valuable irrespective of its other contributions, and finally, some criticisms-not all favorable-of details of the book.

University curricula in engineering and a few other fields receive severe and continuing pressures to introduce new material, move "advanced" material to lower class levels, and otherwise to include in a
fixed span of time more and more. At the same time, raised living standards inculcate ideas of extensive leisure time, less stremuous work, and other concepts which permeate oncoming generations and lead to expectations of corresponding limitations on scholastic work weeks. By radical course and curriculum changes, electrical engineers who have kept in touch with engineering education have seen more and more intellectual material included in fewer and fewer class hours per week without great detriment to the engineering education of the students. Large parts of the "radical curriculum changes" have resulted from changes in emphasis, increases in efficiency, and condensation and elimination.

Dr. Guillemin's book qualifies under the first two categories, as a forerunner of another sharp change. Some material heretofore confined to advanced courses is presented for sophomore or junior use. Some eniphases, such as that on sine-wave Iriving forces, are reduced and postponed in order to introduce new and preceding emphases on transient, pulse, and impulse analysis. Some new material virtually untaught previously is introduced. The total is an increase in efficiency of major proportion and a change in emphasis long overdue to bring first circuit theory courses more in accord with today's electrical engineering.

We consider a few of the items which support this statement. Just as the elements of matrix theory can be taught withont difficulty to sophomores "whose mental attitude is not preconditioned" whereas seniors who have met matrices by backdoor methods and rumors have a psychological barrier which must be overcome, so, the teaching of the response of simple circuits to pulses and impulses can be used as a starting point of circuit analysis. And likewise the impedance concept, far more general than that of the impedance associated with the steady alternating state response of linear systems having impressed sine-wave driving forces, can be introduced early in technical life provided a firm base of general transient analysis precedes. Guillemin does these things, thus recognizing first that although the sine-wave driving force and response still remain paramount, relatively they are considerably less important than formerly and this should be reflected in beginning circuit theory courses; and second, that the impedance concept and the general pole-andzero approach have reached a maturity which calls for introduction in the same beginning circuit theory courses.

The book under review emphasizes certain techniques-not original with the an-thor-which in some respects have had far too little attention. 'To cite but one relatively
simple example, a considerable emphasis is placed on 'Thevenin's theorem so that the student will learn how it joins with duality in enabling many circuit problems to be solved by way of one solution. "The insistence made that circuit theory is an abstraction which may or may not have a one-to-one correspondence with a specilic physical circuit is highly desirable in a beginning circuit theory course. And the insistence on the use of simple numbers and problems and examples (e.g., $K=1, L=2, C=3$, although somewhat of a far cry from some of Dr. Guillemin's previous works, recognizes the current trend to minimize distractions from basic ideas and mechanical effort required of the student to demonstrate his prowess in a highly theoretical field.

New material in the text stands out from the first chapter. Network topology for sophomores or juniors may seem somewhat startling at first glance, but it is actually a topic which should long since have been introduced in beginning circuit courses. Dr. Guillemin's book is probably the first of its kind in this respect, and the job is well done for the level at which the text aims, which is not to say that there are not weaknesses in it. But the general idea that, to take one example, the usual fuzzy introduction of loop currents in a circuit theory course should be superseded by going back an order of magnitude in theory to enable the student to have an intelligent grasp of the problem behind the introduction of loop currents is simply another of those long overclue reforms needed in elementary circuit theory. 'There is much incidental other new material in the book, on which we will not touch.

In the introductory paragraph of this review we reserved our final comments for "some criticisms-not all favorable-of details of the book." For a text of so much potential value, the criticisms are of an order of magnitude less in importance than the favorable comments which have preceded. With this understanding, we mention our first and outstanding complaint-verbosity. The book could have been written in about three-quarters of the space it takes, without loss to the reader. We will not illustrate this, and it can be taken as simply the opinion of the reviewer, but practice condensations have been supplemented by informal comments of others, so that the reviewer is inclined to feel that the defect stands out. Furthermore, occasional poor English and misuses of technical words are jarring. Guillemin's use of "potential" for example is of ten dubious to say the least, and the definition of "passive" has to be caught on the run. Could it be more than coincidence that neither "potential" nor "passive" appear in the index of the book?

Words like "resistor" and "resistance" both appear, but there is no consistency in their use. And words like "clearly" and "surely," those common introductions to a poor argument, appear in good-sized groups (five in one paragraph on page nine). But more important, the author uses " s " instead of " $p$ " for complex frequency, and the reviewer has reason to believe this has been done simply because of ignorance of IRE standards. "Yo use a well-worn cliché, it is regrettable that a text so modern in essence should perpethate outmoded nomenclature rather than join in the attempt to clear con fusion in the field.

There are a few statements-the definition of the "value' of an element" on page three-which seem peculiar; and even the last paragraph of the Introduction, telling of the simplest broad class of networks (to which the book is devoted) omits an essential qualification.

The book is so arranged that chapters four through eight constitute a unit which the author points out can be used for a reduced course. In using this text, it is desirable that at least some instructors experiment, and one suggestion is to have chapters one through three follow the four through eight sequence.
J. (i. J3RAINERD

Moore School of Electrical Engineering
University of I'enssylvania

## Recent Books

Abstracts of the Literature on Semiconducting and Luminescent Materials and Their Applications. Compiled by Battelle Memorial Institute. John Wiley and Sons, Inc., 440 Fifth Ave., New York 15, N. Y. \$5.00.

Booth, Andrew D., Numerical Methods. Academic Press, Inc., 125 East 23 St., Cew York $10, \mathrm{~N} . \mathrm{Y} . \$ 6.00$.
Jacolson, Arvid II., ed., Proceedings of the First Conference on Training l'ersonnel for the Computing Machine Field. Wayne Iniversity Iress, Detroit, Michigan. $\$ 5.00$.
Marcus, William, and Levy, Alex, Elements of Radio Servicing. Mcriraw-Hill Book Company, Inc., 330 West 42 St., New York 36, N. Y. \$6.00.
Petrovsky, I. G., Lectures on P'artial Differential Equations, trans. by A. Shenitzer. Interscience Publishers, Inc., 250 Fifth Ave., New York 1, N. Y. \$5.75.
Rider's Specialized Tape Recorder Manual, Volume One. John F. Rider D'ublisher, Inc., 480 Canal St., New York 13, N. Y. \$4.50.
Shedd, Paul C., Fundamentals of Electromagnetic Waves. Prentice-Hall, Inc., 15 East 26 St., New York 10, N. Y. \$6.00.

# Abstracts and References 

Compiled by the Radio Research Organization of the Department of Scientific and Industrial Research, London, England, and Published by Arrangement with that Department and the Wireless Engineer, London, England

NOTE: The Institute of Radio Engineers does not have available copies of the publications mentioned in these pages, nor does it have reprints of the articles abstracted. Correspondence regarding these articles and requests for their
procurement should be addressed to the individual publications, not to the IRE.
pres abstracted. Correspondence regarding these articles

Acoustics and Audio Frequencies. ..... . 900
Antennas and Transmission Lines...... 900
Automatic Computers.................. . . . . 901
Circuits and Circuit Elements......... . . 901
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Geophysical and Extraterrestrial Phe-
nomena.
904
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Materials and Subsidiary Techniques... 906
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908
Measurements and Test Gear
Other Applications of Radio and Electronics.
Propagation of Waves.
Reception.
Stations and Communication Systems.
Television and Phototelegraphy.
Transmission
Tubes and Thermionics.
Miscellaneous. left of each Abstract is its Universal Decimal Classification number and is not to be confused with the Decimal Classification used by the United States National Bureau of Standards. The number in heavy type at the top right is the serial number of the Abstract. DC numbers marked with a dagger ( $\dagger$ ) must be regarded as provisional.

## C.D.C. CHANGES

In anticipation of a new edition of the Universal Decimal Classification Abridged English Edition (BS 1000 A ), certain changes in U.D.C. numbers will be made in this and subsequent issues. The new numbers used will be:

Radio astronomy: 523.16
Ultrasonics: 534 subdivisions with the special analytical subdivision -8 attached.
Sound recording and reproducing: 534.85
Electroacoustic problems, transduction, etc.: 534.86.

## ACOUSTICS AND AUDIO FREQUENCIES

534.121 .2

1535
Symmetry of Vibrating Square Membrane —M. D. Waller (Proc. Phys. Soc., vol. 67, pp. 895-898; December 1, 1954.) Vibrations of free square plates always conform with the symmetry of the surface; this experimental conclusion conflicts with the accepted theoretical conclusion that two normal modes of vibration of the same frequency can combine regardless of symmetry considerations. This conflict is discussed. See also Acustica, vol. 4, no. 6, pp. 677$680 ; 1954$.

## 534.2:534.833

1536
The Propagation of Sound in Granular Sub-stances-H. Schmidt. (Acustica, vol. 4, no. 6, pp. 6.39-652; 1954. In German.) Theory is developed and neasurements are reported of the loss factor and propagation velocity of granular substances (a) for substances packed in a barshaped container, (b) for a layer of the sub)stance on top of an Al bar. In case (a) sand, witl varying water content. glass spherules 1 mm in diameter, finely broken brick, cinders and coarse iron filings were used, in case (b) only sand and brick. For application as a filling e.g. for sound insulation in buildings, cinders or broken brick are recommended, with wood

The Index to the Abstracts and References published in the PROC. IRE from February, 1954 through January, 1955 is published by the PROC. IRE, April, 1955, Part II. It is also published by Wireless Engineer and included in the March, 1955 issue of that journal. Included with the Index is a selected list of journals scanned for abstracting with publishers' addresses.
shavings or rock wool added where necessary to reduce the resonance frequencies to very low values.
534.21-8

1537
The Propagation of Sound in Solutions of Rubber and Perspex-A. W. Pryor. (Acuslica, vol. 4, no. 6, pp. 658-661; 1954.) Absorption measurements at four frequencies in the range 4-17.3 inc were made on three rubber-inbenzene solutions and on one perspex-inpyridine solution. "The absorption was less than in the pure solvents and there is evidence of relaxation. Measurements of the shear viscosity of a 10 per cent rubber solution showed that the relaxation of the high flow viscosity was complete even at $50 \mathrm{kc} / \mathrm{s}$. The absorption in the solutions must therefore be ascribed to the 'bulk viscosity" of the solvents."
534.22-14

1538
Measurements by Optical Methods of the Sound Velocity in Aqueous Solutions of Electrolytes in dependence on Concentration and Temperature-K. Tamm and M. (.). Haddenhorst. (Acuslica, vol. 4, no. 6, pp. 653-6.57; 1954.)
534.6:621.373.4:621.374

1539
Generation and Use of Single-Frequency Pulses in Electroacoustics and Musical Acous-tics-H. Lackner. (Ost. Z. Telegr. Teleph. Funk Fernsehtech., vol. 8, pp. 141-152; November /December, 1954.) Analysis is presented showing low the sjectrum of a short train of sinusoidal oscillations depends on its duration and on the initial and final phase. A circuit for a generator providing pulses of variable duration and phase is described. Advantages of using such signals for testing loudspeakers and musical instruments are indicated.
534.8461540

Improvements in the Acoustics of the Budapest Civic Theatre-T. Tarnoczy. (Acustica, vol. 4, no. 6, pp. 665-671; 1954. In German.) The upper walls and part of the ceiling were provided with a diffusive surface, made of reinforced plaster of Paris. The rest of the ceiling was plastered over. The lower walls and the new proscenium were lined with wood paneling, which acted as an acoustic resonator. A sound-amplification system is available if required. These measures sufficed to remove flutter echoes and almost all other audible echoes, to increase the reverberation time to 1.6 seconds, and to make the distribution of sound energy completely satisfactory in almost every part of the theater. High frequencies have been enhanced and tonal quality is good.

### 534.86

Advances in the Reproduction of MusicF. Winckel. (Funk u. Ton, vol. 8, pp. 604-607 and 649-652; November and December, 1954.) A series of brief reports on lectures delivered
at three conferences hela in 1954. The subjects covered include a combination loudspeaker init, automatic volume control, a new com-pressor-expander system with variable time delay, and the quality of microphones and loudspeakers.
534.861:534.76

1542
Experiences in Stereophonic Broadcast Transmissions--J. J. Geluk. (Funk u. Ton, vol. 8, pp. 631-634; December, 1954.) Report dealing mainly with Dutch experimental transmissions.
534.862:534.76
1543

Stereophonic Sound-Film Recording and Reproduction-H. Friess. (Funk u. Ton, vol. 8, pp. 622-630; December, 1954.) A review of current problems and teclinique.
621.395 .616

1544
Note on the Stabilization of the Response of a Capacitor Microphone-C. Colin. (Jour. Phys. Radium, vol. 15, pp. 820-822: December, 1954.) It is possible to improve the response of a capacitor microphone by negative feedback if (a) an auxiliary electrode is included, and (b) the voltage output is in phase with diaphragm displaceinent. An analysis is made of the necessary coupling conditions, both mechanical and electrical.

### 621.395.625.3

1545
A Survey of Magnetic Recording-S. J. Begun. [Elec. Engng. (New York), vol. 73, pp. 1115-1118: December, 1954.] Includes a description of operation and an indication of the development and applications.

### 621.395.625.3

1546
Making Magnetic Sound-Records VisibleW. Guckenburg. (Funk u. Ton, vol. 8, pp. 600604: November, 1954.) A review of various methods using suspensions of ferromagnetic dust. Photographs illustrating applications, such as detecting faulty magnetic heads, are briefly discussed.

## ANTENNAS AND TRANSMISSION LINES 621.315.28:621.395.44 <br> 1547

A Transatlantic Telephone Cable-M. J. Kelly, G. Kadley, (r. W. Gilman and R. J. Halsey. (Proc. 1EE, Part B, vol. 102, pp. 117130; March, 1955. Discussion, pp. 130-138.) The inadequacy of radio circuits to cope with transatlantic comnunications is discussed and a general description is given of the cable system for linking the United Kingdom, Canada and the United States; th:s is planned for completion in 1956 and will provide 36 telephone circuits across the Atlar.tic and 60 between Newfoundland and Nova Scotia. The project has been made possible by the development of submerged repeaters containing long-life tubes and other components.

The Launching of a Plane Surface WaveG. J. Rich. (Proc. IEE, Fart B, vol. 102, pp. 237-246; March, 1955.) Fixperiments are described on the propagation of surface waves at about $3 \mathrm{~cm} \lambda$ over plane: brass plates with polystyrene coatings of thickness such that only the TMo mode is rropagated. Various launching devices were tred; the best found was a "double-cheese" device with a launching efficiency of about 50 per cent. Experimentally found values for the efficiency of this device are in agrecment with theory for apertures of dimensions smaller than $\lambda$. Deviations from theoretical values observed for larger apertures indicate that to obtain good launching efficiency the field distribution within the aperture must be to a close approximation exponential.


#### Abstract

621.372.22

1549 The Matrix Equation of Loss-Free Exponential Lines-B. Beghin. [Compl. Rend. Acad. Sci. (Paris), vol. 240, pp. 168-170; January 10, 1955.] A solution of the differential equation for the exponential line is presented in the form of the product of two affine exponential expressions representing a ladder arrangement of two quadripoles, the one a perfect transformer and the other analogous to a uniform line. Reffection can be eliminated by


 terminating the line with a complex impedance.
### 621.372 .8

1550
The Propagation of an Electromagnetic Wave along an Infinite Corrugated SurfaceR. A. Hurd. (Canad. Jour. Phys., vol. 32, no. 12, pp. 727-734; December, 1954.) The problem is analyzed by a method based on the calculus of residues; the slot walls constituting the corrugations are assumed to be vanishingly thin. Exact expressions are obtained for the mode amplitudes and phase velocities; these are valid for frequencies near cut-off, when the number of slots per wavelength is about five or more.
621.372 .8

1551
A Simple Waveguide Directional CouplerP. Andrews. (Jour. Brit. IRE, vol. 15, pp. 112116; February, 1955.) A simply constructed, orthogonally crossed unit is described giving constant coupling and directivity over a wide frequency band. Coupling factor and directivity are calculated for two typical cases for 3 inch $\times 1 \frac{1}{2}$ inch waveguide. An expression is derived for the frequency sensitivity of the device.

### 621.372.8.002.2

1552
Surface Roughness and Attenuation of Pre-cision-Drawn, Chemically Polished, Electropolished, Electroplated and Electroformed Waveguides-J. Allison and F. A. Benson. (Proc. IEE, Part B, vol. 102 pp. 251-259; March, 1955.) "Detailed examinations of certain 3 cm waveguides hare shown that the surface finish of precision-drawn tubing as manufactured at present is quite adequate for most applications. Such surfaces, however, may be improved, if desired, by careful chemical or electrolytic polishing or electroplating in bright baths under closely controlled conditions. Some information on the surface finish of copper guides electroformed on various types of mandrel is also presented. Formulae for calculating the attenuation of any H or E mode in a rectangular waveguide, so as to take account of surface roughness, have been developed from the original expressions derived by Kuhn. A method is given for determining the actual value of attenuation in a waveguicle sample without having to make careful measurements with elaborate equipment on long specimens. A description is included of a new and simple technique, involving electropolishing, for examining the internal surface finish of waveguides; the method cannot, however, be used successfully on silver-plated sections."
621.372.8.029.65.002.2

1553
The Electroforming of Components and Instruments for Millimetre WavelengthsA. F. Harvey. (Proc. IEE, Part B, vol. 102, pp. 223-230; March, 1955.) Electroforming processes considered for producing waveguide components include periodic-reverse-current plating. Either permanent or disposable formers may be used. The electroformed parts are designed to facilitate subsequent machining. Data are tabulated on a range of waveguides with internal cross sections from $0.28 \times 0.14$ inch to $0.034 \times 0.017$ inch which lave been standardized for the Joint Services, and various units incorporating these elements are illustrated.
621.396.67:539.23 1554
Possibilities of Radiation from Thin Metal Films-M. Gourceaux. [Compt. Rend. Acad. Sci. (Paris), vol. 240, pp. 170-172; January 10, 1955.] Systems with rotational symmetry and circular current paths are considered. Expressions are derived for the values of film thickness and current frequency for which radiation can occur with cylindrical and spherical arrangements.
621.396.676.2

1555
The Notch Aerial and some Applications to Aircraft Radio Installations-W. A. Johnson. (Proc. IEE, Part B, vol. 102, pp. 211-218; March, 1955.) Analysis and experiments show that metallic sheets with notehes cut perpendicular to an edge can make very efficient antennas. Wide-band characteristics can be obtained with a notch about $\lambda / 4$ deep, but the detail of the polar diagram may vary with frequency, depending on the shape and size of the sheet. Notches short compared with $\lambda$ termed "nitches"-can be operated either as fixed-tuned narrow-band antennas or, in association with a tuning unit, over a wider band. For reception at frequencies around 100 mc a notch of length about 3 inches is adequate. Tests made with aircraft tail fins and wings are described; hf., vhf and uhf types are practical.

### 621.396.677:523.16 <br> 1556 <br> Aerial Smoothing in Radio Astronomy-

 Bracewell and Roberts. (See 1638.)621.396.677.029.62/.63:621.397.6

1557
U.H.F. and V.H.F. Antenna-R. F. Kolar. (Radio-Electronics, vol. 25, pp. 60-62; December, 1954.) A low-cost television receiver antenna unit is described; the design is based on stacked uhf V's and an uhf dipole and reflector. The power gain relative to a $\lambda / 2$ dipole is a few decibels.
621.396.677.31

1558
Optimum Element Spacing of Uniform Broadside Arrays-R. F. H. Yang. (Wireless Eng., vol. 32, pp. 115-116; April, 1955.) Collinear and curtain arrays are considered. A brief analysis indicates that the optimum element spacing for a collinear array is $\lambda$, irrespective of $n$, the number of elements, while the optimum spacing for a curtain array is given approximately by the formula $d / \lambda=0.6+\frac{1}{3}$ $\log _{10} n$. This result is different from that obtained by Hammond (1909 of 1953).

### 621.396.677.8.029.65 1559 <br> Some Experiments on the Reflecting Prop-

 erties of Metal-Tube Lens Medium-E. M. Wells. (Marconi Rev., vol. 17, no. 114, pp. 7485; 3rd Quarter, 1954.) Experiments described show qualitatively the variation of the reflection coefficient with angle of incidence, plane of incidence, and plane of polarization. With the $E$-vector in the plane of incidence a very pronounced anomalous reflection was observed. This was further investigated. The wavelengths used were in the band $8-10 \mathrm{~mm}$.
## AUTOMATIC COMPUTERS

Wide-Band Analog Function Multiplier-
J. A. Miller, A. S. Soltes and R. E. Scott. (Electronics, vol. 28, pp. 160-163; February, 1955.) Operation of the multiplier is based on the formula $x y=\frac{1}{4}\left[(x+y)^{2}-(x-y)^{2}\right]$. The squaring is effected by means of beam-deflection tubes with parabolic characteristics [2940 of 1950 (Soltes)].
681.142

1561
Analogue Computer with Stepping-Switch Drive for the Solution of Algebraic Equations of up to Sixth Degree-II. IIörner and H. Zemanck. (Ost. Z. Telegr. Teleph. Funk Fernsehtech., vol. 8, pp. 153-158; November/December, 1954.) The analogy used is that between a sinusoidal voltage and a complex quantity. The representation of the polynomials is expected using feedback amplifiers as computing elements. The zero wints corresponding to the solutions are found automatically.
681.142

1562
Some Devices permitting Study of the Variable Regime in the Transpo:tation of Gas in Pipes-A. Blanc. [Compt. Rend. Acad. Sci. (Paris), vol. 2 10, pp. 45-46; January 3, 1955.1 A particular apolication is discussed of the analog-computer arrangements described previously (2317 of 1954).
681.142

1563
Programming a Digital Computer for Cell Counting and Sizing-W. Welkowitz. (Rev. Sci. Instr., vol. 25, pi). 1202-1204; December, 1954.)

### 681.142:538.221

1564
Magnetic Materials for Digital-Computer Components: Part 1-A Theory of Flux Reversal in Polycrystalline FerromagneticsMenyuk and Goodenough. (See 1713.)
681.142:621.3.042

Logical and Control Functions performed with Magnetic Cores-S. Guterman, R. D. Kodis and S. Ruhman. (Proc. IRE, vol. 43, pp. 291-298; March, 1955.) Use of magnetic cores with square hysteresis loops in digitalcomputer systems based on the "single-line" shift register is described.

### 681.142:621.375.23

1566
A Technique for Nonlinear-Function Gener-ation-P. N. Nikiforuk. (Electronic Eng., vol. 27, pp. 118-119; March, 1955.) A method is described for converting a circuit whose output is proportional to the $n$th root of the input to give the $n$th power, or vice versa. A practical circuit, based on a variable-gain feedback amplifier, for dealing with squares and square roots is illustrated.
681.142:621.376

1567
Two New Electronic Analog MultipliersM. A. Meyer and H. W. Fuller. (Rev. Sci. Insir., vol. 25, pp. 1166-1172; December, 1954.) Improved types of four-quadrant multipliers for deriving the product of two time functions are described. One uses a double-amplitude modulation scheme applying the work of Sternberg and Kaufman on the two-frequency-modulation-product problem (2212 of 1954). The other uses successive phase modulation and amplitude modulation of a carrier. The accuracy obtainable is discussed.

### 681.142

1568
Automatic Digital Computation [Book Re-view]-Publishers: H. M. Stationery Office, London, Eng. 296 pp., 21s. (Instr. Practice, vol. 8, p. 1094 ; December, 1954.) Contains discussions and original papers presented at the symposium held at the National Physical Laboratory in March, 1953.

## CIRCUITS AND CIRCOIT ELEMENTS

### 621.314.22.042

1569
Calculation of Losses in Transformer Cores -A. L. Morris.[Engineer (London), vol. 198, pp. 837-839 and 875-877; December 17 and 24,
1954.] Expressions for core losses are develoned which take account of the nonuniformity in the flux distribution across the face of the core plates. The nonuniformity across the stack of plates is neglected.
621.318.4:538.221:621.396.822

1570
The Noise generated in a Coil with a Ferromagnetic Core-G. Builder and D. Haneman. (Aust. Jour. Phys., vol. 7, pp. 654-658; December, 1954.) The discussion presented indicates that when there is no varying or alternating magnctization of the core, the noise can be calculated from the Nyquist formula, provided that the resistance of the coil is measured using an alternating voltage of such magnitude that the magnetization process remains linear.

### 621.318.435.3

1571
Composite Cores for Instrument Transduc-tors-E. H. Frost-Smith and A. E. DeBarr (Proc. IEF, Part II, vol. 101, pp. 663-667; December, 1954. Discussion, pp. 667-671.) The characteristics of self-excited transducers are markedly dependent on the properties of the core material. Cores with satisfactory properties can be produced at reasonable cost by combining different materials. Experimental results are reported for a core comprising a mumetal bridge and a U-shaped grain-oriented $\mathrm{Si}-\mathrm{Fe}$ yoke.
621.318.435.3:621.375.327

1572
Auto-Self-Excited Transductors and PushPull Circuit Theory-A. G. Milnes and T. S. Law. (Proc. IEE, Part II, vol. 101, pp. 643 662; December, 1954. Discussion, pp. 667671.) The modes of operation of transductors with parallel, bridge, and center-tap comections are explained. Expressions are derived for the sensitivity and time-constant under idealized conditions. An examination is made of the use of transductors in push-pull pairs to obtain magnetic amplifiers with outputs whose polarity depends on the signal sense; the transductors are assumed to be of the anto-self-excited type with resistance loads.
621.318 .5

1573
Component Design Trends-Special-Pur pose Relays gain New Uses-F. Rockett. (Electronics, vol. 28, pp. 150-156; February 1955.)
621.318.5:621.318.134 1574

Time Delay in High-Speed Ferrite Microwave Switches-R. C. LeCraw and H. B Bruns. (Jour. Appl. Phys., vol. 26, p. 124; January, 1955.) Measurements indicate that for an $X$-band switch using a rod of a particular $\mathrm{Mg}-\mathrm{Mn}$ ferrite and actuated by a current pulse of given magnitude and rise time $17 \mathrm{~m} \mu \mathrm{~s}$, the time delay of the switch is about $3.2 \mathrm{~m} \mu \mathrm{~s}$.
621.319.4:621.315.614.6

1575
Study of the Dielectric in Paper Capacitors -E. Briganti. (Alla Frequenza, vol. 23, pp 139-156; June/August, 1954.) Measurements were made of the loss angle of wet and dry paper at -5 degrees $C$., and of the permittivity of the impregnant and of the impregnated paper. From these an equivalent circuit was derived for the capacitor, which can be used in estimating the quality of the paper and desirable conditions of manufacture. Breakdownstrength measurements were also made to determine the optimum applied voltage.
621.372 .5

1576
Arc Inductance and Dynatron Capacitance -J. Groszkowski. (Bull. Acad. Polon. Sci., Classe 4, vol. 2, no. 1, pp. 41-45; 1954. In English.) The stability of systems with non linear negative resistances of arc and dynatron types is discussed. Analysis indicates that though the arc can have some inductive properties and the dynatron capacitive ones, these are not essential to an explanation of the behavior of the systems. Transient conditions and steady oscillation states are considered
621.372 .5

1577
Further Bounds existing on the Transien Responses of Various Types of NetworksA. H. Zemanian. (Proc. IRE, vol. 43, pp. 322326; March, 1955.) Five theorems further to those presented previously (2331 of 1954) are proved and illustrated.
621.372 .5

1578
An Approximate Treatment of Cascaded Four-Terminal Networks-FI. L. Arinstrong (Electronic Eng., vol. 27, pp. 130-131; March, 1955.) "An approximate expression is derived for the $n$th power of a $2 \times 2$ matrix. The result is used in an approximate treatment of a ladder network used as a filter."
621.372 .5

1579
The Approximation Problem of Network Synthesis-S. Winkler. (Trans. IRE, vol. C'T-1, no. 3, pp. 5-20; September, 1954.) A review with 240 references.
621.372 .5

1580
Series Resonant Circuit Theory-A. J. Lyon. (W'ireless Eng., vol. 32, pp. 107-108; April, 1955.) Expressions are derived for the fractional errors in tuning capacitance, resonance frequency, maximum current, and selectivity at current resonance due (a) in the case of frequency tuning, to the frequency dependence of the circuit resistance, and (b) in the case of capacitance tuning, to capacitor losses and coil self-capacitance.

### 621.372.5:512.83

1581
The Mesh Counterpart of Shekel's Theo-rem-S. Seshu. (Proc. IRE, vol. 43, p. 342; March, 1955.) It is shown that whereas the determinant of the adinittance matrix of a network is independent of the choice of reference mode [ 2878 of 1954 (Shekel)], the corresponding statement for the mesh determinant is not true.

### 621.372.5:621.3.018.7

1582
Distortion of Arbitrary Waveforms by Resonance Sections-E. Williain. (Funk u. Ton, vol. 8, pp. 592-509; November, 1954.) The method developed earlier (2032 of 19.54) for $R C$ sections is extended to sections including an incluctance.
621.372.5: 621.318.134

1583
Extension of Nonreciprocal Ferrite Devices to the $500-3000 \mathrm{Mc} / \mathrm{s}$ Frequency RangeR. H. Fox. (Jour. Appl. Phys., vol. 26, p. 128; January, 1955.) Calculations indicate that devices, such as circulators, of reasonable dimensions can be designed for operation at frequencies below 3 kmc by using static magnetic fields of intensity greater than required for ferromagnetic resonance.
621.372.512.2.029.65:621.372.8

1584
A Short-Slot Hybrid for $9 \mathrm{~mm}-E . \mathrm{M}$. Wells. (Marconi Rev., vol. 17, pp. 86-87; 3rd Quarter, 1954.) The $X$-band junction described by Riblet ( 1833 of 1952) was redesigned for a wavelength of 9 min . A bricf illustrated note on this junction is given.

### 621.372.512.24

1585
Resonance Conditions in a System of Two Circuits with Inductive Coupling-U. Ruelle. (Alla Frequenza, vol. 23, pp. 157-177; June /August, 1954.) Assuming $M, R, L$ and $C$ constant, conditions are derived for the existence of one minimum and two maximum values of current in the secondary, in the general and in two particular cases. The results are applied to a band-pass filter, and presented graphically for two different parameter values. Low $Q$ values are assumed.

### 621.372 .54

1586
Microwave Filters-E. Willwacher. (Fernmeldetech. Z., vol. 7, pp. 694-704; December, 1954.) Theory of filters comprising coaxial lines and/or waveguides is cleveloped by reference to equivalent lumped-constant circuits. Band-
pass and band-stop filters of ladder, amplifier and bridge types are discussed.

### 621.372.54:621.396.67 1587 <br> Tunable Microwave Aerial Diplexer-

 O. Laaff. (Fernmeldetech. Z., vol. 7, pp, 688-693; December, 1954.) A diplexer for the band 2.1 2.3 kmc is based on the use of continuously tunable resonant circuits comprising coaxialline sections short-circlited at the one end and having a third coaxial conductor insulated from and sliding within the inner conductor. In a particular design illustrated, for $100-\mathrm{mc}$ separation between transmission and reception frequencies, the transmitted energy penetrating into the receiver is attenuated by 65 db .
### 621.372.543.2

1588
Theoretical Investigation of Three-Stage Tchebycheff-Type Band-Pass Filters-B. Betzenhammer and E. Henze. (Arch. elekl. Übertragung, vol. 8, pp. 545-552; December, 1954.)
621.373 .4

1589
Judging the Qualite of Oscillator CircuitsH. Haller. (Funk u. Ton, vol. 8, pp. 565-575; November, 1954.) Discussion of the frequency and amplitude stability of oscillators considered as composed of a slightly nonlinear fre-quency-independent ampliser and a linear frequency-dependent feedback quadripole. Stability is found to depend on the rate of change of the imaginary part of the voltage transfer characteristic of the feedback quadripole. The use of this criterion in the design of bridgestabilized $R C$ oscillators is illustrated. Measurements on and various faults of experimental oscillators are briefly discussed.
621.373.4:621.374:534.6

1590
Generation and Use of Single-Frequency Pulses in Electroacoustics and Musical Acous-tics-Lackner. (See 1539.)
621.373.431.1 +621.375.2.018.756 1591 Experimental Investigations on Multivibrators and Amplifier Circuits with Secondary Electron Emission Valves as described by Kroebel-K. E. Rumswinkel. (Z. angew. Phys., vol. 6, pp. 551-556; December, 1954.) Multivibrators of the type described by Kroebel (383 of February) were investigated. Pulse flank slopes of the order of $10^{10} \mathrm{v} /$ second, up to pulse amplitudes of 67 v , were obtained. The multivibrator can be modified to act as a pulse amplifier without feeciback.
621.373.44:535.33

1592
Equipment for Excitation of Spectra by High Frequency Pulses-L. Minnhagen and L. Stigmark. (Ark. Fys., vol. 8, pp. 471-479; December 14, 1954. In English.) The equipment comprises a pulse generator which controls a Clapp oscillator (2193 of 1948) followed by a frequency multiplier and amplifiers. Average hif power transferred to the discharge tube is about 500 w , with peak power of 3 kw ; the frequency is about 9 mc . Ar spectra obtained are shown.

### 621.373 .52

1593
Transistor Frequency Standard-J. H. Smith, Jr, and M. Camplbell. (Tele-Tech, vol. 13, pp. $90-91,135$; December, 1954.) This unit, designed primarily for geophysical prospecting, has a printed-circuit base and uses transistors in all stages. Accuracy is within 1 part in $10^{4}$ over the range - 40 degrees $F$ to +140 degrees $F$. An $\%-k c$ oscillator is followed by a pulse-forming anplifier and first divider stage, then by two dividers in parallel, whose outputs are mixed to obtain an output frequency of 100 cps .
621.375.1.024

1594
The Transient Response of Direct-Current Amplifier Systems-J. H. Sanders. (Jour. Sci. Instr., vol. 31, pp. 453-455; December, 1954.) Dc amplifiers of the dc-ac conversion type
have an upper frequency response limited by the detector circuit, and vhen negative feedback is used transients are amplified considerably more than the steady signal. The form of the transient response and methorls of reducing its magnitude are discussed.
621.375.2.049 1595

A Long-Lived Packaged Amplifier for Air-craft-J. G. Matthews. (Bell. Lab. Rec., vol. 32, pp. 462-466; December, 1954.) The development and construct.on of units with at probable life of 2,000 hours are described. Selected tubes are pressed into cast Al wells lined with a silicone rubber paste which is hardened by a short curing process; this is more effective for heat tıansfer than Al dust or foil. Deposited carbon resistors and capacitors serviceable at 125 degrees C . are mounted on a phenolic board supported by the Al base and wired to a recessed plug. The assembled unit, after testing, is embedded in liquid plastic which is then solidified. Test figures for heat dissipation under different conditions are given.

### 621.375.221:621.372.512 1596

Amplifier Stages with Transitionally Coupled Two-Stage Band-Pass Filters, particularly for Large Bandwidths - W. Mansfeld. (Funk u. Ton, vol. 8, pp. 576-591; November, 1954.) The amplitude and group-delay characteristics of an amplifier stage consisting of two coupled circuits are analyzed for the case when the amplification is constant over a wide frequency band ("transiticinal" coupling). The case when the damping factors $d_{1}$ and $d_{2}$ of the two circuits are equal is considered first, and formulas are also given tor the cases of either $d_{1}$ or $d_{2}$ tending to zero. Formulas are also given for transforming a filter with indirect inductive coupling into one with direct inductive coupling. Design curves are shown.
621.375.23:621.3.016.35 for Muitiple 1597

Nyquist's Criterion for Feedback Circuits-I, Tasny-Tschiassny. (Wireless Eng., vol. 32. pp. 114-115; April, 1955.) A method of deriving the stability criterion alternative to that of Cutteridge ( 3489 of 1954) and based on conformal transformation is presented. See also 72 of February (Cutteridge).

### 621.375.232

1598
The Effect of Inverse Feedback on Input Impedance-J. B. Earnshaw. \{Radio Elect. Rev. (Wellington, N. Z.), vol. 9, pp. 37-40 and 34-35; December, 1954 and January, 1955.] Formulas are given for the input impedance, gain, and gain without feedback of twelve single-stage amplifiers a ad these, together with the circuits and their equivalents, are tabulated.
621.375 .3

1599
Alteration of Dynamic Response of Magnetic Amplifiers--R. O. Decker. [Elec. Engng. (New York), vol. 73, p. 1088; December, 1954.] Digest of paper to be published in Trans. Amer. IEE, Part I, Communication and Electronics, 1954; pp. 658-664. Analysis is presented indicating how a magnetic amplifier of full-wave self-saturating type can be made to exhibit phase lead or lag by aajusting the parameters of the feedback networks.

### 621.375.4:621.314.7

1600
Analysis of the Common-Base Transistor Circuit-F. Oakes. (Electronic Engng., vol. 27. pp. 120-126; March, 1955.) Simple equations are obtained by choosing a hybrid inverted- $\Pi$ network as the equivalent circuit for investigating the operation of the grounded-base point-contact transistor amplifier.

### 621.375.4:621.314.7

1601
D.C. Stability of Transistor CircuitsF. Oakes. (Wireless W'orld, vol. 6t, pp. 164167; April, 1955.) The design of amplifiers using
junction transistors is discussed, with particular reference to the influence of the base-tocollector leakage current, which increases rapidly with rising temperature. For stable operation, the change of collector current produced by a change of leakage current should be low; circuit arrangements for achieving this are indicated.

### 621.376.3:621.3.018.78:621.372.5 <br> 1602

The Distortion of F.M Signals in Passive Networks-R. H. P. Collings and J. K. Skwirzynski. (Marconi Rev., vol. 17, no. 115, pp. 113-136; 4th Quarter, 1954.) Expressions for the fundamental and first four harmonics of the output instantaneous frequency are given in a series expansion in terms of modulation frequency and modulation index, the coefficients in the expansion depending on the network parameters. Effects of detuning are considered, and detailed results for the Butterworth circuit are presented.
621.376.3:621.3.018.78:621.372.543.2 1603

The Linear Distortion of F.M. Signals in Band-Pass Filters for Large Modulation Fre-quencies-J. K. Skwirzynski, (Marconi Rev., vol. 17, no. 115 , pp. 101-112; 4th Quarter, 1954.) The linear distortion of a FM signal in a highly selective bandpass filter follows almost exactly the static response curve of the network, provided the following conditions are fulfilled:-(a) the modulation frequency is not less than two-thirds of the semi-bandwidth, (b) the modulation index does not exceed unity, (c) the network $Q$ is sufficiently large.
621.376.32:621.318.134

1604
A Ferrite Frequency Modulator-F. Slater. (Marconi Instrumentation, vol. 4, pp. 186-193, 200; December, 1954.) The modulator described comprises a ferroxcube-B4 ring carrying a rf winding and located in the gap of a Ni-Fe core carrying if and polarizing windings. It is suitable for operation over the frequency range from 400 kc to uhf, with slight modifications. Its demodulated distortion is 5 per cent at an nscillator frequency of 400 kc and a $15-\mathrm{kc}$ frequency deviation, and 2 per cent at 170 mc and 100 kc , respectively.

### 621.376 .332

1605
Discriminator Circuit Analysis-F. L. Morris. (Wireless Eng., vol. 32, pp. 93-98; April, 1955. Correction, ibid., vol. 32, p. 142; May, 1955.) Theory and performance figures are given for a simple frequency-discriminator circuit using an asymmetrical arrangement. The conversion efficiency compares favorably with that of the equivalent Foster-Seeley discriminator. It is particalarly recommended for purposes where high efficiency is of greater importance than accurate linearity over a wide frequency range.

### 621.396.049.75

1606
A Universal Printed Circuit-J. R. Goodykoontz. (Tele-Tech., vol. 13, pp. 74-75; December, 1954.) The universal printed-circuit board has a standard pattern of parallel wires on one side and is useful at the design stage, when prototypes are required for testing, since circuit changes may be made with relative ease.

### 621.372 .5

1607
Amplitude-Frequency Characteristics of Ladder Networks [Book Review]-E. Green. Publishers: Marconi's Wireless Telegraph Co., Chelmsford, 1954, 155 pp., 25s. (Brit, Jour. Appl. Phys., vol. 5, p. 457; December, 1954.) Suitable mainly as a reference work for telecommunication engineers.

## GENERAL PHYSICS

53.081

1608
Proposals for Units for Area, Electric Displacement and Magnetic Field StrengthP. Cornelius. (Philips Res. Rep., vol. 9, pp. 444-457; December, 1954.)
$537.21+538.12$
1609
The Two-Dimensional Magnetic or Electric Field of a Single Isolated Pole-Piece-N. H. Langton and N. Davy. (Bril. Jour. Appl. Phys., vol. 5, pp. 431-435; December, 1954.) Theoretical investigation for a pole-piece consisting of a thick plate terminated by a concave semicircular cylinder. Conformal transformations and elliptic functions are used. The variation of the field strength along the edge of the pole-piece and along the external axis of symmetry is calculated and shown graphically.
537.224

1610
Observation of the Costa Ribeiro Effect on the Dissolution of Naphthalene CrystalsE. Rodrigues. (Ann. Acad. Bras. Sci., vol. 26, pp. 381-383; December 31, 1954.) Report of experiments indicating that electric charges are developed on dissolving single crystals of naphthalene in a solvent of very low conductivity such as toluene; the crystals thus treated constitute natural electrets.
537.226:536.421.1:537.29 1611 Electro-fusion: a New Phenomenon observed in the Phase Changes of Dielectrics under the Influence of an Electric FieldJ. Costa Ribeiro. (Ann. Acad. Bras. Sci., vol. 26, pp. 349-355; June, 30 1954.) Experiments are described which indicate that application of a field between the electrodes of a capacitor with the dielectric partly in the solid and partly in the liquid state accelerates the phase change in the dielectric. From measurements of the current and calculation of the Joule energy dissinated, it is shown that this energy is several hundred times smaller than the heat necessary for the normal fusion of the corresponding mass of the dielectric.
537.226:536.421.1: $: 537.29$
Field-Induced Melting of $\quad 1612$ Field-Induced Melting of DielectricsB. Gross. (Ann. Acad. Bras. Sci., vol. 26, pp.
$289-291$; June 30, 1954.) Theory is presented relevant to the phenomena described by Costa Ribeiro (1611 above).

### 537.311 .1

1613
Plasma Oscillations in a Periodic Potential: the One-Zone Theory-J. Hubbard. (Proc. Phys. Soc., vol. 67, pp. 1058-1068; December 1, 1954.) The plasma-oscillation theory of collective electron interactions developed by Bohm and Pines ( 1375 of 1954 and back references) can be applied to the problem of the conduction electrons in metals by modifying it to take account of the periodic potential present. This modification is carried out, neglecting the interactions of electrons in different zones and assuming that the effect of interzone transitions on the collective behavior is small enough to be treated by perturbation theory. The main effect of the potential is to alter the effective mass of the electrons.
$537.311 .31+536.212 .2$
1614
The Electrical and Thermal Conductivities of Monovalent Metals-J. M. Ziman. (Proc. Roy. Soc. A, vol. 226, pp. 436-454; December 7, 1954.) "Numerical calculations for the case of sodium, using the Bardeen (1937) formula for the scattering cross-section and Blackman's (1951) value for a 'longitudinal Debye temperature,' agree better with observation than do the simple Bloch expressions, but there still remain discrepancies." Further study of the scattering cross-section function might remove these.
537.311.31 +537.312 .62
High-Frequency Resistance of Tin and Indium in the Normal and Superconducting State-C. J. Grebenkemper. (Phys. Rev., vol. 96, pp. 1197-1198; December 1, 1954.) Further measurements at 24 kmc confirm results obtained previously [ 3060 of 1952 (Grebenkemper and Hagen)].

The Physical Nature of a Metal Surface in Conduction Theory-H. A. Müser. (Phil. Mag., vol. 45, 1p. 1237-1246; December, 1954.) A discussion is presented in which the reflection of electrons striking the surface from inside is treated as a diffraction phenomenon.

## 537.5:061.3

1617
Gaseous Electronics Conference in New York-1. L. Jones. [ Nature (London), vol. 175, pl. 1.54-155; Jamary 22, 1955.] Papers presented at the conference held in October, 1954 are surveyed briefly. Physics of plasma and ionization processes leading to electrical breakdown were among the subjects discussed.

### 537.52

1618
Study of Gaseous Discharges by Magnetic Resonance-D. J. E. Ingram and J. C. Tapley. [Research (Iondon), vol. 7, supplement, pp. S6.3-S64; December, 1954.] Measurements are reported which indicate that the technique of magnetic resonance may prove useful for stadying the characteristics of low-pressure gas discharges.

### 537.523.5:621.396.822

1619
Relaxation Oscillations and Noise from Low-Current Arc Discharges-M. I. Skolnik and II. K. Puckett, Jr. (.Jour. Appl. Phys., vol. 26, pr. 74-79; January, 1955.) Noise measurements were made on ares, using various electrode materials and gases, over the frequency range $0.1 \mathrm{mc}-4.5 \mathrm{kme}$. The results for Al electrodes in air are plotted in comparison with the computed spectrum. Pulses generated in the circuit formed by the series limiting resistor and the capacitance across the discharge are considered responsible for most of the noise.

### 537.525

1620
Formative Time of the Cathodic Space Charge-I). Brini, O. Rimondi and $P$. Veronesi. ( Nuovo Cim., vol. 12, pp. 915-922; December 1, 1954. In English.) A criterion for evaluating the formative time is based on observations of rise time and overshoot in intermittent discharges. Different kinds of intermittent discharges are reviewed.

### 537.525

1621
On Intermittent Discharges in Air at Low Pressure-D. Brini, O. Rimondi and P. Veronesi. (Nuovo Cim., vol. 12, pp. 948-949; December 1, 1954. In English.) An experimental investigation is reported of the dependence of the discharge frequency on the gas pressure and electrode separation.

### 537.533

1622
Forces between Parallel Electron Streams -J. Webb. [Elec. Rev. (London), vol. 155, pp. 1037-1038; December 31, 1954.) A simplified fundamental analysis is presented, taking account of the forces of acceleration, es repulsion, em attraction, gravitation and any local fields. The results indicate that for all electron velocities up to the immediate threshold of the velocity of light the resultant force between parallel electron streams is repulsive.

### 537.56

1623
Spectroscopic Studies of Highly Ionized Argon produced by Shock Waves-II. E. Petschek, P. II. Rose, II. S. Glick, A. Kane and A. Kantrowitz. (Jour. Appl. Phys., vol. 26, pp. 83-95; January, 1955.)

Kinetic Theory of Weakly Ionized Homogeneous Plasmas: Part $1-\mathrm{M}$. Bayet, J. I. Delcroix and J. IF. Denisse. (Jour. Phys. Radium, vol. 15, pp. 795-803; December, 1954.) A rigorous method is presented for integrating Boltzinann's equation by successive approximations for a weakly ionized gas subjected to an alternating electric field and a static magnetic field. The electron-velocity
to the electric field, and the individual functions obtained are developed as spherical functions. The method is applied to investigation of a Lorentz-type gas. See also 2910 of 1954.

### 537.56:537.311.37

1625
Electrical Conductivity of Highly Ionized Argon produced by Shock Waves-Shao-Chi Lin, E. L. Resler and A. Kantrowitz. (Jour. Appl. Phys., vol. 26, pp. 95-109; January, 1955.)

### 538.22

1626
Antiferromagnetism and Ferrimagnetism of Non-Stoichiometric Compounds-E. W. E1cock. (Proc. Roy. Soc. A, vol. 227, pp. 102114 ; December 21, 1954.) A simple quantitative treatment is given of the magnetic properties of a substance containing vacancies, enabling the most important magnetic properties of many nonstnichionetric compounds to be interpreted.

## 538.3

1627
Representation of Electromagnetic Fields of Any Frequency using the Energy-Quantum Model-H. Zuhrt. (Arch. Elekt. Überlragung, vol. 8, np. 565-577; Decenber, 1954.) Develop)ment of theory presented previously (1004 of May). Wave propagation phenomena including reflection, interference and diffraction are explaited in terms of the energy-quantum model; static fields are also considered. Voltage, current and characteristic impedance are related to mechanical quantitics, and an appropriate system of dimensions and units is presented.

### 538.561:537.533 <br> 1628

Čerenkov Radiation from Extended Electron Beams-M. Danos. (Jour. Appl. Phys., vol. 26, pp. 2-7; January, 1955.) Calculations are made of the radiation emitted by bunched beams passing close to dielectric surfaces; both plane and cylindrical geometries are considered. Sce also Trans. IRE, vol. MTT-2, no. 3, pp. 21-22; September, 1954.

### 538.566:535.42

1629
Calculated Diffraction Patterns of Dielectric Rods at Centimetric WavelengthsC. Froese and J. R. Wait. (Canad. Jour. Phys., vol. 32, pp. 775-781; December, 1954.) Calculations of the field behind the cylinder are made for the case of a normally incident wave with the electric vector (a) parallel and (b) perpendicular to the cylinder axis. The dielectric materials considered include polystyrene, lucite and tenite. The diffraction pattern is only slightly dependent on the dielectric constant of the rod. Small discrepancies between the calculated values and values obtained experimentally [3538 of 1954 (Wiles and McLay)] are probably due to resistive loss in the dielectric.

### 538.569 .4

1630
Absorption of Microwaves by Orygen in the Millimeter Wavelength Region-J. O. Artman and J. P. Gordon. (Phys. Rev., vol. 96, pp. 1237-1245; December 1, 1954.) An account is given of an experimental investigation of the absorption at high and at low pressure. Theories advanced in explanation of the linebroadening effects are discussed.
538.569 .4

1631
Spectral Investigations in the Wavelength Range around 1 mm -L. Genzel and W. Eckhardt. (Z. Phys., vol. 139, pp. 578-598; December $20,1954$. ) A description is given of the construction and method of operation of an infrared-type spectronteter which has been used with thermal radiator and receiver to obtain absorption spectra at wavelengths over 1 mm . The spectrum obtained for water vapor over the range $0.14-1.4 \mathrm{~mm}$ is shown as an example. Results obtained with HCN and $\mathrm{H}_{2} \mathrm{~S}$ are also reported.

Spectrometer-R. J. Collier. (Rev. Sci. Instr., vol. 25, pp. 1205-1207; December, 1954.) Molecular-gas spectra at wevelengths around 10 cm are investigated using a system in which a single coaxial cavity serves as Stark or Zeeman modulation cell and as frequency reference unit.
548.0:53

1633
Simplified Impurity Calculation-G. F. Koster and J. C. Slater. (P/iys. Rev., vol. 96, pp. 1208-1223; December 1, 1954.) The methods developed previousiy ( 699 of March) are used to investigate the case ol a local perturbation in a simple cubic lattice.

## GEOPHYSICAL AND EXTRATERRESTRIAL PHENOMENA

523.16

1634
Abnormal Galaxies as Radio SourcesB. Y. Mills. (Observatory, vol. 74, pp. 248-249; December, 1954.) Results of radio-telescope observations at Sydney indicate that galaxy NGC 1316 is probably a radio source, whereas NGC 1947 is probably not.
523.16

1635
Observations of Galactic Radiation on a Wavelength of 33 cm -J. F. Denisse, É. Leroux and J. L. Steinberg. [Compt. Rend. Acad. Sci. (Paris), vol. 240, DF. 278-280; January 17, 1955.] Report ó measurements made using the Warzburg mirror at Marcoussis. Results are presented in the form of an isophot map of the galaxy with intensities expressed as apparent temperature in arbitrary units. The $33-\mathrm{cm}$ radiation is attributed principally to ionized hydrogen. A table giving the positions of localized sources includes five thought to have been observed for the first time.

### 523.16:523.72

1636
Study of Solar R.F. Radiation on 9350 $\mathrm{Mc} / \mathrm{s}$ around Sunset and Sunrise-I. Kazes and J. L. Steinberg [Compl. Rend. Acad. Sci. (Paris), vol. 240, 1ŋ. 493-495; January 31, 1955.] Report of measurements made at Paris, using a parabolic mirror of diameter 1.5 m . Graphs are presented showing the variation of refraction with solar angle and quasi-periodic intensity variations observed when the angle of elevation is less than 15 degrees. Attenuation due to absorption by oxygen is observed at angles less than about 20 tegrees.

### 523.16:523.72 1637 <br> Fine Structure of Solar Radio Transients-

 G. Reber. [ Nature (Londen), vol. 175, p. 132; January 15, 1955.] Brief report of observations made during 1948-1950. Solar bursts observed at 480,160 and 51 mc were found to be composed of numerous discrete pips with median duration approximately proportional to wavelength. The spectral width of a pip is a few per cent of its mean frequency. The frequency of occurrence of the pips was greatest at 160 mc.523.16:621.396.677

1638
Aerial Smoothing in Radio AstronomyR. N. Bracewell and J. A. Roberts. (Aust. Jour. Phys., vol. 7, pp. 615-640; December, 1954.) Theoretical considerations show that the antenna does not register those spatial Fourier components of the true distribution of radio brightness having frequencies beyond a cut-off determined by the antenna aperture. Components of lower frequency are registered but their relative strengths are altered. The consequences are that (a) there are invisible distributions which produce no response when scanned by the antenna, and (b) in conducting a survey the measuring points must be closer together than half the period of the fonurier component at cut-off.
523.74/.75

1639
Prominence Activity (1905-1952)-R. Ananthakrishnan. (Proc. Indian Acad. Sci.,

Charts prepared from observations made at Kodaikanal show the salient features of solar prominence activity during the last four sunspot cycles. A preliminary examination suggests that during the maximum phase of the cycle geomagnetic activity shows a better correlation with promine aces than with sumspots.

### 523.752

The Emission of Radiation from Mode Hydrogen Chromospheres: Part 2-J. T. Jefferies and R. G. Gioranelli. (Aust. Jour. Phys., vol. 7, pp. 574-585; December, 1954.) An improved method is presented for calculating the characteristics of the radiation field of $\mathrm{II} \alpha, \mathrm{L} \alpha, \mathrm{L} \beta$, and the Lyman continuum emitted by model hydragen atmospheres at kinetic temperatures of $10^{4}-2.5 \times 10^{5}$ degrees K. A useful application of the results would be in interpreting observations of prominences and flares.
550.385
Qualitative Explanation of the Commencement of Some Polar Magnetic Disturbances based on the Theory of Chapman and Ferraro -G. Grenct. [Compl. Rend. Acad. Sci. (Paris), vol. 240, pp. 448-450; Jan'sary 24, 1955.]
550.385:550.375

1642
Geomagnetic and Geoelectric VariationsJ. G. Scholte and J. Veldkamp). (Jour. Atmos. lerr. Phys., vol. 6, pp. 33-45; January, 1955.) Analysis of records of geamagnetic pulsations, especially those with periods of $10-10^{2}$ seconds, indicates that they may be associated with ionospheric vibrations caused by a disturbance of the ionization equilibrium. Information regarding the distribution of ground conductivity is obtained from the relation between the magnetic pulsations and the associated transverse-electric field pulsations.
550.385.25:551.510.535

1643
The Abnormal Variations of the Horizontal Magnetic Intensity at Huancayo (Peru)-Z. Ibrahim. (Proc. Math. Píys. Soc. Egypt, vol. 5, p1. 21-24; 1953.) The magnetic field of the ionospheric current system arising from lunar atmospheric tidal motion is calculated for various points on the earth's surface. The daily variation of the horizontal component at IIuancayo is found to be nearly twice as great as it Batavia, though the latitudes of the two points are nearly the same. Observations at a point between the geographic and magnetic equators in the northern hemisphere shoutd help to explain this anomady.

### 550.385.4:551.510.535

1644
On the World-Wide Disturbance in $\mathrm{F}_{2}$ Region-T. Obayashi. (Jour. Geomag. Geoelect., vol. 6, pp. 57-67; June, 1954.) The average disturbance in the value of $f_{0} F_{2}$ during magnetic storms is separated into the storm-time part $D_{s i}\left(f_{0} F_{2}\right)$ and the local-time part $D_{s}\left(f_{0} F_{2}\right)$. These were calculated in the cases of 10 typical storms from data obtained from 40 stations in the northern hemisphere. Both their range and phase are correlated with the magnetic activity. The phase of $D_{*}\left(f_{0} F^{2}\right)$ during the active stage of the magnetic storm is almost entirely dependent on the local time but after the cessation of activity the disturbance moves with the rotating earth. See also Jour. Radio Res. Labs., Japan, vol. 1, pp. 41-50; June, 1954.

## $550.386^{\prime \prime} 52^{n}$

1645
An Analysis on the Diurnal Variation of the Terrestrial Magnetism, especially on the Day-Time-Variation of Geomagnetically Quiet Days -M. Ota. (Jour. Geomag. Geoelect., vol. 6, pp. 83-98; June, 1954.)

### 551.510 .534 <br> 1646

Vertical Distribution of Atmospheric Ozone at Longyearbyen, Spitzbergen ( 78 degrees N) -S. H. H. Larsen. (Jour. Almos. Terr. Phys., vol. 6, pp. 46-49; January. 1955.)
551.510 .535

1647
The Constitution of the Upper Atmosphere and the Ionosphere Research Station of the Institute of Geophysics at Genoa-M. Bossolasco and A. Masotto. (Geofis. Mel., vol. 2, rp. 80-86; September/I ecember, 1954.) An outline of knowledge on the structure and electrical properties of the ionosphere is followed by a tescription of ionosphere sounding equipment recently put into operation and comprising pulse transmitter for the range $2-15 \mathrm{~mm}$, double-superheterodyne receiver and cro recorder.

### 551.510 .535

1648
Calculation of the Collision Frequency in the Ionosphere-L. Caprioli. (R.C. Accad. naz. I.incei, vol. 17, pp. 365-370; December, 1954.) It is shown that, if $z_{0}$ is the true reflection height corresponding to $\omega_{0}$, the lowest frequency of the sounding sweep, and $z_{E}$ is the true reflection height corresponding to the first critical frequency $\omega_{\mathcal{R}}$ greater than $\omega_{0}$, then provided that the distribution of ion concentration and collision frequency are known for values of $z u p$ to $z_{0}$ and the absorption is observed over the range $\omega_{0}-\omega_{E}$ the function $\nu(z)$ expressing the height distribution of collision frequency can be determined over the range $z_{0}-z_{E}$ by solving an integral equation due to Abel.

### 551.510 .535

1649
Turbulence in the Upper IonosphereA. Maxwell. (Phil. Mag., vol. 45, pp. 12471254 ; December, 1954.) "From the experimental data at present available it is shown that the Reynolds number in the upper $F$ region (300400 km level) is of the order of 300 . The region may therefore be turbulent. It is suggested that the high level diffracting screens which give rise to spread $F$ echoes and to radio star fading are caused by non-laminar flow, and that their non-appearance during the daylight hours may be due to the inhibition of turbulence by large temperature gradients, by lower drift velocities, or by an increase in the kinematic viscosity."

### 551.510 .535 <br> 1650

Motion of Clouds of Abnormal Ionization in the Auroral and Polar Regions-E. L. Hagg and G. H. Ianson. (Canad. Jour. Phys., vol. 32, pp. 790-798; December, 1954.) A study has been made of unusual types of echo exhibited by film records of ratpid-succession sweep observations at several stations in Northern Canada. Three distinct types of echo are identified, probably corresponding to (a) horizontally moving $F_{s}$ clouds, (b) clouds descending vertically from the $F$ to the $E$ layer, and (c) clouds moving at extremely high velocities in the $E$ layer. Type (c) echoes were observed only at stations very close to the auroral zone maximum and may be due to sweeping of the auroral ionizing agent. Type (b) echoes also appear to be related to an auroral ionizing agent.

### 551.510 .5351651

Information obtained from Ionization Charts-R. Eyfrig, E. Harnischmacher and K. Rawer. [Compl. Rend. Acad. Sci. (Paris), vol. 240, pp. 446-448; January 24, 1955.] Ionization charts drawn up monthly for two zones covering (a) Arnerica and (b) Europe, Africa and Asia are discussed. The distributions observed generally confirm the theory of geomagnetic control of the $F_{2}$ layer [2898 of 1946 (Appleton)], but some exceptions are observed; there appears to be a lower limiting value for the height of the sun below which the geomagnetic control does not operate. The control is effective at latitudes as high as 70 degrees. Measurements at other times besides midday are required to study the effect. Asymmetry as between northern and southern hemispheres is observed. The north polar zone was studied in detail; the auroral $E$ zone appears to be centered on the geographic pole.
551.510 .535

1652
Velocity of Movement of Sporadic-E Clouds-M. R. Kundu. (Sci. and Cull., vol. 20, D. 303; December, 1954.) Curves of the diurnal variation of the $E_{n}$-layer critical frequency for four Jamanese stations indicate a progressive retardation of the time of occursence of the maximum on passing from the highestlatitude to the lower-latitude stations. A value of about $70 \mathrm{~m} / \mathrm{s}$ is deduced for the horizontal component of the velocity of the $E_{0}$ clouds; this is of the same order as the velocity of $E$ layer winds found by various methods.
551.510 .535

1653
Interpretation of Measurements on the Ionosphere $\mathrm{F}_{1}$ layer-K. Rawer. [Compl. Rend. Acad. Sci. (Paris), vol. 240, एD. 331-333; January 17, 1955.] The maximum-ionization level of the $F_{1}$ layer is often located within the lower part of the $F_{2}$ layer, so that echo sounding yields a continuous curve for the $F_{1}$ and $F_{2}$ layers, with a more or less marked maximum or merely a point of inflection indicating the virtual height of the $F_{1}$ layer. In consequence, the critical frequency of this layer is not clearly defined. A correction is required, the sign of which depends on the theory accepted for the origin of the $F_{2}$ layer. Mohlers theory (1014 of 1941) gives results in good agreement with observations.
551.510 .535

1654
Lunar Tides in the Ionosphere $F_{2}$ Layer at Dakar-F. Delobeau. [Compl. Rend. Acad. Sci. (Paris), vol. 240, pp. 222-224; January 10, 1955.] Hourly records of $f_{0} F_{2}$ and $h^{\prime} F_{2}$ for the years 1950-1953 have been analyzed: maximum variations and times of occurrence of maxima are tabulated for the different seasons. The semidiurnal component is predominant except in summer; during winter the tidal variations of $f_{0} F_{2}$ and $h^{\prime} F_{2}$ have amplitudes of $\pm 0.2 \mathrm{mc}$ and $\pm 5 \mathrm{~km}$ respectively. The phases and amplitudes are considerably different from those observed at other tropical stations; anomalies may be due to the circumstance that the $F_{2}$ layer at Dakar is very thick in summer; the ionization maximum occurs at the equinoxes and not in winter.
551.510 .535
On the Meridional Distribution of the
Minimum Virtual Height of the Fs Minimum Virtual Height of the $F_{2}$ LayerT. Shimazaki. (Jour. Radio. Res. Labs. Japan, vol. 1, pp. 15-25; March, 1954.) Results of an analysis of daily variations in $h^{\prime} F_{2}$ from 35 observatories indicate the presence of two periodic terms: (a) the diurnal variation, which is dependent on the geographic latitude, and (b) the semidiurnal variation, which depends on the geonagnetic latitude. The former may be considered as the thermal effect, the latter as the tidal effect in the $F_{2}$ region.

### 551.510 .535

1656
The Effect of the Solar Tides and the Temperature Change on the Daily Variation in Electron Density and Height of the F2-Layer T. Shimazaki. (Jour. Geomag. Geolecl., vol. 6, pp. 68-82; June, 1954.) The departure $\Delta n_{m}$ of the maximum electron density from the norm of a static Chapman region is inversely proportional to the temperature variation, excent in the morning when the electron density is small. The semidiurnal variation in $\Delta n_{m}$ appears even in the case when account is taken of the diurnal temperature variation only. As regards $\Delta n_{m}$, the departure from the norm of the height of $n_{m}$, the result is similar to that obtained by Weiss ( 1.348 of 1953) except that it is more strongly temperature dependent. There is some interaction between the temperature and the tidal effects.

### 551.510.535:523.3

Influence of the Moon on the Maximum Ionization of the Ionosphere $E$ LayerE. Harnischmacher. [Compt. Rend. Acad. Sci.
(Paris), vol. 240, pp. 553-555; January 31, 1955.] An analysis of $f_{0} E$ variations over a period of four years, observed at six stations, indicates that after taking account of seasonal, solar-cycle and equation-of-time effects, there is evidence of a lunar variation of amplitude 1 per cent centered on the full moon.

### 551.510.535:523.72 <br> 1658

Soft X Radiation from the Quiet Solar Corona-G. Elwert. (Z. Nalurf., vol. 9a, pp. 637-653; July/August, 1954.) The mechanisms involved in the production of the soft $X$-rays are discussed, and the significance of this radiation for the formation of the normal ionosphere $E$ layer is indicated. See also 726 of 1954.

### 551.510.535:621.317.3

1659
Equipment for the Measurement of Changes of the Phase Path of Ionospheric Echoes-H. Yuhara, T. Koseki and Y. Aono. (Jour. Radio Res, Labs, Japan, vol. 1, pp. 1114; March, 1954.) The equipment described is built with units conventional in pulse applications and is developed from that described by Findlay (397 and 404 of 1952). Oscillograms of ground pulse beats and echo patterns of the $E$ and $F$ regions are shown.

### 551.510.535:621.396.11

1660
Reflection Conditions for Vertical Propagation in the Ionosphere in the Presence of Collisions and of the Earth's Magnetic Field. Case of the E Layer-Lepechinsky and Durand. (See 1767.)
551.510.535:621.396.11.029.51

The Development of an E-Region Model consistent with Long-Wave Phase-Path Meas-urements-R. E. Jones. (Jour. Almos. Terr. Phys., vol. 6, pp. 1-17; January, 1955.) An electron-density model is evolved by modification of the Chapman theory to include the effects of variable scale height, dissociation of $\mathrm{O}_{2}$, and variable recombination. The model is checked against phase-height data for 150 kc and 2.4 mc and against $f_{0} E$ values.

### 551.510.535:621.396.812

1662
Measurement of Attenuation in the Iono-sphere-Ochs. (See 1764.)

### 551.510.535:621.396.812.3 <br> 1663 <br> The Autocorrelation of Randomly Fading Waves-Banerji. (See 1772.)

551.510.535(98): 621.396.11

1664
Statistical Studies of Polar Radio Black-outs-J. W. Cox and K. Davies. (Canad. Jour. Phys., vol. 32, pp. 743-756; December, 1954.) "A statistical study of high frequency radio blackouts in Canada is made from records taken at several ionosphere sounding stations. Both vertical incidence and communication data are examined to determine the geographical, seasonal, and diurnal distributions of the frequency of occurrence of blackout. It is found that blackouts are most abundant in the morning and that the time of maximum occurrence increases with increasing latitude."

### 551.578:621.396.11:621.396.96 <br> 1665

The Microwave Properties of Precipitation Particles-K. L. S. Gunn and T. W. R. East. (Ouart. Jour. Roy. Mel. Soc., vol. 80, pp. 522545 ; October, 1954.) "The theory of scattering and attenuation by rain, snow and cloud is reviewed and theoretical results are presented in the form of equations, tables and graphs, so that the radar response to meteorological particles can be calculated at six wavelengths ( 10 , $5.7,3.2,1.8,1.24$ and 0.9 cm ) and various temperatures. Particular emphasis is placed on developments since Ryde's comprehensive paper in 1946. Published experimental results are compared with the theory. All results computed from the theory are contained in Tables 4 and 5 . The attenuation by water vapour and oxygen is given in an Appendix."
551.594.6 1666

An Attempt to observe Whistling Atmospherics near the Magnetic Equator-J. R. Koster and L. R. O. Storey. [ Nalure (London), vol. 175, pp. 36-37; January 1, 1955.] According to the theory of Barkhausen and Eckersley, the mode of propagation of whistling atmospherics is such that they should not occur near the magnetic equator. A report is given of observations made at Achimota, over the period from December, 1951 to March, 1954; no whistlers were detec ed, though other types of atmospheric were frequent.

## LOCATION AND AII'S TO NAVIGATION

621.396.93:551.594.6

1667
Low-Frequency Direction Finder-C. Clarke and V. A. W. Harrison. (Wireless Eng., vol. 32, pp. 109-114; A ril, 1955.) A more detailed account of the instrument described previously [ 3226 of 195 (Horner)].

### 621.396 .933

1668
Radio Installations of the Danish Airways System-K. Svennings :n. [Teleteknik (Copenhagen), vol. 5, pp. 391-400; December, 1954.] The navigation aids and communication systems used are described.
$\begin{array}{rr}\text { 621.396.96:551.578:621.396.11 } & 1669 \\ \text { The Microwave Properties of Precipitation }\end{array}$ Particles-Gunn and Eist. (See 1665.)
621.396.962.33

1670
Radar Receiver with Elimination of FixedTarget Echoes-H. Tanter. (Elec. Commun., vol. 31, pp. 235-248; December, 1954.) English version of paper abstrat ted in 2407 of 1954.

## MATERIALS AND SUBSIDIARY TECHNIQUES

## 535.5

1671
Design and Operaticn of Evapor-ion Pumps -R. H. Davis and A. S. Divatia. (Rev. Sci. Instr., vol. 25, pp. 1193-1197; December, 1954.) Operation of the pump described depends on the gettering action of continuously evaporated Ti in conjunction with is n pumping. The lowest pressure attainable is about $2.10^{-7} \mathrm{~mm} \mathrm{Hg}$. The dependence of pumping speed on the temperature of the getterin\& surface, the pressure, and the rate of evapor ation of Ti is investigated.

### 535.215:537.311.33:546.817.23

1672
Response Time of Photoconductivity of Lead Selenide-L. Sosrowski and M. Chmielewski. (Bull. Acad. Po'on. Sci., Classe 3, vol. 1, nos., 3/4, pp. 119-121; 1953. In English.) An oscillographic method for investigating response times of less tha $1 \mu \mathrm{~s}$ is described. An exponential timebase is used and the specimen is illuminated in synchronism by light pulses at repetition rates up to 50,000 per second. The response times of three different PbSe cells were $0.25,0.35$ and $0.9 \mu \mathrm{~s}$ within $\pm 0.1 \mu \mathrm{~s}$, their respective resistances and sensitivities being 23 , 51 , and $100 \mathrm{k} \Omega$, and 12,15 , and 40 arbitrary units.
535.215:537.311.33:546.817.23

1673
Photoconductive an 1 Photovoltaic Layers of Lead Selenide- $I f$, Checifiska. (Bull. Acad. Polon. Sci., Class' 3, vol. 1, nos. 3/4, pp. 123-135; 1953. In English.) The method of preparing PbSe layers 'xhibiting these effects is described and some results of determinations of limiting sensitivity and spectral sensitivity in the range $0.5-3.6 \mu$ a e given. At room temperature the sensitivity to radiation from an ordinary incandescent lamp is less than one tenth that of PbS .

### 535.37

1674
The Shape of the Enission Bands of Luminescent Solids-C. C Vlam. (Bril. Jour. Appl. Phys., vol. 5, pı. 443-446: December, 1954.)
535.37 1675
Luminescence in High Polymers-H. Hinrichs. (Z. Naturf., vol. 9a, pp. 617-630; July /August, 1954.) An investigation of organic phosphors embedded in polystyrol.
535.37:537.311.33

1676
Temperature Dependence of the EnergyGap in ZnS - C. Z. van Doorn. (Physica, vol. 20, pp. 1155-1156; December, 1954.) Measurements on a single crystal showed that the temperature variation of the energy gap varied between $4.6 \times 10^{-4} \mathrm{ev}$ degrees K at 77 degrees $\mathrm{K}^{-1}$ and $8.5 \times 10^{-4} \mathrm{ev}$ degrees K . at 800 degrees K .
535.3761677

The Edge Emission of $\mathrm{ZnS}, \mathrm{CdS}$ and ZnO and its Relation to the Lattice Vibrations of these Solids-F. A. Kröger and H. J. C. Meyer. (Physica, vol. 20, [f*. 1149-1156; December, 1954.)

### 535.376:538.615

1678
The Effect of Intense Magnetic Fields on Electroluminescent Powder Phosphors-A. N. Ince. (Proc. Phys. Soc., vol. 67, pp. 870-874; December 1, 1954.) No quenching of electroluminescence was observed in phosphors in magnetic fields of up to $1.3 \times 10^{5}$ oersted. This is contrary to the prediction: of Destriau ( 110 of 1949). The significance of this result is discussed.
537.226

1679
High-Frequency Polarization of a Spherical Body and of an Assemblage of Particles of a Perfect Dielectric-A, Colombani. [Compl. Rend. Acad. Sci. (Paris), vol. 240, pp. 503 505; January 31, 1955.] Analysis is given for a single particle, using Maxwell's equations in spherical-coordinate form. If the permittivity is real and very high the electric field and es energy within the particle may become very large. An expression is derived for the apparent susceptibility of an assemblage of the particles.
537.227:546.431.824-31 1680

Observation of Paramagnetic Resonances in Single Crystals of Barium Titanate A. W. Hornig, E. T. Jaynes and H. E. Weaver. (Phys. Rev., vol. 96, p. 1703 ; December 15, 1954.)

### 537.311 .33

1681
A Possible Mechanism of the Scattering of Current Carriers in Semiconductors-T. A. Kontorova. (Zh. Tekh. Fiz., vol. 24, pp. 22172220; December, 1954.) The unusual form of the temperature dependence of the mobility of current carriers for certain semiconductors has not yet been satisfactorily explained. By taking account of the interaction of the current carriers not only with acoustic vibrations of the lattice but also with optical vibrations which are possible when the adjacent particles vibrate in opposite phase, the theoretical results are brought into agreement with experiment.

### 537.311 .33

1682
Use of $p-n$ Junctions for Solar Energy Con-version-E. S. Rittner. (Phys. Rev., vol. 96, pp. 1708-1709; December 15, 1954.) Calculations of the efficiency to be expected from a Si cell lead to results in agreement with those of Cummerow ( 145 of January ${ }^{\dagger}$, except that for maximum power conversion efficiency the energy gap of the semiconductor should be about $1.5-1.6 \mathrm{ev}$; thus AlSb is preferable to Si in this respect.
537.311 .33

1683
Determination of the Optical Constants of Type-AIIIBV Semiconductors at Infrared Wave-lengths-F. Oswald and R. Schade. (Z. Naturf., vol. 9a, pp. 611-617; July/August., 1954.) Reflection and transmission measurements were made on a series of compounds of elements of Groups III and V of the periodic system, at wavelengths from 0.8 to $15.2 \mu$. Values found for absorption constant, refractive index and
width of energy gap are tabulated; energy-gap values determined by other workers, using electrical methods, are given for comparison

### 537.311 .33

1684
Theory of Conduction in Isotropic Semicon-ductors-O. Madelung. (2. Nathrf., vol. 9a, pp. 607-674; July/August, 1954.) The classical theory is developed and the equations for electrical and thermal current density are established. From these equations the coefficients of all the galvanomagnetic, thermoelectric and thermomagnetic effects are derived. Expressions for these coefficients are tabulated for small values of magnetic: field strength, for degenerate and nondegenerate semiconductors with given scattering mechanism.
537.311.33:537.323

1685
Theory of the Thermoelectric Power of Semiconductors - C. Ilerring. (Phys. Rev., vol. 96, Di. 116.3-1187; December 1, 1954.) The marked rise of thermoelectric power which has been observed in some semiconductors at temmeratures below room temperature [1093 of 1954 (Freclerikse)] has been attributed to the effect of thermal lattice vibrations, the total thermoelectric power being composed of an electron term and a plonon term. This theory is developed; the results are supported by experimental data on Ge.

### 537.311.33:546.23

1686
Crystallization of Selenium under Pressure -D. N. Nasledov and P. T. Kozyrev. (Zh. Tekh. Fiz., vol. 24, pp. 2124-2135: December, 1954.) The effect of pressure up to 4,000 atm was investigated. I'ressure retarels crystallization at low temperatures below 110 degrees C .) and accelerates it at higl temperatures (above 170 degrees C.). A curve shows the relation between temperature and pressure at the point of fusion. large single crystals (up to 5 mm ) were formed from liquid Se under pressure. The electrical properties of the crystallized material were investigated. I'ossible mechanisms explaining the observed phenomena are discussed.
537.311.33:546.24

1687
Tellurium Single Crystals prepared by the Czochralski Process-J W'eidd. (Z. Nalurf. vol. 9a, p. 697; July/August, 1954.)

### 537.311.33:[546.28+546.289 1688

Theory of Donor and Acceptor States in Silicon and Germanium-C. Kittel and A. II. Mitchell. (Ihys. Rev., vol. 96, pp. 1488-1493; December $15,1954$. ) 'T'e applicability of the Wannier equation to the donor states in Si and Ge is examined with patticular reference to the multiple energy minima in the conduction band. The theory is extended to include degenerate bands, and it is shown that the Wannier equation is to be replaced by a set of coupled wave equations. The theory is also applied to acceptor states. Theoretical and experimental results are in fair agreement for both donors and acceptors.
537.311.33:546.28

1689
Single Crystals and $p-n$ Stratified Crystals of Silicon-H. Kleinknecht and K. Seiler. (Z. Phys., vol. 139, np. 539-618; December 20, 1954.) Measurements are reported on specimens produced as described previously [422 of 1953 (Kleinknecht)]. The voltage rlependence of the capacitance of the $p-n$ specimens indicates a linear variation of impurity-center concentration. The diffusion voltages and the slope of the forward chatracteristic are not consistrnt with a carrier cencentration conforming to the Boltzmann law at the boundary layer. The backward currents do not fit Shockley's theory; the discrepancy is explained by assuming the presence of traps acting as recombination centers. The concentration and effective cross section of these is evaluated from the increase of capacitance at low frequencies.
537.311.33:546.28:621.314.7

1690
Electronic Behaviour of Certain Grain Boundaries in Perfect Crystals-H. F. Mataré. (Z. Naturf., vol. 9a, ค. 698; July/August, 1954.) Some experiments are reported on the effect of incorporating structural inhomogeneities in semiconductor crystals; grain boundaries with highly nonlinear resistance were studied. A current gain of 50 was obtained with a threeelectrode arrangement using a $5-\Omega \mathrm{cm}$ Si specimen. A practieal form of grain-boundary transistor is illustrated.
537.311.33:546.289

1691
The Production of Germanium from Zinc Residues-[.Metallurgia (Manchester), vol. 50, pp. 27ī-278; December, 1954.] Some details are given of methods used in industry and in laboratory assay work for extraction of (ie and for the preparation of the metal. A different recovery scheme is required for each different complex residue. loison hazards are mentioned.
537.311.33:546.289 1692

An Observation of Circular Patterns in the Vicinity of Small-Area Alloyed Germanium $p-n$ Junctions-N. Holonyak, Jr. (Jour. Appl. Phys., vol. 26, pp. 121-123; January, 1955.) Rings observed on the crystal surface in the neighborhood of small alloyed junctions are explained in terms of the action of bubbles accompanying the etching process
537.311.33:546.289 1693
Optical Studies of Injected Carriers: Part 3Infrared Absorption in Germanium at Low Temperatures-R. Newman. (Phys. Kev., vol. 06, pp. 1188-1190; December 1, 1954.) Measurements are reported which confirm that carriers produced by injection and by impurity doping produce the same absorption effect. Anomalies observed in samples containing Fe may be due to trapping effects. Part 2: 1084 of 1954
537.311.33:546.289 1694

A Photoelectric Method for the Simultaneous Determination of Lifetime and Mobility of Injected Current Carriers in Semicon-ductors-G. Adam. (Z. Naturf., vol. 9a, pp. 607-611; July/August, 1954.) Additional carriers are formed by illumination with an elongated spot of light which is swept along the specimen. A probe in the vicinity of the illuminated region picks up a voltage proportional to the additional carrier concentration; this is recorded oscillographically. Methods of evaluating the oscillograms are described and illustrated by examples. Carrier lifetime and diffusion constant can be determined from a single oscillogram. Results are presented for Ge specimens containing various impurities.

### 537.311.33:546.289

1695
Determination of the Relation between Mobility and Diffusion Coefficient for Photoholes in $n$-Type Germanium-S. M. Ryvkin. (Zh. tekh. Fiz. vol. 24. pp. 2136-2149; December, 1954.) Detailed report of an experimental investigation of the diffusion and drift of photoelectrically jroduced holes. The results confirm the theoretical prediction that Einstein's formula (1) relating mobility to diffusion coefficient applies in this case. The stationary clistribution of minority carriers in a partially illuminated semiconductor is discussed in an appendix.

### 537.311.33:546.289

1696
Resistivity and Hall Effect of Germanium at Low Temperatures - C. S. Hung and J. R. Gliessman. (Phys. Rev., vol. 96, pp. 12261236; December 1, 1954.) Report of an extensive experimental investigation at temperatures from room temperature to that of liquid IIe. Anomalies in the Hall-constant curves at low temperature are explained on the assumption of small but finite mobility of carriers in the impurity states; the contribution of these carriers to the total conduction becomes impor-
tant at low temperature because the concentration of carriers in the conduction band is then very low.
537.311.33:546.289 1697
Transverse Hall and Magnetoresistance Effects in $p$-Type Germanium-R. K. Willardson, T. C. Harman and A. C. Beer. (Phys. Rev., vol. 96, pp. 1512-1518; December 15, 1954.) Calculations based on modification of the two-band model to take account of a small number of high-mobility holes give values for the magnitude, temperature dependence and magnetic-field dependence of llall and magnetoresistance effects in good agreement with experimental results. The importance of making measurements both at large and small magnetic-field strengths is indicated.

### 537.311.33:546.431-31:535.215

1698
Ultraviolet Absorption in Barium Oxide Films-K. Okumura. (Phys. Rev., vol. 96, pp. 1704-1705; December 15, 1954.) Measurements are reported briefly; results are compared with those of Tyler and Sproull (148 of 1952).
537.311.33:546.48.241.1

1699
Semiconducting Cadmium Telluride-D. A. Jenny and R. 1F. Bube. (Phys. Rev., vol. 90, pp. 1190-1191; December 1, 1954.) In general, $n$ type specimens are obtained by adding GroupIII or Group-VII impurities and p-type by adding Group-I or Group-V, the activation energies of the $p$-type impurities being much larger than those of the $n$-type impurities. The intrinsic energy gap is about 1.45 ev . Electron and hole mobilities are at least 30 cm per $\mathrm{v} / \mathrm{cm}$.
537.311.33:546.482.21:535.215 1700

Determination of Trap Distribution from Interrupted-Illumination Measurements on Photoconductive Cadmium Sulphide Single Crystals-E. A. Niekisch. (Z. Naturf., vol. 9a, pp. 700-701; July/August, 1954.)
537.311.33:546.682-31

1701
Investigations of Electrical and Photoelectric Conductivity of Thin Films of Indium Oxide G. Rupprecht. (Z. Phys., vol. 139, pp. 504-517; December 20, 1954.) The conductivity of thin films of the $n$-type semiconductor $\operatorname{In}_{2} \mathrm{O}_{3}$ is markedly dependent on the surrounding atmosphere. Specimens prepared by evaporating In on to a quartz plate and heating in air at 700 degrees $-1,000$ degrees $C$. had thicknesses between 50 and $250 \mathrm{~m} \mu$ and conductivities between 10 and $10^{-5}(\Omega, \mathrm{~cm})^{-\mathrm{t}}$. Measurements are reported of the temperature variation of conductivity and of the effect of an oxygen atmosphere. Above 500 degrees $C$, a balance is reached bet ween the concentration of impurity centers and the external oxygen concentration. The photoelectric conductivity exhibits an irreversible increase in vacuum.
537.311.33:546 682.231

1702
Electrical and Optical Properties of Indium Selenide-R. W. Damon and R. W. Redington, (Phys. Rev., vol. 96, pp. 1498-1500; December 15, 1954.) Measurements mainly on single crystals are reported. The optical absorption edge was not sufficiently well defined for the energy gap to be estimated unambiguously. The photoconductive response was mostly in the visible region, the sensitivity being comparable with that of grey Se . Attempts to determine the carrier type gave conflicting results; the material may not be a single-carriertype semiconductor, at least within the surface region.

### 537.311.33:546.86:539.234 1703

Fermi Level in Amorphous Antimony Films -l:. Taft and L. Apker. (Phys. Rev., vol. 96, pp. 1496-1497; December 15, 1954.) Photoelectric experiments are reported which confirm that the amorphous form of Sb common in thin evaporated films is a semiconductor with Fermi level about 0.1 ev above the occupied band.
538.221

1704 Investigations of Irreversible Magnetization and After Effect-J. Kranz. (Z. Phys., vol. 139, pp. 619-637; December 20, 1954.) An experimental arrangenient is described for measuring the Barkhausen jumps, produced on reversing the field applied to a ferromagnetic specimen, by amplitude-analyzing and counting the pulses induced in a solenoid. Results for various materials are presented and discussed.

### 538.221

1705
Large Magnetic Kerr Rotation in BiMn Alloy-B. W. Roberts and C. P. Bean. (Phys. Rev., vol. 96, pp. 1494-1496; December 15, 1954.) Brief illustrated note on observations of ferromagnetic domain patterns in large grains of BiMn.
538.221

Kinetics of Magnetization in Some Square Loop Magnetic Tapes-C. P. Bean and D. S. Rodbell. (Jour. Appl. Phys., vol. 26, pp. 124125: January, 1955.) Curves showing flux reversal characteristics for permalloy tapes are discussed in relation to the domain-wall processes.
538.221

1707
The Effect of Particle Shape Variations on the Coercivity of Iron-Oxide Powders-W. P. Osmond. (Proc. Phys. Soc., vol. 67, pp. 875882; December 1, 1954.) Calculations show that a Gaussian distribution of particle shape factors about an observed mean value can satisfactorily explain the clifference between measured values of coercivity of dispersed magnetic powders and the theoretical values for assemblages of identical particles.
538.221: 621.318.134:537.311.33

1708
The Nature of the Insulating Layers in Ferromagnetic Semiconductors-R. Parker. (Physica, vol. 20, pp. 1314-1315; December, 1954.) Recent experimental results [ 3604 of 1954 (Volger)] can be explained on the follow ing assumptions:-(a) that the appearance of spontaneous magnetization is the cause of the deviation from the normal relation between resistivity and temperature, and (b) that the insulating layer may be identified with the region in the material that is not spontaneously magnetized.

### 538.221:621.318.134

1709
Saturation Magnetization and Crystal Chemistry of Ferrimagnetic Oxides-E. W. Gorter. (Philips Res. Rep., vol. 9, pp. 295-320, 321-365 and 403-443; August-December, 1954.) A thesis in which measurements of saturation magnetization, $\sigma$, against temperature, 7 , are reported for various mixed crystal oxides with spinel structure. Results are in agreement with Neel's theory; some of the anomalous $\sigma / T$ curves predicted have been found. Single ferrites investigated of the type $\mathrm{Me}^{\mathrm{II}} \mathrm{Fe}_{2}{ }^{I I I} \mathrm{O}_{4}$ belong to a group with complete parallelism of the ionic moments inside each sublattice; mixed crystals of the type $\mathrm{Me}_{1-a} \mathrm{Zn}_{a} \mathrm{Fe}_{2} \mathrm{O}_{4}$ with $a>0.4$ belong to a group with angles between the ionic moments inside one of the sublattices; for $\mathrm{Ca}_{0.36} \mathrm{Zn}_{0.65} \mathrm{Fe}_{2} \mathrm{O}_{4}$ the magnetic moment is higher than that of any MgZn ferrite; this is discussed with reference to Anderson's theory. Other materials investigated are ferrimagnetic spinels containing Ti and Al , and ferrimagnetic oxides containing Cr .
538.221:621.318.134

1710
Low-Temperature Acoustic Relaxation in Ni-Fe Ferrites-M. E. Fine and N. T. Kenney. (Phys. Rev., vol. 96, pp. 1487-1488; December $15,1954$.$) "An acoustic relaxation effect occurs$ near 40 degrees K in $\mathrm{Ni}_{0.75} \mathrm{Fe}_{2.25} \mathrm{O}_{4}$ and is attributed to a stress-induced change in distribution of $\mathrm{Fe}^{++}$and $\mathrm{Fe}^{+++}$similar to that occurring in magnetite. The process involves electron diffusion. The activation energy is
between 0.026 and 0.055 eV per electron
jump."

### 538.221:621.318.134

1711
Magnetic and Crystalline Behavior of Certain Oxide Systems with Spinel and Perovskite Structures-L. R. Maxwell and S. J. Pickart. (Phys. Rev., vol. 96, pp. 1501-1505; December 15, 1954.) Experiments are reported in which nonmagnetic trivalent io 1 s were substituted for $\mathrm{Fe}^{3+}$ in Ni ferrites.

### 538.221:621.318.134

1712
Conference on Ferrites, Leningrad, 1st-5th February 1954-(Bull. Acad. Sci. URSS, sér. phys., vol. 18, pp. 307-416 and 419-520; May/June and July/August, 1954.) The text is presented of more thar 20 papers covering theoretical and experimental investigation subjects discussed inclu•le Faraday effect at centimeter wavelengths and temperature dependence of electrical properties of ferrites.

### 538.221:681.142

Magnetic Materials for 1713 Components: Part 1-A T versal in Polycrystalline Ferromagnetics N. Menyuk and J. B. Goodenough. (Jour. Appl. Phys., vol. 26, pp. 3-18; January, 1955.) Output-voltage waveforms of computer storage elements are consistent with the assumption that the flux reversal is attributable to the creation and growth of 180 degree Bloch walls. A switching coefficient is defined having one component dependent ou eddy current and another depending on spin relaxation; for ferrites and thin metal tapes the first of these components is much smaller than the latter. Consideration of various parameters involved indicates that it is better $t$, produce hysteresisloop squareness by grain orientation or magnetic anneal than by apolication of external stress or variation of cher acal composition.

### 538.23

A Relation between Hysteresis Coefficient and Permeability: Part 2-Further Experimental Results-M. Kornetzki. (Z. Angew. Phys., vol. 6, pp. 547-550; Decesaber, 1954.) The investigation on ferrites reported earlier ( 756 of 1953) was extended to cover permanent magnets and various $\mathrm{Fe}, \mathrm{Fe}-\mathrm{Si}, \mathrm{Fe}-\mathrm{Si}-\mathrm{Al}$ and $\mathrm{Ni}-$ Fe alloys with initial permeability, $\mu$, between 1.25 and 120,000 and hysteresis coefficient, $h$, between 0.65 and $9,000,000 \mathrm{~cm} / \mathrm{ka}$. A doublelogarithmic plot of $h /(\mu-1)$ against $(\mu-1)$ shows that the points for most of the materials lie between a pair of parallel lines of slope 1.15 and separation, measured on the $h /(\mu-1)$ scale, corresponding to a ratio of 40 . All groups, except the Fe and $\mathrm{Fe}-\mathrm{Si}$ alloys, include lowhysteresis materials for which $h /(\mu-1) \approx 3$ $\mathrm{cm} / \mathrm{ka}$. Values of $\mu$ and $h$ are tabulated for over 40 materials, and graphs are plotted relating $(\mu-1)$ to $h, h /(\mu-1)$, and $h^{\prime}(\mu-1)^{2}$.

### 621.315.613.1:537.529

1715
Phenomena preceding Dielectric Breakdown in Mica-B. Fallou (Rev. Gén. Elecl., vol. 63, pp. 643-653; Noveıaber, 1954.) Report of an experimental investigation. Oscillograms are reproduced and discu sed in relation to charge conditions at the surfaces of separation in the mica.

### 621.315.616:537.226

1716
Dielectric Breakdown of Thermosetting Laminates-N. A. Skow. (Mlod. Plash., vol. 32, pp. 152, 240; December, 19.4.) A report is presented on short-time and indurance tests on laminates bonded with phenolic resin. The grades tested included thriee based on paper and one each on asbestos, cotton fabric, glass and nylon. The variation of the dielectric strength with temperature, clirection of applied field, thickness of laminae and conditioning of the specimens is tabulated and some results are also presented graphically. The most suit-
able grades for use under various conditions (e.g. high humidity) are indicated.
621.315.614.4 1717
Forest Products Research Special Report No. 8. The Dielectric Properties of Wood [Book Review]-R. F. S. Hearmon and J. N. Burcham. Publishers: II. M. Stationery Office, London, 1954, $19 \mathrm{pp}$. 1s. 6d. (Elec. Times, vol. 126, p. 848; December 9, 1954.) An investigation of the influence of grain direction, density and moisture content on the permittivity and loss tangent of 12 species of wood over the frequency range $2 \mathrm{kc}-60 \mathrm{mc}$ is reported.

## MATHEMATICS

517.5

1718
The Approximation to a Characteristic Function by its Fourier Series-D. Dugué. [Compl. Rend. Acad. Sci. (Paris), vol. 240, pp. 151-152; January 10, 1955.]

### 517.512.2

Moving-Strip Fourier Analyzer-H 1719 Grenville-Wells ( $R$ ew. Sci Analyzer-H. J 1156-1161; December, 1954.) A modified form of the device described by Robertson (Phil. Mag., vol. 21, pp. 176-187; January, 1936.) Two-dimensional and three-dimensional summations can be performed.

### 517.564.3

1720
The Asymptotic Expansion of Bessel Functions of Large Order-F. W. J. Olver. (Phil. Trans. A., vol. 247, pp. 328-363; December 28, 1954.)
517.9

1721
The Asymptotic Solution of Linear Differential Equations of the Second Order for Large Values of a Parameter-F. W. J. Olver. (Phil. Trans. A, vol. 247, pp. 307-327; December 28, 1954.)

### 519.272 .119

Qualitative Evaluation of Correlation 1722 efficients from Scatter Diagrams-T. M. Burford. (Jour, Appl. Phys., vol. 26, pp. 56-57; January, 1955.) Analysis shows that the result previously obtained by Sugar (2457 of 1954) for the case of a Gaussian distribution can be applied to any distribution.

## MEASUREMENTS AND TEST GEAR <br> $621.3 .018 .41(083.74)+621.396 .91$

 Essen. (Proc. IEE, Part B, vol. 102, pp. 173178; March, 1955.) Discussion on 3289 of 1954.
### 621.3.018.41(083.74)

Standard Frequency Transmission 1724 ment at Rugby Radio Stansion Equip(Proc. IEE, Part B March 1955. Disu. 102, pp. 166-173 March, 1955. Discussion, pp. 173-178.) "The transmissions are controlled by one of three highly stable oscillator-clock chains, which are checked daily in terms of a frequency standard, and are continuously intercompared in order that instability may be quickly detected. Automatic shut-down features are incorporated to reduce the risk of broadcasting incorrect frequencies under fault conditions. The vital parts of the equipment are protected against mains failure. The transmitted frequencies are normally kept within 1 part in $10^{3}$ of the currentlyassessed nominal frequency, based on predicted clock performance, so as to keep within a tolerance of $\pm 2$ parts in $10^{8}$ in terms of finallycorrected time determinations. Day-to-day frequency variations are usually less than $\pm 2$ parts in $10^{\circ}$."
621.3.081.41(083.74): 621.317.761 1725 The Standard Frequency Monitor at the National Physical Laboratory-J. McA. Steele. (Proc. IEE, Part B, vol. 102, pp. 155-165; March, 1955. Discussion, pp. 173-178.) Description of equipment for the automatic
measurement and recording of the MSF standard frequency transmissions on $60 \mathrm{kc}, 2.5 \mathrm{mc}$, 5 mc and 10 mc and of the Droitwich 200kc transmission. The measurements on Droitwich and on MSF 60 kc are made by an extension of the established method for intercomparison of the N.I.L. frequency standards; the standard deviations for measurements lasting a few seconds are 3-4 parts in $10^{\circ}$ and $2-3$ parts in $10^{\circ}$ respectively. " Measurements of the phase of the $1 \mathrm{c} / \mathrm{s}$ pulse modulation on MSF $60 \mathrm{kc} / \mathrm{s}$ can be made with very high precision. In daily comparisons, the scatter of a group of 60 readings does not usually exceed $30 \mu \mathrm{~s}$. By a process involving frequency changing by continuous phase shifting, the received frequencies of MSF 2.5 and $5 \mathrm{Mc} / \mathrm{s}$ are recorded on a frequency meter of range $\pm 1 \mathrm{c} / \mathrm{s}$, the discrimination being 2 parts and 1 part in $10^{8}$ at these frequencies, respectively. The records obtained show well-defined diurnal variations on $2.5 \mathrm{Mc} / \mathrm{s}$, particularly at sunrise, where the deviations may amount to several parts in $10^{7}$. The received frequency of $5 \mathrm{Mc} / \mathrm{s}$ is subject to continuous variations over a range of about 1 part in $10^{7}$ during daylig't hours; in darkness larger changes are recorded."

### 621.3.018.41(083.74): 621.373.4.029.5

1726
Quartz Resonator Servo: a New Frequency Standard-N. Lea. (Marconi Rev., vol. 17, pp. 65-73; 3rd Quarter, 1954.) The response to an applied oscillator signal of a briclge detector circuit including a quartz crystal resonator is used to control a servodrive-operated variable capacitor in the oscillator tuned circuit. In an experimental instrument, using a $5-\mathrm{mc}$ crystal, frequency variations were reduced to less than 1 part in $10^{10}$ for variations of +20 per cent in the applied hv, a 1 v change from 5 to 6 v , or an oscillator capacitance change which would, if uncorrected, cause a change of 200 parts in $10^{10}$. A rate of frequency correction of 1 part in $10^{9}$ per second was easily achieved. The circuit diagram given is briefly commented on. Experimental results are presented graphically.
621.3.018.41(083.74):621.396.82

1727
Effect of Interference by Other StandardFrequency Transmissions upon the Accuracy of Frequency Calibration by a Standard-Frequency Transmission-K. Matsumoto, T. Nagatake and Y. Suguri. (Jour. Radio Res. Labs, Japan, vol. 1, pp. 41-48; March, 1954.) Experimental results indicate that interference with the JJ Y standard-frequency transmissions by station WWVH has a negligible effect on the accuracy of frequency calibrations when the reference beat method is used.

### 621.3.018.41(083.74): 621.396.91 <br> 1728

An Experiment on the Types of Time Signals superposed on the Standard-Frequency Transmission-K. Matsumoto and T. Nagatake. (Jour. Radio Res I.abs, Japan, vol. 1, pp. 49-56; March, 1954.) Coding of a JJY standard-frequency transmission by a $20-\mathrm{ms}$ interruption of the $1-\mathrm{kc}$ tone each second and a 200 -ms interruption each minute, was compared with coding by the CCIR method of substituting a $1.4-\mathrm{kc}$ tone for 5 ms each second. Using simple receivers at $1,000 \mathrm{~km}$ from the transmitter, the former method was found to be the more suitable one.

[^63]621.317.3:621.315.61

1730
Measurements of Electrical Polarization in Thin Dielectric Materials-R. W. Tyler, J. H. W'ebb and W. C. York. (Jour. Appl. Phys., vol. 26, pp. 61-68; January, 1955.) A method suitable for measuring electrical effects such as are produced in film moving over a roller system consists in arranging the dielectric material in contact with a grounded metal backing plate and placing a field meter at a short distance in front of the dielectric. Tests made with a short strip of cellulose acetate film without the emulsion coating are described.

### 621.317.321:538.632

1731
Apparatus for Measurement of Hall Effect and Magnetic Change of Resistance with Alternating Current-K. A. Muller and J. Wieland. (Ilelv. Phys. Acla, vol. 27, pp. 690-696; 1)ecember, 31, 1954. In English.) The null method described is operated at 73 cps and is suitable for detecting voltage changes down to $2 \times 10^{-8} \mathrm{v}$ in specimens of resistance between $10^{-8}$ and $10^{-2} 22$. The pd across the specimen or the Hall emf is determined by compensating it by an equal and opposite voltage derived from a fixed resistor in series with the specimen. The null detector comprises an amplifier, filter, phase discriminator and and an indicator instrument which is either a cro or an aperiodically damped galvanometer.

### 621.317.331:621.385.2

1732
A Method for recording Logarithmic Variations of Resistance-H. A. Vodden. (Jour. Sci. Instr., vol. 31, pp. 475-476; December, 1954.) A simple circuit of an ohmeter recording resistance logarithmically is based on the fact that the logarithin of the anode current of a diode is proportional to the anode voltage over a range of negative voltage values. The useful range of the instrument described is between about $2 \times 10^{2}$ and $10^{9} \Omega$.
621.317.336:621.317.755

1733
Visual Impedance-Matching EquipmentR. Dalziel and A. Challands. (Wireless Eng., vol. 32, pp. 99-107; A pril, 1955.) A cro method is described for indicating the degree of match between a load and a cable, as e.g. in antenna feeding. The test oscillator is mechanically swept over the whole frequency range of $80-$ 250 mc . The oscilloscone face is calibrated in terms of swr. Impedance measurements made with the equipment have yielded results in good agreement with those obtained by other methods. See also 2852 of 1948 (Libby).

### 621.317.361:621.385.029.65

1734
Cold Measurements of 8 mm Magnetron Frequency and Pulling Figure-Barrington. (See 1839.)
621.317.382:538.632:537.311.33 1735
The Application of the Hall Effect in a Semiconductor to the Measurement of Power in an Electromagnetic Field-M. E. M. Barlow. (Proc. IEE, Part B, vol. 102, pp. 179-185; March, 1955. Discussion, pp. 199-203.) Analysis shows that the mean value in time of the Hall emf is a direct measure of the power traversing the semiconductor in steady or varying fields. Residual rectifier effects are eliminated by operating with a strong magnetic field. Various types of wattmeter embodying the princijle are described. Experiments with an $n$-type Ge crystal mounted between the inner and outer conductors of a coaxial line indicate that the Hall effect is approximately the same at 50 cps and at 300 mc , so that instruments for use at high frequencies can be calibrated at low frequency.
621.317 .443

1736
An Improved Precision PermeameterC. D. Mee and R. Street. (Proc. IE I., Part II, vol. 101, pp. 639-642; December, 1954.) A modified form of the de permeameter de-
scribed by Armour et al. (3499 of 1952) uses a saturable-inductor type of field-measuring device to give automatic and continuous indication of the required compensating-coil current at all points on the $B / I I$ curve. Values of $I I$ from $10^{-3}$ oersted upward can be measured.
621.317.715:621.383.2

1737
Photodianode and Galvanometer Feed-back-L. Deloffre, É. Pierre and J. Roig. [Compl. Rend. Acad. Sci. (I'aris), vol. 240, 1pl). 59-61; January 3, 1955.] Analysis is presented relevant to galvanometer measurements using the twin-anode fhotocell device previonsly described ( 2543 of 1954 and back reference) in a feedback arrangement. The galvanometer sensitivity can be multiplied by a factor as great as 10 in this way:
621.317.72: 621.3.018.3

1738
Two-Frequency Waveforms: Effects on Rectifier Instruments-J. E. Parton and W. D. Sutherland. (Trans. Soc. Instrum. Technol., vol. 6, pj. 147-161; December, 1954.) Asymmetrical waveforms resulting from the presence of even harmonics are included in this study. For a waveform with two frequency components, a voltmeter with full-wave metal rectifior is found to read within 2 per cent of the value calculated on the assumption of perfect rectification. The mean value differs from that given by a rms instrument in being dependent on the relative phase of the two components.
621.317.733:621.317.4

1739
Mutual Inductance Bridge and Cryostat for Low-Temperature Magnetic Measurements -R. A. Lirickson, I. D. Roberts and J. W. T. Dabls. (Kev. Sci. Insir., vol. 25, pp. 1178 1182; December, 1954.)
621.317.755: 621.314.7

1740
An Alpha Plotter for Point-Contact Transis-tors-T. P. Sylvan. [Elec. Engng., ( Neve York), vol. 73, pD. 1094-1098; December, 1954.] Description, with detailed parts list, of a cro test set.

### 621.317.78.029.5/.64

1741
Broadband R.F. Power Meters-1. Strauss. Radio-Electronic Engng. vol. 23, pp. 10-11, 36; December, 1954.) Equipment for measuring average rf power from $5 \mu \mathrm{w}$ to 5 w in the frequency range 20 mc to 10 kmc is briefly described; three frequency sub-ranges are covered by separate instruments. The energy dissipated in a bolometer element in a Wheatstone bridge is kept constant by (a) varying the dc current through it, and (b) attenuating the rf energy.

### 621.317.78.029.65

1742
A Calorimeter for Power Measurements at Millimeter Wavelengths-W. M. Sharpless. (Trans. IRE, vol. MTT-2, pp. 45-47; September, 1954.) Description of an instrument suitable for measuring power of the order of 1 mw , in which equal temperature rises are produced in two waveguide-section power absorbers, one of which is heated by de and the other by the rf power.

### 621.317 .784

1743
Audio-Frequency Power Measurements by Dynamometer Wattmeters-A. H. M. Arnold. (Proc. ILEE, l’art B, vol. 102, pp. 192-199; March, 1955. Discussion, 1p. 199-203. The screening necessary to obtain accuracy in the upper af range comparable to that at power frequencies is discussed. An account is given of methods used at the NPL to calibrate wattmeters. The useful uprer frequency limit is taken as the frequency at which a significant deflection is obtained with voltage only or current only applied to the terminals.
621.317.784:538.632:537.311.33

1744
The Design of Semiconductor Wattmeters for Power-Frequency and Audio-Frequency

Applications-II. E. M. Barlow. (Proc. IEE, Part B, vol. 102, pp. 186-191; March, 1955. Discussion, pp. 199-203.) Design details and performance characteristics are given for two wattmeters based on the Hall effect in semiconductors ( 1735 above), for use at frequencies up to 150 cps and 20 kc respectively. The power-frequency instrument incorporates an iron-cored magnetizing coil; the af instrument uses air-cored coils and includes screening arrangements. These instruments offer advantages over other types for measurements of high power.

### 621.37.029.6.049.001.4

1745
A Surface-Terture Comparator for Microwave Structures-A. F. Harvey. (Proc. IEE, Part B, vol. 102, pp. 219-222; March, 1955.) The dependence of the attenuation coefficient of microwave components on the relation between surface roughness and skin depth is discussed and a description is given of a simple comparator scale covering the various classes of finish in normal use. Measurements are stated in terms of the center-line-average figure, in microinches, obtained on traversing a small stylus over a sample of the surface. Comparisons are made by sight and by touch.
621.373.42.001.4:621.317.361

1746
Testing Precision Oscillators - M. P. Johnson. (W'ireless World, vol. 61, pp. 179-182; April, 1955.) The frequency stability of $124-$ kc oscillators used as masters for carrierfrequency telephone systems is determined by comparison with the $2.5-\mathrm{mc}$ standard-frequency transmissions from Rugby, the measurement being made at a point 10 miles away Counts of the difference frequency are made over regularly recurring sampling periods, and an output current proportional to the count is obtained. A linear recording meter is used: full-scale detlection is produced by a count of 58 , corresponding to a frequency difference of 2 parts in $10^{7}$ for a sampling period of 116 seconds. The equipment is described and specimen records are shown.

### 621.373.52:621.314.7 <br> 1747

Transistorized F.M. Signal GeneratorJ. J. Hupert and T. Szubski. (lilectronics, vol. 28, pp. 133-135; Fehruary, 1955.) An instrument covering the frequency range $20-100 \mathrm{mc}$ has been designed giving an output of 10 mw across 10s?. The vhif section conurises a IrM oscillator onerating at a half or athird of the output frequency, followed by a harmonicselector stage. The frequency modulation is effected by a transistor acting as variable reactance. The relative merits of point-contact and junction-tetrode transistors for this circuit are discussed. The saving in bulk as compared with equipment using thermionic tubes is to some extent offset by the need for a constantdemperature enclosure for the oscillator.
621.375.2.024.083:681.142

1748
Gain Measurements on Computing Amplifiers $-A$. B. Johnson. (Electronic lingng., vol. 27, pp. 127-129; March, 1955.) Special techniques are required for measurements on amplifiers for de amalog computers, which usually have very high gatin. Mathols are divided into two broad classes, (a) direct, in which the drift output is reduced without affecting the test signal, the amplifier operating effectively without feedback for the latter, and (b) indirect in which the drift output is small becaluse of negative feedback, and the gain is deduced from some other property

### 621.397.5:535.623].001.4

1749
Phase Measurement for Color TV and F.M —K. Schlesinger. (Electronics, vol. 28, It), 142-146; February, 1955.) The principles of operation with the vectorscope are described.

### 621.317.3.029.6 <br> 1750

Handbook of Micsowave Measurements, Vols. 1 and 2 [Book F eview]-M. Wind and H. Rapaport, Eds. I ublishers: Polytechnic Institute of Brooklyn, Vew York, 616 pp . and 320 pp. (Wireless Eng. vol. 32, p. 116; April, 1955.) A useful compesidium not only for the technician but also, as a supplement to more theoretical treatments, for the engineer or physicist. The text is all in Volume 1, the diagrans are in Volume 2

## OTHER APPLICA IONS OF RADIO AND ELECTRONICS

## 621.3:61

 1751Electricity in Medicine-S. N. Pocock. (Proc. IEE, l'art 11, vol. 101, pp. 629-o38 December, 1954.) A su vey of applications of electronic techniques in diagnosis and therapy.

### 621.317.39:531.71

1752
A Direct-Reading Instrument for the Measurement of Small Displacements-W. 1). Corner and G. H. Hunt (Jour. Sci. Instr., vol 31, pı, 445-447; Deceuber, 1954.) Displacements down to $2 \times 10^{-4} \mathrm{~cm}$ were measured with an accuracy of within $10^{-7} \mathrm{~cm}$ by means of a bridge circuit using a differential capacitor, with a $100-\mathrm{v}, 10-\mathrm{kc}$ suppis. The unbalance was measured by a direct-indicating tube voltmeter. The apparatus vias designed for mag netostriction measuremeats.

### 621.384.612

1753
Generating R.F. Energy for $6-\mathrm{kMeV}$ Beva-tron-C. N. Winningstad. (Electronics, vol. 28, pp. 164-169; leebruary, 1955.)

### 621.385 .833 <br> 1754

Extension of the Electron-Optical Theory of the Deflecting Electrostatic System to the Case of Relativistic Particles-A. M. Strashkevich. (Zh. Tekh. Fiz., vol. 24, pp. 2264-2270; December, 1954.)
621.385 .833

1755
Aberrations of Relativistic Electron Beams -A. M. Strashkevich and N. G. Gluzman. (Zh. Tekh. Fiz., vol. 24, p'). 2271-2284; December, 1954.) A mathemat cal cliscussion is pre sented of the operation of a system with a curved axis. Equations (9, and (10) are derived for a wide beam in an : rbitrary electrostatic field for the relativistic case. Equations are also derived for the particular cases of a wide beam in an axially symmetric field (28) and in a plane field (30) as well is in the fields of a cylindrical lens (33) an a cylindrical condenser (35). The aberrations of axially symmetric lenses are calculat :d for the relativistic case (40)-(45), and also of cylindrical lenses (50)-(53).

### 621.385 .833

1756
Simple Presentation of the General Theory of Systems of Revolution in Electron Optics (covering Relativity and Aberrations) - $\hat{\mathbf{E}}$ Durand. (Rev. d'()ptique, sol. 33, pp. 617-629) I ecentber, 19.54.) Calculat ions are simplified by introducing complex con bbinations into the Lagrangian.
621.387.424:537.52 1757

A Cloud-Chamber Study of some Aspects of the Geiger Discharge-l'. J. Campion. (Proc. Phys. Soc., vol. 67, pp. 1095-1102; December 1, 1954.)

### 621.396:623.451.8

1758
Launching Control for Guided MissilesJ. 13. Schrock. (Electronic: vol. 28, pp. 122127; February, 1955.) A description of the cir cuits which control the fring of the missile actuate the guiding and telemetering equipment, and alert rocket-range control units.

## PROPAGATION OF WAVES

Coupled Wave Equations for Inhomogeneous Anisotropic Media-K. Suchy. (Z. Naturf., vol. 9a, pp. 630-636; July/August, 1954.) "A special system of coordinates has been introduced for the calculation of electromagnetic wave propagation in an inhomogeneous, anisotropic medium. One of the coordinate axes is parallel to the wave normal, the two others (perpendicular to it) are defined by the relation between the $\mathbf{E}$ and $\widetilde{\mathbf{D}}$ vector. In the coupled wave equations it is shown that the coupling terms can be neglected under certain conditions."
538.566

1760
On the Possibility of Electromagnetic Surface Waves-P. S. Epstein. [Proc. Nat. Acad. Sci. (Washington), vol. 40, pp. 1158-1165 December, 1954.] An independent surface wave is defined as comprising two inhomogeneous waves which are independent of each other and run along the surface dividing two media, one wave in cither medism. A discussion of the conditions for the existence of such a wave indicates that it could exist only at the boundary of two nonconducting media, one with a positive dielectric constant, the other with a negative one. Sommerfeld's solution for the field of an electric dipole at the surface of a plane earth is briefly coramented on

### 538.566537 .56

1761
General Expression for the Absorption of Electromagnetic Waves in Lorentz-Type Plasmas (Ionosphere)-M. Lozzi, R. Jancel and T. Kahan. [Compl. Rend. Acad. Sci. (Paris), vol. 240, pp. 162-164; January 10, 1955.] A general expression is derived from the theory presented previously [2364 of 1954 (Jancel and Kahan)].
538.566.2 1762

Reflection and Refraction of Electromagnetic Waves in a Stratified Medium-K. Försterling. (Hochfrequenzlech. u. Eleklroakust., vol. 63, pp. 112-116; April, 1954.) Exact analysis is presented and the effect of certain approximations is discussed. Splitting due to the presence of magnetic fields is considered.
621.396 .11

1763
Theory of Radio Transmission by Tropospheric Scattering using Very Narrow BeamsH. G. Booker and J. T. de Bettencourt. ('roc. IRE, vol. 43, pp. 281-290; March, 1955.) When propagation is effected by scattering from turbulences in the troposphere [1757 of 1950 (Booker and Gordon)]. the energy arriving at the receiver may be expected to extend over. a substantial angle; it should be possible to demonstrate the effect by using narrow beans, of width less than about 1.5 degrees. Calculations are made for a communication path of length 300 km between raraboloidal antennas of diameter $100 \lambda$; a beare width of 0.73 degree is assumed. A study is made of the elfects to be expected on swinging the two beams in synchronism so that their axes always intersect; a hateral swing of 1 dregree off the greatcircle path would reduce the received power by about 7 db if propagation is controlled by scattering, whereas it wouid reduce the receiver power by about 40 db if propagation is controlled by refraction. The distortion of pulses to be expected as a result of beam swinging is evaluated. Choice of communication bandwidths is discussed.

### 621.396.11:551.510.535

1764
Measurement of Attenuation in the Iono-sphere-A. Ochs. (Arch. elekl. Überiraging. vol. 8, pD. 535-544; December, 1954.) i fixedfrequency method based or a continuous photographic record of echo amplitude is discussed,
and measurements made during the period October, 1952-January, 1953 are reported. Difficulties are introduced ly ground reflection, ground irregularity, fluctuations of transmitted power, the effect of the extraordinary component, and nonuniformity of the ionosphere. The last-mentioned factor is especially important, and its effects are illustrated by field-strength records of a signal (a) once reflected and (b) twice reflected from the $I$. liyer. The absence of correlation between these records can be explained by assuming that the reflected surface is inclined or curved, and/or that within the usual ionosphere strata there are local regions, or clouds, with higher concentrations of electrons. To eliminate errors due to the focusing effect of the curved surface, measurements must be averaged over a suitable period, must be made at suitably spaced points, and must be made simultaneously at different frequencies. If only a single frequency is used, this should be as low as possible with the available power. Details are given of the equipment used. Carrier frequencies of 2 mc and, later, 1.6 mc were used, with pulse duration $100 \mu \mathrm{~s}$, pulse power $15-20 \mathrm{kw}$, and pulse repetition rate 50 or 1 per second. The results indicate considerable interdiurnal differences of the diurnal variation of attenuation.

### 621.396.11:551.510.535

1765
Ionospheric Absorption Measurements at Prince Rupert-K. Davies and E. L. Hagg. (Jour. Almos. Terr. Phy's., vol. 6, pp. 18-32; January, 1955.) Report and discussion of measurements made nea: the northern auroral zone between April, 1949 and March, 1950. Monthly median noon values of total absorption, $-\log \rho$, do not fit the inverse square law $-\log \rho \propto\left(f+f_{L}\right)^{-2}$; scasonal variation indicates little dependence on solar zenith angle $x$. Diurnal variation is very approximately represented by the relation $-\log \rho \propto(\cos \chi)^{0.5}$, maximum absorption generally occurring about 20 minutes after local noon. There is a pronounced correlation between night-time absorption at 2 me ansl the 3 -hour-range $K$ index for $K>4$. High night-time absorption is often associated with intense sporadic $E$.
621.396.11:551.510.535 1766 Influence of the Inclination of the Earth's Magnetic Field on the Absorption of Radio Waves in the D Layer- $\mathrm{I}^{\prime}$. Lejay and $D$. Lepechinsky. [Compt. Fend. Acad. Sci. (Paris), vol. 240, pp. 136-13.3; January 10, 1955.] Analysis is presented and its application is illustrated by evaluating the absorption index for three particular directions of propagation. For the extraordinary ray the absorption increases considerably with increase of the angle between the direction of the earth's magnetic field and the direction of propagation; for the ordinary ray the varistion of absorption is in the opfosite sense. The results indicate that caution is necessary in applying absorption values obtained from: vertical soundings to conditions along actual radio communication paths; the cosine law is not directly applicable.
621.396.11:551.510.535

1767
Reflection Conditions for Vertical Propagation in the Ionosphere in the Presence of Collisions and of the Eartn's Magnetic Field. Case of the E Layer-D. I.epechinsky and J. Durand. [Compl. Rend. Acad. Sci. (Paris), vol. 240, pp. 333-3.36; January 17, 1955.] The type of propagation, i.e. quisi-transverse or quasilongitudinal, occurring in a given region of the ionosphere is determined by $\theta$, the inclination of the earth's magnetic field, and the collision number. A graph shows the height $h$ at which the transition occurs, as a function of $\theta ; h$ coincides with the $E$;ayer ( 100 km ) for a value of $\theta$ about 6.5 degrees, i.e. in France at about the latitude of Paris; propagation there is quasitransverse or quasi-'ongitudinal according as the collision number increases or decreases.

Consideration of the relation between refractive index and electron concentration for the two types of propagation indicates that neglect of collisions introduces error in the calculated values of $f_{0} E$ for values of $\theta>6.5$ degrees. These considerations are shown to explatin the apparent flattening of the daily variation curve of $f_{0} E$ around midday.
621.396.11.029.55

1768
The Prediction of Short-Wave Propagation -IB, IBeckmann, (Tech. Hausmill. NordwDisch. Kdfunks, vol. 6, nos, 9/10 and 11/12, pp. 211-219 and 247-259; 1954.) Concepts fundamental to ionospleric and radio weather ( 3605 of 1954) and the relations between them are discussed. The relevance of solar, geomagnetic and ionospheric factors in forecasting is indicated. An examination is made of the accuracy of forecasts and of the causes of error, e.g. scattering. Statistics are presented showing differences in reception of WVWV transmissions at Norddeich and Munich.

### 621.396.11.029.6

1769
An Experimental Study of the Propagation of 10 cm Radio Waves over a Short Non-optical Sea Path-E. F. Stack-Forsyth. (Proc. IEE, Part B, vol. 102, pp. 2.31-236; March, 19.35.) Tests over a path 1.14 times the opticalhorizon distance were made off the coast of Natal during the winter months Ajril-August, 1953. Vertical polarization was used. The effect on signal strength of variations in the structure of the refractive-index profile of the atmosphere in the first few hundred fert above sea level was studied. The results indicate that a duct about 120 feet high was present for a considerable part of the winter. The signal strength at a height of 47 feet in the duct was $6-10 \mathrm{db}$ above the free-space value. The absolute value of the signal strength and its variation with duct height are in moderately good agreement with values given by the mode theory, using only the first mode and assuming either a square-law, fifth-root or bilinear profile.
621.396.11.029.6:551.510.535

1770
Study of Long-Distance Propagation of V.H.F. Waves by Sporadic-E IonizationT. Kono, Y. Uesugi, M. Hirai and G. Abe. (Jour. Radio Res. Labs, Japan, vol. 1, pp. 1-10; March, 1954.) A report is presented of verticalincitlence measurements and propagation tests carried out in Japan between June and August, 1952 , over distances of $500-1,100 \mathrm{~km}$ at frequencies of about 31,43 , and 65 mc . Results were analyzed statistically; empirical expressions are given for the relation between $f E_{s} / f \cos i$ and $F$, where $F$ is the field strength relative to the free-space field strength, $f$ is the transmission frequency, and $i$ is the angle of incidence. The probability of the calculated value of $F$ exceeding its actual value can be estimated. Results are presented graphically.

### 621.396.11.029.6:[551.524+551.57

1771
What Role does the Nocturnal Cooling play in the Ultra-short-Wave Propagation?-K. Ilirao. (Jour. Radio Res. Labs, Japan, vol. 1, p12. 27-39; March, 1954.) Variations of the temperature and humidity in the lower atmosphere at the transminter are compared with the field-strength variations at the receiving station. The frequency used was 150 mc . Results are tabulated and also presented graphically: Nocturnal cooling has both a direct effect and an indirect effect through the resulting changes of humidity.
621.396.812.3:551.510.535 1772

The Autocorrelogram of Randomly Fading Waves-R. B. Banerji. (Jour. Almos. Terr. Phys., vol. 6, pp. 50-56; January, 1955.) The received power spectrum for the case of a completely rough ionosphere having superposed steady and randon motion is deduced assuming transmitting and receiving antennas omni-
directional. Autocorrelograms of fading patterns corresponding to pure drift and pure turbulence are compared. An observed fading pattern need not contain more than 250 independent points for these two extreme cases to be distinguished. See also 2120 of 19.53 .

## RECEPTION

621.396 .621

1773
Logarithmic-Amplifier Simplifications and fmprovements-D. E. Sunstein and T. II. Chambers. (Proc. IRE, vol. 43, np. 343-344; March, 1955.) Coinment on 3345 of 1954 (Chambers and lage) and reply.

### 621.396 .621

1774
Statistical Survey of the Engineering Construction of Broadcast Receivers [in Western Germany]-W. W. Dietenbach. [Funk-Technik (Berlin), vol. 9, pp. 674-676; 1)ecember, 1954.]
621.396.621:621.376.3

1775
Reception of Frequency-Modulated Oscillations with Automatic Matching of Receiver Bandwidth to the Dynamic Range of the Modulation-K. Lamberts. (Fernmeldetech. Z., vol. 7, pp. 605-669; December, 1954.) For weak signals in wide-band noise, the af signal /noise ratio of a FM receiver can be improved by restricting the IF bandwidth. A circuit is described for varying this bandwidth automatically in proportion to the variations of the frequency deviation. An improvement of 14 db is (o)tained by $1: 4$ reduction of handwidth. The system does not eliminate impulsive noise.

### 621.396.621.54:621.314.7

1776
Transistor Broadcast Receivers-A. P', Stern and J. A. Raper. [Elec. Engng. (New York), vol. 73, pp. 1107-1112; December, 1954.] General design principles are discussed and circuit diagrams are shown of some experimental AM superheterodyne receivers using point-contact and junction transistors. Good quality reception is attainable.

### 621.396.82:621.376.3

Reception of an F.M. Signal in the Presence of a Stronger Signal in the Same Frequency Band, and other Associated Results-R. M. Wilmotte. (Proc. IELE, Part B, vol. 102, pp. 260-261; March, 1955.) Discussion on 1899 of 1954.
621.396 .828

1778
Interference Suppression-R. Davidson. (Wireless World, vol. 61, pp, 173-176; April, 1955.) Techniques for dealing with small commutator motors are described. Details are given of recently developed lead-through capacitors.

## STATIONS AND COMMUNICATION SYSTEMS

621.376 .2 1779
Tables of Bennett Functions for the TwoFrequency Modulation Product Problem for the Half-Wave Linear Rectifier-R. L. Sternberg, J. S. Shipman and W. B. Thurston. (Guarl. Jour. Mech. Appl. Math., vol. 7, part 4, p1. 505-511; December, 1954.) For previous work see 3028 of 1954 (Sternberg) and 2212 of 1954 (Sternberg and Kaufman).
621.376.5:621.39

1780
Average Spectrum of a Periodic Series of Identical Pulses Randomly Displaced and Dis-torted-IR. M. Fortet. (Elec. Commun., vol. 31, מ口. 283-287; December, 1954.) See 544 of February.
621.376.56:621.39

1781
Signal/Noise Ratio in Pulse Code Modula-tion-N. L. Yates-Fish and E. Fitch. (Proc. IEEE, Part B, vol. 102, pp, 204-210; March, 1955.) Formulas are derived for the output signal/noise ratio in simple pulse-code modulation syatems for all values of input signal/noise
ratios. The output ratio improves very rapidly with increasing input ratio provided the latter exceeds a certain critical valuc. The system is useful for links connected in tandem, since regradation of the overall performance below that for a single link niay be avoided by a relatively small increase of power in each link.

### 621.376.56:621.39

1782
Study of Pulse-Code Modulation-C Villars. (Tecih, Mill. schweiz. Telegr.-TelephVerze., vol. 32, pp. 449-472; December 1 1954. In French.) An account is given of an experimental installation developed in collaboration with C. Margna. A binary counter is used as coder, and 32 discrete amplitude levels are recognized on each side of zero. Two types of receiver were constructed, one in Which each pulse is treated separately and the other in which pulses are treated in groups, as described by Meacham and Peterson (2366 of 1948). Measurements of signal quality and of the improvement of signal/noise ratio between the hif and af channels are reported and compared with international standard requirements; the af bandwidth is quite satisfactory but quantization noise may be excessive. For multichannel communications the system compares well with others in respect of reliability, freedom from crosstalk, and ease of providing secrecy.
621.391 .1

1783
Prospects for the Development of Transmission Paths-E. Hölzler. (Fernmeldelech. Z. vol. 7, [1). 647-651; Decenber, 1954.) Review and comparison of various types of line and radio paths. Manufacturing difficulties appear to limit the bandwidth of tubular lines to about 10 mc . The possibilities of surface-wave lines at usw are discussed. Special types of dielectric and metal waveguides may prove suitable at frequencies above that ( $10-15 \mathrm{kmc}$ ) for which atmospheric absorption on radio paths becomes excessive.

### 621.395.44:621.315.28 <br> 1784

A Transatlantic Telephone Cable-Kelly, Radley, Gilman and Halsey. (See 1547.)
621.396.41:621.376.3

1785
Linearity Requirements for Multichannel F.M. Radio-Link Systems-G. Bosse. (Fernmeldelech. Z., vol. 7, pp. 678-682; December, 1954.) The relations between the measured distortion factor and the noise in frequencydivision inultichannel systems are calculated by substituting an equivalent noise voltage for the sum of the voltages in the channels. Nonlinear distortion due to curvature of modulator and demodulator characteristics and to transmission-time variations is considered. Diagrams are presented from which a determination can be made of the naximum permissible distortion factor for a given signal/noise ratio, and of the optimum amount of preemphasis

### 021.396.41:621.376.3

Problems of Frequency Modulat 1786 Multichannel Radio Links-H. Meinke in meldelech. Z., vol. 7, pp. 670-677; December, 1954.) The four principal criteria to be considered in deciding on the type of modulation to use in a multichannel system are (a) signal /noise ratio, (b) crosstalk, (c) bandwidth requirements, (d) equipment requirements. An account is given of experimental work on systems using individual-channel FM . Comparison is made with p.ph.m. systems. The results with the FM system are promising.

### 621.396.41:621.396.65

1787
Wide-Band Radio Links: Deliberations of Study Commission No. 9 of the C.C.I.R. at Geneva (from 10th to 22nd September 1954) W. Klein. (Tech. Mill. Schweiz. Telegr. TelephVerw., vol. 32, pp. 497-499; December 1 1954. In lirench.) Proposals for standardizing multichannel links are reported; both fre-
quency-division FM systems and time-di vision PPM systems a e considered.
621.396.41.029.62:621., 96.822.1

1788 Intermodulation Nise in V.H.F. Multichannel Telephone Systems-J. 1. Slow: (Jour. Brit. IRE, vol. 15, pp. 67-83; February, 1955.) Intermodulatior noise due to various forms of distortion in the radio circuits of a frequency-division mutiplex system is analyzed. Expressions are derived of the form $\bar{N}_{n}=\bar{\Pi}_{n}+n \bar{P}+C$, whe e $\bar{N}_{n}$ is the $n$ th-order noise power and $\bar{P}$ the multichannel speech power in db referred to $1 \mathrm{mw}, \bar{\Pi}_{n}$ is the $n$ thharmonic ratio of a test tone in db and $C$ is a constant. The analysis i; valid for FM systems handling up to 60 ch.innels. Formulas and curves are given and siugle-tone and two-tone tests are described for letermining the intermodulation noise due to (a) modulator/demodulator distortion, (L) phase distortion, (c) feeder mismatch.
621.396.61:621.396.66

1789
Operational Measurements on U.S.W [f.m.] Broadcast Transmitters-L. Merkl. (Arch. lech. Messen, nu. 227, pp. 269-272; Deceniber, 1954, and no 228, pp. 7-10; January, 1955.) Measurements considered include the monitoring of voltages and tube currents, rf output power, frequency and distortion.

### 621.396.665.1:621.396.65:621.376.3

1790
Transmission of Speech with Dynamic Compression-G. Hässler. (Fernmeldelech. Z., vol. 7, pp. 659-664; December, 1954.) The method of operation of the syllable compandor is described. Design advantages resulting from use of these compandors in multichannel carrier-frequency systems are indicated. The example of a $1 \mathbf{H}$ radio link is treated numerically.

## $621.396 .712 .029 .62+621.396 .61$

1791
Some Aspects of V.H.F. Sound Broadcasting and F.M. Broadcast Stations-P. A. T. Bevan. (Electronic Eng., vol. 27, pp. 96-101 and 147-153; March and April, 1955.) The relative nerits of systems using AM, AM with limiting, and FM are discussed, mainly on the basis of their effectiveness for suppressing various types of noise ar.d interference; field tests indicated the superiority of FM. A detailed account is given of the Wrotham highpower experiment. The cesign of FM transmitters is considered, with particular reference to modulators and monitoring. The antenna and transmission-line systems, parallel operation of FM transmitters, and unattended operation of transmitters are also discussed.

## 621-526

1792
Closed Expansion of the Convolution Integral (A Generalization of Servomechanism Error Coefficients)-E. Arthurs and L. H. Martin. (Jour. Appl. Phys., vol. 26, pp. 5860; January, 1955.)
621.311.6:681.142

1793
Precision High-Current Computer Power Supplies-A. B. Rosenstein. [Trans, Amer. IEE, Part I, Communicalion and Electronics, pp. 405-409; September, 1954. Digest, Elec. Engng. (New York), vol. 7 i, p. 1080 ; December, 1954.] A unit supplyin; 225 V dc at 15 a uses a magnetic amplifier-regulated Se rectifier.
621.319 .339

1794
A Portable Van de Graaff Generator- 1794 T. R. Foord. (Jour. Sci. Insth,, vol. 31, pp. $440-$ 441; December, 1954.) The generator described develons a maximun open-circuit volttage of about 200 kv and has a short-circuit current of $15 \mu \mathrm{a}$.
621.396.63:621.314.7
Practical Local Calling Circuit-(Shorl Wave Mag., vol. 12, pp. 55;-558; December, 1954.) The circuit/diagramocfan experimental
local-station calling device designed for operation on the $160-\mathrm{m}$ amateur band is given and discussed. The unit is basically a transistor receiver which operates a calling bell via a relaty. l'ower consumption is of the order of 0.5 kwh per annum.

## TELEVISION AND PHOTOTELEGRAPHY

### 621.397.5(44)

 1796The French Television Network-(Télévision, no. 49, pp. 305-307; December, 1954.) Details are given of revised frequencies for channels 1-12, and basic operational data for the various stations are tabulated.

### 621.397.5: 535.623].001.4

1797
Phase Measurement for Color TV and F.M -K. Schlesinger. (Electronacs, vol. 28, pp. 142-146; February, 1955.) The principles of operation with the vectorscone are described.

### 621.397.61:621.372.54

## Filters for Television Transmitter Diplezers

 -G. Meyer-Brötz. (Fernmeldelech. Z., vol. 7, pp. 683-688; December, 1954.) The requirements for separating filters used with combined sound and vision antenna systems differ from those for other diplexers because the bandwidth of the vision signal is large conlbared with the frequency separation of the two carriers. Various types of diplexer are surveyed, and the design of notch diplexers composed of coaxial lines is discussed particularly.
### 621.397.611:535.623

1799 and P 11 Bacoder Colorcasting-C. G. Lloyd and P. 11. Boucheron. (Radio-Electronic Engng. vol. 23, pp. 7-9, 35; December, 1954.) A description is given of a system employing emitron storage tubes for conversion of the se quential color signal, obtained from a monochrome television camera with a rotating color-segment disk in the lens system, into a NTSC-standard signal. See also 275 of leebruary.
$621.397 .7+621.397 .26$
1800
The Television Transmitter and Relay Installations at Antwerp-(Radio Rev. TV, vol 6, pp. 610-613; December, 1954.) The station, installed at the top of the 23 -story Torengebouw, oferates as a two-way microwave link between Brussels and Breda, at the same time broadcasting the received program in band I at a mean power of 2.5 kw . Sound is transmitted by cable. An outline description of the broadcast transmitter is given.
621.397.7:778.5

Considerations on the Operation of Vidi-
1801 graphs-Y. Angel. (Onde élect., vol. 34, pp). 958-973; December, 1954.) The term "vidigraph" is proposed for apparatus for the cinematographic recording of television programs from the face of a receiver tube. A particular system is described using a long-persistence screen. Probleins of obtaining correct contrast are discussed.
621.397.7:778.5:621.395.625.3

1802
The $16-\mathrm{mm}$ Substandard Film with Magnetic Stripe [for sound] as used in the Südwestfunk Television Service-Equipment and Operating Methods-H. Lauer and O. Schulze. (Tech. Hausmill. Nordzw Disch. Rdfunks, vol. 6, nos. 9/10, pp. 203-210; 1954.)
621.397.7.029.62:621.372.51

1803
V.H.F. Power Transmission Equipment for Band III Television Broadcast-B. M. Sosin. (Marconi Rev., vol. 17, pp. 88-100; 3rd Quarter 1954.) A descriptive account including some technical details on the construction and the characteristics of a system for linking television sound and vision transmitters to a common antenna. The system included a vestigialsideband filter, a frequency-discriminating combining filter, test load and feeder monitor-

### 621.397 .8

Various Factors affecting Picture Qualit Television. Possibilities of Picture Quality in F. Below. (Tech. IIausmill. NordwDisch. Rdfunks, vol. 6, nos. 9/10, pp. 195-202; 1954.) Deleterious effects due to bandwidth limitaltion and overshoot are cousidered. Methods of reducing defects due to the vestigial-sideband system of transmission are indicated. Crispening technique described by Goldinark and Hollywood (828 of 1952) and spectrum equalization methods described by Gouriet (1936 of 1953) are discussed. Improvements can be effected by reshaping or replacing the synchronizing pulses and by correct adjustment of level and gamma.

### 621.397.81:621.397.26

1805
Propagation on Bands I and III-F. W. R. Strafford and I. A. Davidson. (Wircless World, vol. 61, pp. 171-172; April, 1955.) A direct comparison has been ma-le of propagation in the two bands by radiating 180.4 -mic signals from the BBC mast at Suiton Coldfield, as well as the television waveform on 61.75 mc . Two receivers were installed in a mobile unit, and continuous records of signal strength were made. The receiving antennas were at a mean height of 25 feet, thus the difference between the local variations at the two frequencies could be investigated. Rapid variations due to reflecting objects and slaw variations possibly due to ground irregula-ities were observed. The significance of the latter for calculations of service area is discussed. The mean level of the band-III signal decreases with increasing distance faster than that of the band-I signal, as predicted theoretically:

## TRANSMISSION

621.396.61:621.314.7:621.311.6:621.383.5 1806 SPTTX [sun-powered transistor transmitter] Demonstration for N.P.L.--(Shorl Wave Mag., vol. 12, p. 557; December, 1954.) Brief note on a demonstration of the transmitter referred to in 1165 of May.

### 621.396.61:621.372.2 <br> 1807

Frequency Stability of Self-Excited Transmitters with Long Aerial Feeders-i. Kiich. (Arch. elekt. Übertragung, vol. 8, pp. 491-498 and 553-561; November and Decenber, 1954.) A stability criterion is derived whereby the influence of the transmitter is reduced to that of a single equipment constant which can be determined experimentally. The magnitude of this constant and the data of the feeder line uniquely determine the load conditions at the stability limit. For a given type of line the retroaction of the load on the transmitter is greatest for a line lengtl. giving a total attenuation of about 3 d b. The corresponding critical mismatch at the line termination is shown in normalized load curves. Experimental methods are described for determining the transmitter constant, and an indication is given of the maximum frequency shift to be expected. Agreement between theoretical and practical results is good.

### 621.396.61:621.376

1808
Phase-to-Amplitude Modulation 13. D. Virmani. (Wireless World, vol. 61, pp. 183-187; April, 1955.) High efficiency can be obtained with phase-to-amplituce modulation because (a) the phase modulation is performed at low level, and (b) the tubes in the two ph.m. rf channels can be contiruously driven to their limits. Details are given of a 400-w transmitter covering the frequency bands $3.58,13-30$ and $26-56 \mathrm{mc}$ and permitting either phase-to-amplitucle-modulation or ssb operation; the conventional oscillator and phase-shifting network are replaced by a polyphase oscillator, which retains the correct phase displacement when the oscillator frequency is varied.
621.396.61:621.396.712.029.62 1809

Some Aspects of V.H.F. Sound Broadcast-
ing and F.M. Broadcast Stations - Bevan. (See 1791.)

## TUBES AND THERMIONICS

## $621.314 .63+621.314 .7$

1810
Saturation Current in Alloy JunctionsW. M. W'ebster. (Jroc. IRE, vol. 43, pp. 277280; March, 1955.) Theory is dleveloned for diodes made by alloying circular junctions on to thin base wafers. An equation is derived from which the value of $I_{s}$, the saturation current obtained with reverse biasing, can be calculated. Most of this current originates from thermal generation at the free surfaces of the base. $I_{s}$ increases linearly with base resistivity and exponentially with temperature; it also increases, but more slowly, with base thickness and surface recombination velocity. The equation for the collector of an alloy-junction transistor is basically the same as for a diole, with a correction for the emitter. For the einitter junction the equation requires modification.
621.314.63:546.289

1811
Inductive Behaviour of $p-n$ Junctions in the Forward Direction-G. Kohn and W: Nonnonmacher. (Arch. elekl. Überiragung, vol. 8, pp. 561-564; December, 1954.) The observed time lag of the forward conductance of Ge dionles [. 3282 of 1952 (Einsele)] cannot be explained in terms of the usual equivalent circuit comprising parallel voltage-dependent resistance and capacitance. A circuit comprising resistance in series with parallel inductance and resistance fits the observations and is also consistent with the small-signal frequency variation of the diode impedance.

### 621.314.63:546.289

1812
Admittance Measurements on Alloyed Germanium-Indium Rectifiers-H. L. Rath. (Z. Nalurf., vol. 9a, pp. 699-700; July/August, 1954.) A method is described for determining the admittance from the characteristic variation of the capacitance with frequency at different temperatures.
$621.314 .632+621.314 .7] \cdot 002.2$
1813
Transistors and Germanium Diodes[Elect. Rev., (London), vol. 155, pp. 791-795; November, 1954.] An account of procluction techniques used in Britain for the manufacture of point-contact crystal tubes on a comparatively large scale.

### 621.314 .7

1814
Transistors and their Applications-a Bibliography, 1948-1953-A. R. Krull. (Trans. IRE, vol. ED-1, pp. 4D-77; August, 1954.)

### 621.314 .7

1815
Electrical Characteristics of Power Transis-tors-A. Nussbaum. (Proc. IRE, vol. 43, pp. 315-322; March, 1955.) Mcasurements have been made on $p-n-p$ junction transistors with collector dissipation of about 20 w [3391 of 1954 (Roka et al.]. The results relating to parameters which are a function of current strength do not agree with those obtained theoretically for low-power operation [3390 of 1954 (Rittner)]. Plans for further investigations are indicated.
621.314.7:537.311.33:546.28 1816

Electronic Behaviour of Certain Grain Boundaries in Perfect Crystals-Mataré. (See 1690.)
621.314.7:621.317.755

1817
An Alpha Plotter for Point-Contact Transistors T. 1’. Sylvan. [Elec. Engng. (. Yew York), vol. 73, pp. 1094-1098; December, 1954.] Description, with detailed parts list, of a cro test set.
621.383 .27

1818
The Transit Time of Electrons in Photo-multipliers-E. H. Rhoderick and R. W. P. McWhirter. (Jour. Sci. Instr., vol. 31, p. 475;

December, 1954.) Experimental results on commercial 11- and 13-stage photomultipliers indicate transit times ranging from about 5.5 to $8 \times 10^{-8}$ seconds at voltages between 185 and 105 v per stage.

### 621.383 .5

1819
Electronic Interpretation of Inertia Phe~ nomena in Photocells (in particular, Internal Capacitance)-G. Blet. (Jour. Ihys. Radium, vol. 15, pip. 823828 ; December, 1954.) See 2801 of 1954.

### 621.383 .5

Barrier-Layer Photocells: Part 1-D. Giest. (Arch. lech. Messen, no. 227, pp. 281284; December, 1954.) A brief account of the mechanism and operational characteristics of $p-n$ junction and metal-semiconductor barrierlayer photocells. 35 references.
621.385.029.6

1821
The Traveling-Wave Tube-a Record of its Early History--R. L. Wathen. (Jour. Frank. Insl., vol. 258, pp. 429-442; December, 1954.)
621.385.029.6

1822
A Large-Signal Theory of Traveling-Wave Amplifiers -P. K. Tien, L. K. Walker and V. M. Wolontis. (Proc. IRE, vol. 43, pp. 260277; March, 1955.) Analysis presented by Nordsieck (2497 of 1953) is extended to cover space-charge effects at high operating levels. In adelition to the space-charge parancter $Q C$ and the other parameters used in linear theory, a parameter $k$ is introduced such that $1 / k$ is proportional to the beam radius and gives an indication of the range of action of the space-charge forces. Computations have been made for a number of typical cases; anplitucle and phase of the circuit wave are given as functions of distance. The limiting efficiency increases with electron injection speed, increases first and then decreases with increase of $Q C$, and increases with $1 / k$, assuning uniform distribution of fied and current over the cross section. The electron motion is analyzed.
621.385.029.6

1823
The Performance of Travelling-Wave Valves at High Input Levels W. Klein and W. Friz. (Fernmeldelech. Z., vol. 7. pp. 349-357; July, 1954.) The helix is considered as a succession of elenentary zones, in each of which linear amplification theory is applicable, and the power output is determined in successive steps. Losses are taken into account, and graphs show the dependence of other parameters on the velocity and loss parameters. The effect of an attemating section on amplification and the optimum choice of its position are examined. The expressions developed are interpreted with reference to a particular tube and the output-power/input-power curve discussed in detail. The importance of the coupling factor $C$ in predicting tube performance is stressed. Amplification at low and high input levels is compared, and the effects of operating voltage and type of helix considered. An estimate is made of the greatest power output compatible with linear operation. Computed and measured values are in reasonable agrecment in all cases.

### 621.385.029.6

1824
Traveling-Wave-Tube Characteristics for Finite Values of $C$-C. K. Birdsall and G. R. Brewer. (Trans. IRE, vol. ED-1, pp. 1-11; dugust, 1954.) Values of the preformance parameters of traveling-wave tubes are presented in the form of curves for values of the gain parameter $C$ up to 0.5 and for relatively large values, up to 2 , of the space charge parameter $Q C$.

### 621.385.029.6

1825
Cross-Wound Twin Helices for TravelingWave Tubes-M. Chodorow and E. L. Chu. (Jour. Appl. Phys., vol. 26, pp. 33-43; Janu-
ary, 1955.) When a single-helix structure is used with voltages above about 10 kv , a large proportion of the energy is diverted into noninteracting space harmonics, with a corresponding reduction of the impedance for the fundamental interacting mode; undesired backward waves may also be generated. Use of a twin helix climinates these drawbacks. For a particular twin helix with radius $0.4 \lambda / 2 \pi$, the impedance for the fundamental mode was more than twice that of a single helix of the same radius. Dispersion is greater with the twin helix.
621.385.029.6

1826
Electronic Resonance Effect in Valves with Crossed Electric and Magnetic Fields-W. Willshaw, G. Mourier and G. Guilbaud. [Compt. Rend. Acad. Sci. (Paris), vol. 240, pp. 283285; January 17, 1955.] A formula is presented expressing the conditions under which electronic resonance is liable to disturb the normal operation of traveling-wave tubes using lines with periodic structure. A large effect is produced when the resonance corresponds to a space harmonic with high-intensity field. Measurements of the variation of efficiency of al carcinotron over the electronic tuning range are discussed in the light of the theory.

### 621.385.029.6

1827
Excitation of the Carcinotron $M$ Valve- G . Mourier. [Compt. Rend. Acad. Sci. (Paris), vol. 240, pp. 406-408; January 24, 195.5.] An investigation is made of the current build-up in back-ward-wave oscillator tubes of the type discussed by Guenard et al. (3616 of 1952); the calculations are based on the energy exchanges between beam and delay line. A coefficient $Q_{a}$ is introduced defining the excitation quality; an expression is given for $Q_{a}$ depending on delay-line length, operating wavelength, electron and wave velocities, and ratio of instantaneous current to current in the oscillating state. The excitation time varies in the same sense as Qa.
621.385.029.6

1828
Influence of Space Charge on the Excitation Current of a Carcinotron-Type Magnetron Oscillator-B. Fpsztein. [Compl. Rend. Acad. Sci. (Paris), vol. 240, pp. 408-410; January 24, 1955.] Values found experimentally for the excitation current in the carcinotron $A I$ tube are generally lower than values derived from the formula of Guénard et al. (3616 of 1952) and Warnecke et al. ( 3085 of 1954). A new formula is developed taking account of space charge effects; satisfactory agreement is then obtained between the theoretical and experimental values.
621.385.029.6

1829
Resonant Behavior of Electron Beams in Periodically Focused Tubes for Transverse Signal Fields-R. Alder. O. M. Kromhout and P. A. Clavier. (Proc. 1RE, vol. 43, pp. 339341; March, 1955.) Transverse-field travelingwave tubes are discussed in which a ribbon beam is passed between pairs of plates at alternate high and low potentials; resonance is observed over a range of focusing conditions. A formula is derived for the resonance frequency; which is equal to the plasma frequency.

### 621.385.029.6

1830
Magnetic Focusing of Electron BeamsJ. T. Mendel. (Proc. 1 RE, vol. 43, pp. 327331 ; March, 1955.) Analysis is developed applicable to various types of focusing used for pencil beams as in traveling-wave tubes. The semi-shielded cathode offers practical advantages, in conjunction with either uniform or periodic focusing fields. Expressions are derived for the electron trajectories. Balance conditions yielding minimum beam ripple are determined; these require a high percentage of flux threading the cathode for a relatively small increase in magnetic field above the Brillouin value.
621.385.029.6

1831
Hysteresis in Klystron Oscillators-T. Moreno and R. L. Jepsen. (Proc. IRE, vol. 43, p. 344 ; March, 1955.) A possible explanation of electronic hysteresis is advanced based on an electronic-admittance/rf-voltage characteristic which is not monotonic.
621.385.029.6:621.372.413

1832
Stabilization of Reflex Klystrons by High- $Q$ External Cavities-S. J. Rabinowitz. (Trans. $I R E$, vol. MTT-2, pp. 23-26; September, 1954.) The effective $Q$ of the klystron oscillator is raised by associating a high- $Q$ external cavity with it. A suitable design of stabilization cavity and coupling network is illustrated, and experimentally determined characteristics of some stabilized klystrons are presented.
621.385.029.6:621.396.822

1833
Noise in Transverse-Field Traveling-Wave Tubes-G. Wade, K. Amo and D. A. Watkins. (Jour. Appl. Phys., vol. 25, pp. 1514-1520; December, 1954.) Analysis indicates that in a transverse-field traveling-wave tube the noise can only be kept low if the beam is well collimated. In a particular tube for operation at 1 kmc , with a collimator of width 0.004 inch, the theoretical noise figure is 2 db ; without the collimator the noise figure would be 11 db .
621.385.029.63/.64:621.372.2

1834
Theoretical Study of Traveling-Wave Tube -K. Udagawa. (Rep. Elect. Commun. Lab., Japan, vol. 2, pp. 34-52; August, 1954.) Expressions are derived for the propagation constants of a coaxial arrangement of (a) two helices, and (b) a helix and dielectric cylinder. Series of curves are plotted showing the influence of the thickness of the dielectric cylinder on the phase velocity, and the characteristic impedance and attenuation of a traveling-wave tube helix with a dielectric support. Gain is calculated taking account of thermal effects; space-charge effects and transverse electron motion are investigated.
621.385.029.63/.64:621.396.822

1835
Influencing Space-Charge Waves of Fluctuating Beams by Resonator Circuits-K. Pöschl. (Frequenz, vol. 8, pp. 284-288; September, 1954.) Theoretical considerations show that the noise factor in a two-stage klystron can be reduced below the minimum given by Robinson ( 1618 of 1954) by using an additional resonator ahead of the input stage.

### 621.385.029.65

1836
An Experimental Broad-Band Helix Travel-ing-Wave Amplifier for Millimeter Wave-lengths-S. D. Robertson. (Trans. IRE, vol. MTT-2, pp. 48-54; September, 1954.) A tube with a helix of diameter 0.015 inch has given 19 db gain at $6 \mathrm{~mm} \lambda$ and 9 db at 5.2 nmm . Design problems are discussed. An anode potential of 1 kv and a beam current of 10 ma are used.

### 621.385.029.65:621.317.361

1837
Cold Measurements of 8 mm Magnetron Frequency and Pulling Figure-A. E. Barrington. (Proc. IEE, Part B, vol. 102, pp. 247-248; March, 1955.) A square-wave signal is applied to the magnetron, and voltages proportional respectively to the input and reflected signals are applied to the $x$ and $y$ plates of a cro, giving a straight-line trace whose slope is minimunt at resonance. The change of resonance frequency with adjustment of the position of a puller probe is observed.

### 621.385.032.21

1838
Influence of Space Charge in Spherical Electron Guns-N. B. Aizenberg. (Zh. lekh. Fiz., vol. 24, pp. 2079-2082; November, 1954.) While it is usually assumed that the intensity of the field at the tip of the cathode in a spherical electron gun is proportional to the anode voltage, for large values of the discharge cur-
rent the space charge must be taken into account. The effect of this charge is investigated experimentally and the minimum value of the discharge current for which the effect becomes noticeable is determined.

### 621.385.032.21

New Forms of Thermionic Cathode[Nature, (London), vol. 174, pp. 1176-1177; December 25, 1954.] Report of a colloquium held in October, 1954 and sponsored by the Institute of Physics and the Physical Society. The discussion dealt mainly with efforts to produce cathodes giving high emission over long periods.
621.385.032.21:537.29

1840
Progress in Electron Emission at High Fields-W. P. Dyke. (Proc. 1 RE, vol. 43, pp. 162-167; February, 1955.) The properties of cold and hot field-emission cathodes yielding high current densities are surveyed and methods of stabilizing their performance are discussed.

### 621.385.032.216

1841
A New Type of Diffusion Cathode-A. H. Beck, A. B. Cutting, A. D. Brisbane and G. King. [ Nature (London), val. 174, pp. 10101011; November 27, 1954.] Brief details are given of a cathode made by molding a mixture of Ni powder and ( $\mathrm{Ba}, \mathrm{Sr}$ ) carbonate powder with a small percentage of reducing agent. The pulsed emission is shown as a function of temperature in comparison with the conventional oxide cathode and with published data for the L cathode. From determinations of the work function it is inferred that the emission originates from an incomplete monolayer of Ba ions with about 70 per cent surface coverage. For a fuller account, see Le l'ide, vol. 9, pp. 302-309; November, 1954.

### 621.385.032.216

1842
Contaminated-Metal Sintered Thermionic Cathodes-R. Uzan and G. Mesnard. [Compl Rend. Acad. Sci., (Paris), vol. 239, pp. 1613 1615; December 8, 1954. . New techniques are outlined for preparing high-emission molded cathodes combining metal and coxide powders in such a way that the surface is entirely of metal. Good results are obtained with Ni when the oxide content is only just sufficient to provide the necessary diffusion to the surface; the sintering may be performed in vacuum or in a hydrogen atmosphere. W gives lower emission than Ni at low temperatures but greater emission at temperatures over 1,250 degrees K ; sintering in hydrogen lowers the emission with W. See also 3728 of 1955.

## MISCELLANEOUS

$621.396 .6+621.396 .712$
1843
Development Work of [West] German Broadcasting Institutes-H. Rindfleisch. (Elektrotech. Z., Edn A, vol. 75, pp. 587-590; September 11, 1954.) A brief survey of equipment and techniques developed since the war, and particularly since 1951. 44 references.

### 621.396 .97

1844
Recording and Tabulating the Radio-TV Audience-A. C. L. Brown. (Electronics, vol. 28, pp. 126-129; January, 1955.) Listener research is conducted by means of film records indicating the periods during which the receiver is tuned to different channels. The recording system is applicable for ather purposes, e.g. investigation of atmospheric variations.

## $413=30=20$

1845
English-German Technical and Engineering Dictionary [Book Review]-L. de Vries. Publishers: McGraw-Hill, London, 997 pp., £7. (Wireless Eng., vol. 32, p. 90; March, 1955.) A companion volume to the already published German-English Dictionary. Includes many new terms in the field of radar, television, nuclear engineering, etc.

# Gulton abstracts 



Spot-checking of prototype acrelerometers on electro-mechanical vibrators

## Cooperative Program for Hi-Temp Shock and Vibration Measurement

Up to the present, engineers have been severely handicapped in shock and vibration measurement problems under conditions of high temperature owing to the inability of available accelerometers to withstand the excessive heat generated under test. Cognizant of the tremendous need of designers for assistance in this respect, particularly in the aircraft and missile development fields, Gulton Mig. Corp. is announcing the allocation of facilities and staff for cooperative research and work on high temperature measuring problems. Shortly, Gulton is planning to extend its line of high-temperature pirzoelectric accelerometers so that they will be available as stock items for both industry and the military. Meantime. during the design stages of the newer type accelerometers. inquiries are solicited from engineers who are facing serious high temperature measuring problems of any kind.

As a result of this intense program for the development of accelerometers for high temperature work. Gulton is now making availallle a new unit which promises to provide data on severe environmental conditions previously impossible
to ascertain. This new accelerometer, the first of a newly-designated AHT series, is based on the use of new techniques in mechanical design plus cancentrated research on piezoelectric ceranic materials relatively unaffected by temperatures from below $-70^{\circ}$ to above $500^{\circ} \mathrm{F}$.
Construction of the accelerometer involves special temperature-stable housing materials and a new cable fabrication to withstand the high temperatures. In conjunction with Glenco Corporation. a new piezoelectric ceramic, formulated and fired by unique techniques, maintains its response well above operating range. The new unit does not require cooling fins or liquid circulation systems. and can be operated continuously at these elevated temperatures. Owing to its high temperature characteristics, in addition to its miniature size, it will make possible the measurement of shock and vibration in many types of devices heretofore impossible to consider with existing instruments.
From prior experience with these elevated temperatures, Gulton engineers are probing the higher heat spectrums to develop further accelerometer designs

## New Thermistor Mountings Preserve Control Characteristics Under Large Power Loads

A limiting factor in temperature compensation problems and other thermistor applications has been undesirable temperature rise of the thermistor itself due to electrical self-heating.
Large power handling ability is a feature of two new thermistor mountings developed by the Thermistor Corporation of America. Both types are particularly useful for temperature compensation of transformers, small motors. coils, relays, and resolvers. The upper one pictured is a thin ceramic thermistor soldered directly to a metal nounting plate which serves as one terminal. The thermistor is about $1 / 100$ inch thick; much thinner than anything previously available. The lower unit is comprised of a thin thermistor embedded in a copper bracket and is designed for circuits where neither side can be grounded.


The new mountings are used to provide close thermal coupling between the thermistor and the device to be compensated, as well as to minimize errors of compensation from electrical self-heating of the thermistor. The resulting high thermal dissipation constants have been achieved by an exclusive patented process for manufacture of extremely thin ceramic thermistors.

For further information about these thermistor units, write on your firm letterhead to the Thermistor Corporation of America, Metuchen, $\mathbf{N}$. J.
that will provide even better characteristics than the new series. If these high ranges are now affecting your work, you are most urgently invited to write now to the Director of Engineering, Gulton Mfg. Corp., Metuchen, N. J. You are under no obligation, and you are assured of a prompt, competent evaluation.

[^64]


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## beathkit iv

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Woodworth, J. D., 20) Cedar St., Filberon. N. J. Yamagami. Y., 25 Prospect Ave., Montclair, N. J. Yang, C. C., 40-07-193 St., Flushing, L. I., N. Y. Yaw, D. F., 734 Melrose Ave., Columbus 11, Ohio Yokelson, B. J., Bell Telephone Laboratories, Inc., Whippany, N. J.
Zaslavsky, S., 740 E. Gun Hill Rd., New York 67. N. Y.

The following admissions to the Associate grade were approved to be effective as of June 1, 1955:

Abel. A. O., 100 Memorial Dr., S. 2-20A,Cambridge, Mass.
Abeson, I. S., 236 Keller St., Monterey Park, Calif, Adams, W. I... 7413 Parkwood Dr., St. Louis 16, Mo.
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Aeronautical*
Chairman Louis B. Rothschild of the Air Coordinating Committee has announced that the group has unanimously approved a program to be followed in connection with the current controversy surrounding the common system of shortdistance radio navigational aids. "In view of the military's planned implementation of a tactical system," the ACC anmonerment said, "the clivergence between the common civil/military nem-tactical system and the tactical military system becomes obvious. Thus, a course of action must be devised which will minimize disruption to all aviation interests during this divergence, and arrive, to the extent possible, at a common civil/military systom of navigation which provides for cisil/military mon-tactical, as well ats hasic tactical operations." The controversy was touched off earlier this year when the Xir Navigation Development Board issucd a report in which it favored the TiACAN system developed by the military over the Vor 1)ME system recommerded loy the Radion Technical Commission fo- . Deronantios and accepted as the common eystem. The matter has been probed by several committees on Capitol Hill and a resolution is now pending 10 establish a joint commitere to investigate the matter. The ACC: program was reported in part as follows: . I. Interim Military Tactical Program (1) The military will proced immediately to imple. ment the minimum amonnt of TACAN necessary to meet military tantical require-
(Continuct on paze 92.t)

* The data on which these Notes are based were selected by permission from Iridustry Reports. issines of April 18, 25, May 2, and 9, puslished by the Radio-Flectronics Television Minnufacturers Association, whose hulpfulness is gratenully acknowl-
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## (Continued from fage 90A)

ments. These installations must avoid interference with the Radar Safety Beacon and DME. (2) The proposed use of TACAN must be coordinated through the ACC agencies, and aircraft flights based on the system must be capable of being integrated into the operation of the Federal Airways to minimize operational conflicts. Implementation of TACAN will be in accordance with policy to be developed by ACC. B. The Common System Program (1) The CAl continue the loR /I)ME system, as hercinafter provided, until some succeeding common system has been adopted and installed on the Federal Airways system. During the transition period when a succeeding common system is being installed, priorities in frequencies or other areas of confict shall be given to the succeeding system. (2) The agencies of the government responsible for the implementation or operation of any phase or phases of the present or a future common systen program will consult with other interested agencies through the ACC in the discharge of these responsibilities. Any implementation of INME will be carried out only in proper priority relationship to the other needed improvements in the system. (3) Based on information :ow available, if TACAN is adopted for use in the succeeding common system, it is estimated that under such a program, lOR will be continued until 1965 and that DMIE will be continued until 1960. The A.NDI3 also was directed to study several "undetermined factors" in connection with the use of TACAN and "immediately plan and direct a program to complete the development of TACAN for possible common system use." It also was directed to "conduct studies to determine the feasibility of developing a third rho-theta system which would meet all of the stated common system reguirements." The ACC statenent also covered the international program to be followed.

## FCC Actions

The Federal Communications Commission has issued a public notice soliciting information which can serve as the basis for determining whether a rule-making procedure should be instituted concerning radio-astronomy frequency requirements. The Commission stated that in view of the widespread interest and work being done in the field of radio-astronomy, both in this country and many others, it considers it expedient to develop at this time additional information and hats listed some six points upon which comments are solicited from interested parties boy July 1.

## Electronics

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| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\begin{aligned} & 6.05 \times 104 \\ & 2 \times 104 \end{aligned}$ | $41^{\circ} / \mathrm{hr}$ | $\begin{aligned} & \text { A.C.Vane } \\ & \text { A.C.Vane } \end{aligned}$ | $\begin{aligned} & 0.17 \% \\ & 0.1 \% \end{aligned}$ | $\begin{aligned} & .0035 \mathrm{Sec} . \\ & .0025 \mathrm{sec} . \end{aligned}$ | 3-3/4" Diam. $\times 6-1 / 8^{\prime \prime}$ tong <br> $2^{n}$ Diarn. $\times 3-7 / 8^{n}$ long | - $6-1 / 2$ |

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(Continued from Page 22A)

plans which look toward the expansion of America's early warning system and modernization and expansion of the aircraft program. Air Förecoblignt ionls during fiscal year 1056 for the procuremont of electronics and communications erpapment are experced to total mearly So 25.9 million. This compares with obligations of about St27.6 million in 1955 and $\$ 326.4$ million obligated for this type of equipment in fiscal year 1954.

## MOBH.IZATION

Atomic Energy Commission Chairman Lewis Strauss has announced a new program under which organizations or individuals may be given access to non-military "confidential" and "secret" restricted data on atomic energy technology for their own private purposes. Thder the new program, information classified as "confidential, restricterl data" may be made available to any person who can evichence a potential use or applieation of the data in his business, profession or trade. The ot her conditions to this arcess of information are that the applicant obtain a simplified security "L" clearance and agree in writing to conform to all AEC security regulations. Aso, so-called "study agrements" now in effect will be comwerted to the mew type arrangements. I'nder similar con. dit ionts, limited aceess also may be granted to everain specific information classified ats "secrev, restricted data" if the appliant proves that the information is significantly important to his business provided, however, that the applicant obtains a full security" O " charance. The government will retatio rowaltyfree, non-exclusive rights for gremomental purposes in imentions and discoveries which result from such access. Present AEC contractors will be granted access to both categories of restricterl data for private purposes on the same basis and conditions and will be franted the same rights as other applicatus.

## RETMA

Climaxing a three-day industry conference, the RETMA Board of Directors approved proposals for broad administrative changes in the Association's organizational structure, subject to approval by the membership at the June convention, and selected Director Leslie F. Muter, pioneer radio manufacturer and veteran RETMA Treasurer and Past President, to receive the 1955 Medal of Honor at the industry banquet on June 16 in Chicago.

## Tecinilical

The Office of Technical Services, Commerce Department, recently released several reports dealing with scientific discoveries and developments of interest to the
(Cominmed on paze 97A)


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electronics industry, and in addition made available to the public the "AN" (ArmyNavy) nomenclature system for communications and electronics equipment, devised by the Joint Communications-Electronics Committee of the Defense Department. "Properties of Large Slot Antennas" P'B 111523, and "A High Vacumm Gold Cathorle Gaseous I ischarge-. I Cyclotron Eiffect" PB 111522, are both available from the OTS, Commerce lepartment, Washington 25, D, C., for $\$ 1.25$ each. "The Summary of Joint Nomenclature System ("AN" Systen) for Commuilica-tions-Electronic Equipment* Pl 111581, also is available from the Commerce I epartment, for 25 cents per copy. This report presents a coordinated system of nomenclature for communications-electronic equipment, jointly used by the three military services, in a handy chart form, 'The system is useful, it was said, in identifying equipment references where reference is made to a new item, or a complete listing is not at hand. The coding system used in the charts indicates at once the type of equipment, where it is normally installed, and its functional purpose.... Technical details on the system for mass production of electronics-known as the "Modular Design of Electronics" and the "Mechanized Production of Electronics" -have been made available to industry in five reports published recently by the Office of Technical Services, Commerce Department. The reports include a summary description of the system, techniques of conversion from conventional to modular design, hand fabrication techniques, mechanized production, and manufacturing cost determination. The system of producing electronics equipment mechanically is not limited to large-scale production, it was said. It may be applied equally well to model-shop or laboratory practices. It is suggested, in fact, that hand assembled electronic models be developed and performance tested before large-scale production by MIPE techniques is attempted. Through the application of this system it is possible to substantially recluce the lead time normally required before full production is attained and to stock-pile production facilities, it was noted. It requires no proportionate increase in skilled manpower when the standly plant is put into operation, OI'S said. The five volumes describing this system are listed below. They may be obtained from the OTS, Commerce Department, Washington $25, \mathrm{D}$. C. at the prices indicated. "Vol. 1, Summary of Modular Design of Electronics and Mechanized Production of Electronics," 113 111275, \$2. "Vol. 2, Techniques for Converting from Comventional Desiqut of Eilectronics to Nodular 1)esign of Electronics," PB 111276, \$2. "Vol. 3, Hand Fabrication Technique and Photographic Processing for Nodular Design of Electronics," P'B 111277, \$2. "Vol. 4, Mech-

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（Continucd from page 97A）
anized I＇rolnction of Feleetronics，＂IV 111278，\＄t．＂Vol．5，Manufacturing Cost Determination，＂［＇I3 111315，\＄4．

The Office of Technical Services， Commerce Department，has listed stud－ ies in the field of electronics in its March 1955 issue of the＂ U ．S．Government Research Reports．＂The following re－ ports can be purchased from the l＇hoto－ duplication Section，Lilurary of Congress， Washington 25，D．C．，for the reported price：＂Design and Calibration of Micro－ wave Antenna（Bain Standards＂I＇l3 116133，microfilm，$\$ 2.50$ ；photocopy， \＄5．25．＂Comparison of Linear and Circular Polarization at X－hend by Means of a CIV Doppler Radar Operating Over Water＂ PB 116180，microfilm，\＄1．50；photocopy， \＄1．50．＂Aerodymamic and Radar Trans－ missivity Properties of Screen Materials＂ Pl3 116249，microfilm，$\$ 3.25$ ；photocopy， \＄9．＂Folded Antennas＂I＇l3 116294 ，micro－ film，\＄t：photocopy，\＄11．50．＂General Study of Rectangular Waveguide（pres－ surized）＂l＇l3 116171，microfilm，\＄2．50； photocopy，$\$ 5.25$ ．＂Noine Studies on CW Klystrons＂P＇B 116250：microfilm，\＄1．．50）； photocopy，$\$ 1.50$ ．＂On the Theory of Wave propagation in Nom－Homoge：neous Media＂ I＇B 116185 ，microlilm，\＄2．25；photocopy，St． ＂On the Perturbation Theory of Filectro－ magnetic Cavity Resonators＂P13 115744， microfilm，\＄2；photocope，\＄2．7．5．＂Practical Transmission line Networh Design for VIF and IHHF Filtor Applications＂PIS $1159+2$ ，microfilm，$\$ 5.25$ ；photocopy， \＄16．50．＂Nicrowave Noise Study＂I＇B 116251，microfilm，\＄3．75；photucopy， $\$ 10.25$ ．＂Study＂of R－F I＇erformanere Meas－ urements－Final Repert＂PB 11607．3， microfilm，\＄5．25；photoropy，\＄16．50．＂Sy＇n－ thesis．Final Report＂I＇13 116102，micro－ film，\＄2．25；photocop！：\＄t．＂Trouble－ Shooting in Ekectronics Equipment－． l＇roposed Method＂PB 116207，microfilm， $\$ 5$ ；photocopy，$\$ 15.25$ ．＂I＇se of Real Gases in a Shock＂Tube＂I＇B 116211，microfilm， \＄2．75；photocopy＇，\＄77．5．．．The Air Force has recently released the results of a program of basic research in the field of nonlinear servomechanisms．The results are described in a research feport made available to indastry by the Office of Technical Services，Commere Depart－ ment．The report，＂Research in Nonlinear Mechanies as Applied to Servomechat－ nisms，＂is the result of research done by the Cook Electric Co．L．aboratories under Wright Air levelopment（enter contract． The work was almod at and improsement of servomechanism response through the use of nonlinear techmiques，development of practical nonlinear elements to instrument these technifues，extention of nonlinear serromechanism theory through theoret i － cal analysis，and formularion of monlinear servo theory into practica！design informa－ tion for use by design enginesers．The re－ port is availathe from the orss，Commerce Department，W：ashington 25．D．C．，and should be ordered by nutolser－1＇に $11158+$ —priced at \＄3．75．

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Literature

Antennas and I'ropagation
Washington Chapter-March 28
"Ionospheric Propagation" by Alvin McNish, National Bureau of Standards.

## Audio

## Boston Chapter-April it

"The Past. Present and Future of Magnetic Recording" by John S. Boyers, Magnetic Memory I levices Division, National Company.

Houston Chapter-March 22
"A Multi-Loop Self Balancing Power Amplifier" by J. Ros: MacDonald, Texas Instruments, Inc.

Philadelphia Chapter-March 17
"Experiences and Observations Along the Road to Improved Sound Reproduction" by E. W. Kellog, RCA (retired).

## Chicago Chapter-February 18

"Toward Flutter Measurements of Magnetic Recorders" by U. R. Furst, Furst Electronics.

## Circuit Theory

Albuquerque Chapter-March 2.3
"Circuit Synthesis with Particular Reference to "Transistors" by Walter Brown, Sandia Corporation.

Chicago Chapter-January 21
"RF Spectra as Kelated to Non-Linear Circuit Elements" by William Firestone, Motorola, Incorporated.

Seattle Chapter-February 10
"Digitai Computers and Automatic Control" by 1). R. Firown, M.1.T.

Syracuse Chapter-March 15
"Active Filters" by J. J. Suran, General Electric.

## Communications Systems

Washington Chapter-March 30
"Disaster Planning in the Field of Telecommunications" by Horace R. Hampton, Cand P Telephone Company.

COMPONENT PAETS
Los Ingeles Chapter-March 14
"The Metal Film Potentioncter" by D. W. Moore, Servomechanisms, lic.
"Instrument Switches" by C. Broneer and G. Edwards, hoth of Aerovos Corporation.


# 23 Fields of Special Interest - 

## The 23 Professional Groups are listed below, together with a brief definition of each, the name of

## ACTIVITIES

The IRE Professional Group has the responsibility of providing the individual with the advantages of a small, select society in the field of his specialization, with its own magazine, just as IRE provides him with the advantages of a large, general society. The advantages of the small society relate primarily to meetings and to publications. Specialized symposia may be arranged either to coincide with IRE Conventions or to occur where there are places of large activity in the field of interest.

The Group is concerned with the advancement of scientific engineering leading to increased professional standing in its field and serves to aid in promoting close cooperation and exchange of technical information among its members. It provides a form for discussion and presentation of papers on subjects of mutual interest, and provides smaller, more compact Groups who may meet on the common basis of professional interests.

## ORGANIZATION

The IRE Professional Group is established under a constitution within the framework of the IRE. The constitution defines the technical field of interest of the Group, establishes committee structures, describes broadly its functions and procedures, and fixes a minimum level of activity. The management of an IRE Professional Group is in the hands of its Administrative Committee, the officers and members of which are elected annually. The IRE provides financial assistance to the Groups in accordance with their activity and current needs.

## PUBLICATIONS

Every Group publishes a magarine which is called TRANSACTIONS of the Professional Group, generally on a regular quarterly schedule. The TRANSACTIONS serve to preserve and disseminate the body of knowledge that constitutes the fields of interest of the Groups. All editions are distributed without additional cost to members who have paid the annual assessment.

The CONVENTION RECORD covering the sessions presented at the IRE National Convention is furnished without further charge to the members of Groups who have paid assessments.

## Circuit Theory

Design and theory of operation of circuits for use in radio and electronic equipment.

Dr. Herbert J. Carlin, Chairman, Microwave Research Institute, Poly-
technic Institute of Brooklyn, 5
Fee \$2.7 Transactione. "1, "2, "Vol. CT-1, Nos. 1-4; CT-2, No. 1 .

Fee $\$ 2.9$ Transactions. "1, "2, "3, "5, "6, "7,
$8 ;$ BTR-1, No.
8,

Fee \$2. 24 Transactions, 4 Newsletters. ${ }^{-5}$, *7, "10. "Vol. AU-1, Nos. 1-6; "Vol. AU-2, Nos. $1-5$; Vol. AU-3, Nos. 1-2.

## Broadcast Transmission Systems

Broadcast transmission systems engineering, including the design and utilization of broadcast equipment.

> Mr. Oscar W. B. Reed, Jr., Chairman, Jansky \& Bailey, 1735 DeSales
man, Jansky \& Bailey, 1735 DeSales
St., N.W., Washington, D.C.
Fee \$2. 1 Transaction, No. 1.

## Audio

Technology of communication at audio frequencies and of the audio portion of radio frequency systems, including acoustic terminations, recording and reproduction.

Mr. Winston E. Kock, Chairman,
Bell Telephone Laboratories, Inc.,
,

## Communications Systems

Radio and zire telephone, telegraph and facsimile in marine, aeronautical, radio-relay, coaxial cable and fixed station services.

Mr. Arthur C. Peterson, Jr., Chair-
man, Bell Telephone Laboratories,
463 West Street, New York 14, N.Y.
Foo $\$ 2.4$ Transactions. 5 Newsletters. "Vol. CS-1, No. 1; Vol. CS-2, Nos. 1-2; CS-3,

## Component Parts

The characteristics, limitation, applications, development, performance and reliability of component parts.

Mr. Floyd A. Paul, Chairman, Ben-
dix Development Lab., 116 W . Olive Ave., Burbank, Calif.
Fee \$2. 3 Transactions. *PGCP-1-2-3.

## Aeronautical and Navigational Electronics

The application of electronics to operation and traffic control of aircraft and to navigation of all crafi.

Mr. Edgar A. Post, Chairman,
United Air Lines, Operations Base,
Stapleton Field, Denver 7, Colo. Fee $\$ 2.13$ Transactions, 4 Newsletters,
 2,3 and 4.

## Automatic Control

The theory and application of automatic control techniques including feedback control systems.

Mr. Robert B. Wilcox. Chairman,
Raytheon Mfg. Co., 143 Callfornia
St., Newton 58, Mass.
Fee $\$ 2$.

## MEMBERSHIP

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## Electron Devices

Electron devices, including particularly. electron tubes and solid state devices.

Dr. John S. Saby, Chairman, Elec-
tronics Laboratory, General Electric Company, Syracuse, New York Fee \$2. 9 Transactions, 3 Newsletters, 2 Technical Bulletins. ${ }^{1} 1,{ }^{2} 2,{ }^{4}$, "Vol, ED-1, Nos. 1-4; ED-2, No. 1.

## -IRE's 23 Professional Groups

the group chairman, and publicafions to date.

| Electronic Computers <br> Design and operation of electronic computers. <br> Mr. Jean H. Felker, Chairman, Bell Telephone Laboratories, Whippany, N.J. <br> Fee $\$ 2.13$ Transactions, 5 Newsletters. *Vol. EC-2, Nos. 2-4; *Voi. EC-3, Nos. 1-4; EC-4, No. 1 . | Engineering Management <br> Engineering management and administration as applied to technical, industrial and educational activities in the field of electronics. <br> Mr. C. J. Breitwieser, Chairman, Lear, Inc., 3171 S. Bundy Drive, Los Angeles 34, Calif. <br> $\underset{* 1,{ }^{*} \text { 2. }}{\text { Fi. } 2}$ Transactions, 8 Newsletters. | Industrial Electronics <br> Electronics pertaining to control, treatment and measurement, specifically in industrial processes. <br> Mr. George P. Bosomworth, Chalrman, Firestone Tire and Rubber Co., Akron 17, Ohio. <br> Fee \$2. 2 Transactions, "PGIE-1-2. |
| :---: | :---: | :---: |
| Information Theory <br> Information theory and its application in radio circuitry and systems. <br> Mr. Louls A. DeRosa, Chalrman, Federal Telecommunications Lab., N.J.' <br> Fee \$2. 4 Transactions, 1 Newsletter. "2, *3, 4. | Instrumentation <br> Measurements and instrumentation utilizing electronic techniques. <br> Mr. Robert L. Sink, Chairman, Cone solidated Engineering Corp., 300 N . Sierra Madre Villa, Pasadena, Calif. <br> Fee \$1. 3 Transactions. "2, "3. | Medical Electronics <br> The application of electronics engincering to the problens of the medical profession. <br> Dr. Julia F. Herrick, Chairman, <br> Mayo Foundation, Rochester, Minn. <br> Fee \$1. 1 Transaction. 3 Newsletters. "1. |
| Microwave Theory and Techniques <br> Microzvave theory, microzeave circuitry and techniques, microwave measurements and the generation and amplification of microzaves. <br> Mr. Alfred C. Beck, Chairman, Bell <br> Telephone Laboratories, 463 West <br> Street, New York 14, N. Y. <br> Fee \$2. 6 Transactions. ${ }^{\text {V }}$ ol. MTT-1, No. 2; <br> -Vol. MTT-2, Nas. 1-3; MTT-3, No. 1. | Nuclear Science <br> Application of electronic techniques and devices to the nuclear field. <br> Dr. Donald H. Loughridge, Chairman, Northwestern Tech. Inst., Evanston, Ill. <br> Fee \$2. 1 Transaction, 3 Newsletters. | Production Techniques <br> New advances and materials applications for the improvement of production techniques, including automation techniques. <br> Mr. R. R. Batcher, Chairman, 240-02 <br> 42nd Ave., Douglaston, L.I., N.Y. <br> Fee $\$ 1$. |
| Reliability and Quality Control <br> Techniques of determining and controlling the quality of electronic parts and equipment during their manufacture. <br> Mr. Leon Bass, Chairman, Jet En- <br> gine Department, General Electric <br> Co., Cincinnati 15, Ohfo <br> Fee \$2. 4 Transactions, 1 Newsletter. "1, "2, *3, 4. | Telemetry and Remote Control <br> The control of devices and the measurement and recording of data from a remote point by radio. <br> Mr. Conrad H. Hoeppner, Chairman, Stavid Engineering, Plainfield, N.J. <br> Fee \$1. Transactions, Newsletter. 1-2. | Ultrasonics Engineering <br> Ultrasonic measurements and communications, inchuding underwater sound, ultrasonic delay lines, and various chemical and industrial ultrasonic devices. <br> Mr. Morton D. Fagen, Chairman, <br> Bell Telephone Laboratories, Whippany, N.J. <br> Fee \$2. 2 Transactions, 4 Newsletters. "1, 2. |
| Vehicular Communications <br> Commminications problems in the field of land and mobile radio services, such as public safety, public utilities, railroads, commercial and transportation, stc. <br> Mr. W. A. Shipman, Chairman, Columbia Gas Systems Service Corp., <br> 120 East 41st St., New York 17, N.Y. <br> Fee \$2. 4 Transactions, 2 Newsletters. 2, | Miss Emily Sirjane <br> IRE-1 East 79th St., New York 21, <br> Please enroll me for these IRE Pro <br> Name <br> Address <br> Place <br> Please enclose remittance with thi | COUPON PG.7.55 <br> ional Groups $\qquad$ $\qquad$ <br> der. |

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\section*{capacitance a attenuation <br> | TYPE | NNFF | IMPED $\Omega$ | O.D. |
| :--- | :---: | :---: | :---: |
| C 1 | 7.3 | 150 | $.36^{\circ}$ |
| C 11 | 6.3 | .173 | $.36^{\circ}$ |
| C 2 | 6.3 | 171 | $.44^{\circ}$ |
| C 22 | 5.5 | 184 | $.44^{\circ}$ |
| C 3 | 5.4 | 197 | $.64^{\prime}$ |
| C 33 | 4.8 | 220 | $.64^{\circ}$ |
| C 4 | 4.6 | 229 | $1.03^{\circ}$ |
| C 44 | 4.1 | 252 | $1.03^{\prime}$ |}


(Continucd from page 101.A)
Los Angeles Chapter-January 10
"Manufacturing and Application Techniques of "Transistors" by Don Combes and Leslie King, Hydro-Aire, Inc.

Philadelphia Chapter- A pril 20
"Significant Testing of Super-Reliable Components" by John A. Connor and Richard H. Baker, Radio Corporation of America.

Washington Chapter-March 16
"Behavior of Ferrites in Microwave Components" by John C. Cacheris. Diamond Ordnance Fuze Laboratory.

## Eiectron Devices

Boston Chapter-March 30
"Low-Noise UHF Ceramic-Metal Triolle" by G. C. Downing ard IV. C. Wicke, both of Bomac Laboratories, Inc.

Engineering Management
Philadelphia Chapter-Octcber 27
"Interaction Between Top-Level Management and Engineers in a Large Corporation" by J. 'T. Cimorelli, RCA

## Chicago Chapter-Feloruary 18

"Management Considerations for New Product Introduction" by E. H. Wavering, Motorola, Inc.

## New York Chapter-April 21

"Supervising Engineering Programs from the Cost P'oint of liew" by F. X. Lamb, Weston Electric Instrument Corporation.

Electronic Comiuters

## Akron Chapter-April 26

"Linear Programming" by Joseph E. Flanagan, Applied Science Rep. of IBM Corporation.

## Boston Chapter—ipril 21

"Digital Machines for Nationwide Dialing" by John Meszar, Bell Telephone Laboratories.

Boston Chapter-February 24
"Panel Discussion: Requirements and Applications of Computers in Business" by Milton Brand, Nowland and Company, and Edward L. Wallace, Harvard Business School.

Chicago Chapter-February 18
"Fundamental Considerations in the Design of Magnetic Core Storage Systems" by Robert Schuman, Argonne National Laboratories.

Chicago Chapter-January 21
"Teletype High-Speed Equipment and Systems" by IV. P. Byrnes, Teletype Corporation.
(Continued on page 151A)


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## 




Section Irecines

Akron
"The Weapon Systems Concupt," by Col. E. N. Ljunggren, Air Research end Development Command; April 19. 1955.

General Electric Lighting Shrow, and Election of Officers: May 17, 1955.

## Albuquerque-Los Alamos

a Fstablishment of Keliabilities for Vacuum Tubes in Complex Electrone Derices," by R. O. Frantik, and "Quality Assurance of Fiectron Tubes for Maximum Keliability Application," by A. F. Hurford, both of Sandia Corporation; May 12. 1955.

## Atlanta

"Frequency Control of UHF Oscillators," by E. D. Holmes, Cicorgia Institute of Tech.; April 29, 1955.

Bealmont-Por: Arther
"Antenna Applications in Two Way Radio Systems," by T. J. McMillin, Communications Engineering Company; May 18, 1955.

## Bingham:on

"Development of Automation," by T. W. Zebley. Gencral Electric Corp.; May 9, 1955.

## Buffalo-Niagara

"The Evolution of Broad-Band Mixer-Duplexers." by T. N. Anderson, A•ttron, Inc.; April 20, 1955.

## Cincinnati

"Engineering Problems in the Nuclear Age" by Rear Adm. H. G. Kickover, U. S. Navy; February 24,1955
"The New WILW Cathanode Modulation System," by K. J. Kockwell, Crosley Bestg. Corp.; March 15, 1955.

Spring Technical Conference; April 15. 16, 1955.

## Clevelakd

"The Weapon System Concept." by Col. E. N. Ljungeren, ISAl; April 19. 1955.

## Davton

"Semiconductor Plyysics as Applied"to Junction Transistors and Rectifiers." by N. B. Nichols, Raytheon Mfg. Co.; April 26, 1455.

## Detroiy

"Automatic Fabrication of Eilectronic Equipment." by D. F. Melton. Feneral Mills, Inc.; April 15, 1955.

## Elmira-Corving

"Automatic Street Lighting Control," by (i.W. Nagel. Westinghonse Electric Corp.: April 18, 1955.

## El Paso

"Engineering Training," by Dr. B. W. Holcombe. Texits Western College: February 21, 1955.
"Audio Engineering," by P. W. Falipach, Klipsch and Associates; March 8, 1955.
"Intercommunication Svstems. Aural and Visual," by 11. Markowitz, Custom Electronies; April 21, 1955.

## Emporicm

"Traveling W:ave Tubes." by Jobri (raenzle, Sylvania Blectric Products; May 17. 1955.

## Fort Wav:e

"Basi- Concepts of Irtormation Theory," by Dr. B. M. Oliver, Her-lett-Packard Corp: March 31, 1955.
"Status of Traveling-Wive Tube Development," by A. K. Wing, F.T.1..; April 7. 1955.

## Houston

"Instrument Engineering." by J. V. Sigford, Minneapolis Honeywell Kegllator Co.; May 17, 1955.


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# Aircraft Radio Corporation BOONTON, NEW JERSEY 

[^65](Continued from page 106A)

HuNTSVILI.E:
"A Contribution to Microwave Measurements, by Dr. F. I. Tischer, Redstone Arsenal; April 26, 1055.

General Filectric "llouse of Magic"; May 1.3. 955.

## InYokern

"Digital Techniques Applied en Aircraft Fire Control Systems." by Robert M-Intyre. Librascope. Inc.: March 21, 1955.
"Electrical Ceramics." by W: A. East. Consult ing Fingineer; April 18, 1955.
"A System for Autonatic Target Acquisition by a Phototheodolite Litizing a kemote Tracking Kadar," by James Sherwin, II, S. Naval Drdnanme Test Station; May 9. 1955.

## Ithaca

"Pulse "ode Modulation." by $S$. Slıriner. Federal Telecommunications I.alss; May 5, 195.5.

> LIITIE ROCR
"High Fidelity Systems for Honke I se." by I. Spilman, WV. M. McClanahan and Harry Cooke; May 10, 1955.

## IGNDON

"The Application of Transistors in Comphter Ci-cuitry, " by (:, D. Florida, The Defence Restareh Board; April 26, 1955.

## Long Islani

"Xincleat Instrumentatieon," by K. L. Chase. Brookhaven National Laboratories: May 10. 1955.

## Les Angeles

"The Role of Scientific Kesterarh in the Development of Missile Systems," by Dr. Ernst Krause, Lookheed Missile Systems Division, and . 1 Review of 1'rogress in the Land Mobile Communication Service," by J. $k$. Byme. Motorola Kiverside Keseareh Iaboratory; May 3, 1955.

## Loresvilife

"The Application of IBM io Enginerering and Statistical Problems," by P. H. Sterbenz. International Business Machines Corp; Fibluary 10. 1955.
"Recent .dvances in the Keproducing Art." by . . M. Wiggins and Howad Souther, both off Eilectroviele, Inc: March 10, 1055.
"lligh Fidelity - I'ast and Fiuture, " by Marvin Camats, Armour Rescareh Voundation; . Iprit 14. 1955.

## Miami

Student grapers by the following I niversity of Miami students: S. Afagomes, R. Watts. Dave Wensley. Ray IIarpet. Harry Hoperoft and $k$. N. Stock: April 29. 1955.

## Nf.W YORK

*. Ipplication of Communication Consept.s en Infrared Problems," by Dr. M. J. E. Golary. Squier Signal Iab.; April o, 1955.
"An Experimental Automobile Reweiver Fimploying Transistors." by T. O. Stanley, and "5 Watt Transistor Amplifier," by A. I. Aronson. both of David Sarnoff Reseurch Center: May 4. 1955.

## Oliawa

"〈'urrent Problems in UHF" and Microwave Multiplex Communication Systems," By Dr. H. J Von Bateyer, Dept. of Defence Production; Election of Officers; $A_{\text {pril 14. }} 1955$.

Field Trip to RCA Vietor Television and Hone Receiver Plant and Fort Welington at I'rescott. Ontario; May 13.1955.

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(Continued from page 198.-1)

Philadelphia
"Fundamentals of Color Television," by J. W Wentworth, KC. $1:$ Match 1, 1955.
"Color Signal Generating Fquipment," by Dick O'Brien, CBS; March 8, 1955
"Color Reproducang Tubes and Associated Components," by B. Loughlin, lazeltine Corp. March 15, 1955
"Color Decoding Circuits," by Jack Avins, RC. Industry Service Lab.; March 22. 1955
"Measurement and Lqualization of Amplifiers and Transmission Systems for Color TV Service," by Ilugh Kelly. Bell Labbotatories; March 29, 1955
"Colorimetry Problems in Color TV' and the Effect of Transmission Errors on Color Keproduc tion." by IIarold Weiss General Electric Company; April 5, 1955.
"Automatic Electronic Production," by Dr J. I. Ryder, President. IKl:; May 4, 1055.

## Phoenix

"Series Peaked Amplifier Analyzed on Analogue Computer," by Fired Srhwepn, Student, Inversity of Arizona. and Tapeseript "The I'hysies of Music and Hearing, " by W. A. Koch. Bell Telephone Labs.; April 8, 1955

## PORTLAND

"Engineering Asp+cts of U11F Booster Instal lations," and "Engineers are People," both by Dt G. II. Brown. David Sarnoff Research Center; April 21, 1955.
(Continuea on fage 112.1)

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(Continted from page 111A)

Student Papers: "A Portable Sodium-Flash Recorder," by M. H. MacKenzie, "Elementary Magnetic Digital Computer Component," by R. W. Austad, and "Air-Borne Measuren,ents of I.ow Frequency Effective Ground Corductivity in Alaska." by G. M. Stanley; April 30, 1955.

## Rochester

"Modern Loudspeaker Design," April 21, 1955

## Rome-Utica

"More Significant Characteristics of Nonlinear Circuits," by Dr. Ernst Weber. Polytechnic Institute of Brooklyn; May 4, 1955.

## Sacramento

"Planning and Constructing a Television Station," by P. K. Onnigian; KBET-TV; lilection of Officers; May 13, 1055.

## Salt Iake C:ty

Demonstration of Boeing Analog Computer, by Dr. R. E. Stephenson, University of Utah; Spril 13, 1955.

General meeting; May 17. 1055.

## San Diego

"Color Broadcasting Specifications and Applications," by George Jacobs, Wrather-Alvarez; April 2, 1955.
"Bio-Technology in Engineering." by Dr. A. M. Small, Navy Electronics Lab.; May 4. 1955.

## Toronto

"Thirty Years in Canadian Radio." by Miss Jane Gray. Commentator, CHML; April 18. 1955.

## Tulsa

"Dynatnic Instrumentation," by C. M. Hathaway. Hathaway Instrument Co.; April 21. 1955.

## Twin Cities

"Feedback Control Systems-Past. Present and Future," by Prof. T. J. Higgins, University of Wisconsin; April 26, 1955.

## Washington

"The Navy 'Jim Creek' Transmitting Station," by Cdr. R. G. Bywater, USN and II. E. Dinger, Naval Research Lab.; Mey 9 9, 1955.

## Williamsport

"Guided Missiles are Stmarter than People." by D. E. Mullen, General Electric; March 11. 1955.
"Antomation," by Ben Warriner, General Electric Lab.; April 21, 1955.
"Traveling Wave Tubes," by Dr. R. C. Hutter. Sylvania Electric; May 17, 1955.

## SUBSECTIONS

## Buenayentura

"Early Guided Missile Development in Germany," by Dr. Willy Fiedler. L.. S. Naval Air Missile Test Center; April 14, 1955

## Centre County

"Automation." by Charles Godwin, Cornell University; election of officers; April 21, 1959.

## Charleston

Tour of Television Station and Radio Studio w.CSC conducted by Dick Hart, Station Engineer; April 26. 1955.
East bay
"Radio Astronomy and Its Engineering Aspects, ${ }^{n}$ by Dr. K. N. Bracewell, University of California; May 3, 1955,

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## Section Ilectings

(Continued from page 112A)

## Erie

"The Theory and Application of Quartz Ulitrasonic Delay Lines, ${ }^{"}$ by J. M. Wolfskill and "The Theory and Application of Quartz Crystal Units." by R. H. Tuznik, both of Bliley Electric Company; April 28, 1955.

## Lancaster

"An Excursion in Flectronics." by C. N. Hoyler, RCA Laboratory, May 11, 1955.

Mid-Hudson
"Magnetic Core Circuitry," by Dr. An Wang. Wang Labs; April 19, 1055.

Monmouth
"Ferrites at Microwave Frequencies," by A. C. Fox, Bell Telephone Labs.; May 18, 1955.

Oranger Belt
"Stereophonic (3D) Sound in Your Home," by C. M. Brainard, Master Electronics Company; April 25, 1955.

## Palo Alto

"Pulse Applications of Junction Transistors," by John Linvill, Stanford University; April 21, 1955.

## Wichita

"Design of Airborne Automatic Antenna Tuners," by L. Hutton, Boeing Airplane Co.; April 27. 1955.

## Student Branch Meetings

University of Akron (IrE-AiEE Branch)
"Engineering Ethics" by R. D. Landon, Dean. College of Engineering. $\mathrm{U}^{\top}$, of Akron, and Election of Officers; April 25, 1955.

Alabama Polytechnic Institute (IRE Branch)
"Electronics in Guided Missles" by S. Johnson; April 25, 1955

University of British Columbia (IRE-AIEE Brasch)
"Some Aspects and Future Trends of Electrical Engineering" by D. Carpenter, President, Research Industries, Ltd., and Election of Officers; March 30. 1955.

Polytechnic Institl'te of Brookiyn (IRE-AIEE Branch) Eive. Div.
"The New High Speed Electronic Printer" by Edmund Diginlio, Field Engrg. Mgr., Control Instrument Co., Brooklyn. N. Y.., and Busiriess Meeting; April 28, 1955.

Film, "A is for Atom" and Election of Officers: May 11, 1955.

Brown University (! RE-AIEE Branch)
Election of Officers; March 16, 1955.
"An Inexpensive Voltage Kegulator," by James Davis, Indergraduate, Brown U.; March 30, 1955.

Clarkson College of Technology (IRE-A1EE Branch)
Election of Officers: April 21, 1955
"Motors and Motor Control," by D. B. Seymore. Westinghouse; April 28, 1955.
"I.B.M. Computers \& Acrounting Machines," by Mr. Markle, I.B.M. and Film. "Piercing the Unknown." and Business Meeting: May 12, 1955,
(Continued on page 117A)

## BALLANTINE Sensitive,

## Wide Band Electronic Voltmeter

measures 1 millivolt to 1000 volts from 15 cycles to 6 megacycles

## Accuracy $3 \%$ to $3 \mathrm{mc} ; \mathbf{5}$ above <br> Input impedance 7.5 mmfds shunted by $\mathbb{1} 1$ megs



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It features an output impedance of 100 ohms ct a terminal box at end of $3^{3}$-cable: frequency range of 10 cps to 1 mc continu. ously variable over decade steps; rise time of 0.02 microseconds at the low impedance output.

Write for catalog



UNI'ERSIty of Colorado (IR F.-AIFE Branch)
"The Resistor-A Simple Flement?" by Harry Bishop, Hytronic Measurement Assoc.; May 11, 1955.

Coll'mbia L'Niversity (IRE-AlEE Branch)
"Magnetic Recording" and "Method for Time or Frequency Compression-Expansion of Speech" (P'(iA Tape Scripts); April 18-19. 1955.

Film, "Research and Development at Hughes Aircraft Co.": May 3. 1955.

I niversity of Connectice-t (IRE-AIEE Branch)
"X-Ray and High Power Flectron Tubes," by Chester A. Kirka, Machlett Labs. and Election of Officers: March 31, 1955.
"Instrument Serves \& Their Application to Aircraft Auto-Pilots," ly F. W. Campbell. Sperty Gyroscope; April 28, 19.55.
(niversity of Delaware (IRE-AIfE Branch)
"Africa," by Professor Eiarl Parker Hanson; March 7, 1955.
"The Gernanium Story." by Dr. S. M. Christian, "A P-N.P Alloy Junction Transistor for Radio Frequency Amplification," by Dr. C. W. Mueller and "An Experimental Transistor Personal Broadrast Receiver," by I.. F.. Barton; April 19, 1955.

CNidersity of Detroit (IRE-ilele Branch)
Student Paper Contest: "An Introduction to the I'se of Symbolic Logic in the Design of Switching Circuits," by J. Dennis Kennedy and "Electronic Comparator," by Albert Vanschaemelhout; April 20, 1955.
"Success Story" (tape and slide talk) and "An Introduction to Miniature Low Pass Filters for Telemetering in Guided Missile Research, ${ }^{\text { }}$ by Victor Schutzwhol; May 5, 1955.

Drexel Institlte of Technology (1RE-AlEE BRANCH)
"Carbon in Electrical Engineering." by Bernard Silver, Flectronite Carbon Co.: April 5, 1955.
"Microwaves," by Richard A. Dibos, Philco Corp-; April 28, 1955.

University of Fiorida (IRE-AIEF Branch)
"The Trends in Industrial Electrical Distribu. tion Systems," by C. F. Kucera, Allis-Chalmers, and Business Meeting; February 14, 1955.
"Research \& Development of Transformers" (Film); March 14, 1955.
"Digital Computers," by Mr. Zyrak and Mr. Rich, both of Lincoln L,ab., M.I.T.; March 28. 1955.

Election of Officers; May 9. 1955.
Iowa State College (IRE-AIEE Branch)
"Hints in Interviewing for Jobs," by J. J. Jondle, Student and Election of Officers; April 27. 1955.
"Development of Radar," by C. J. Marshall. IRE Regional Director, Region 5; May 6, 1955.

Lafayette (ollege IIRE-AIEE Branch)
"Tests on 69 Kv Horned Air Brake Switches," by Mr. Wesley Smith and Mr, Jack McDonald, Pennsylvania Power and Light Company; March 10. 1955
"Electronics in Medicine," by Ray Wiech. Lafayette Student '56 and Election of Officers; April 14, 1955.
"High Speed Electronic Flash Photography," by Richard B. Manbicki, Student, Class '56 and Ceneral Meeting; May 121955.
(Continued on page 118.A)

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## HOW TO ZERO-SET RANDOM PHASE VARIATIONS

Many modern control devices are designed for applications where sensed input signals fluctuate randomly about an approximately known frequency, In some of these applications, the information is conveyed by the phase relationship within one cycle, and the random cycle-tocycle phase variations often submerge the signal in noise. Filtering, or averaging, techniques may be extremely difficult to devise because of the requirement for use within one cycle.

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In this mamer, the resolver is reset to a prescribed phase relative to the signal at a fixed point of every cycle of the generated signal.

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3. Decreased overtravel. Pre-tension of contact springs against the support springs provides proper contact pressure with minimum travel after closure. Thickening of confact springs is not a proper alfernative as this tends to increase the frequency of contact chatter and bounce without providing suitable frictional damping io minimize arcing and mechanical wear.
4. Less mechanical wear. Motion means wear. Support springs provide adequate contact pressure without excessive spring fexure, thus increasing relay life.
5. Increased sensitivity of adjustment. Reduced spring over. travel permits a smaller armalure air gap for greater over-all sensitivity.


The principle of contact support springs is one of many exclusive features contributing to the use-tested superiority of performance in North Relays of the type shown above. Some specific control type switching applications are safety controls - switchboards - elevator controls - power control circuits - carrier application - intercommunication systems - fire alarms - airport lighting centrols and computers.

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| 1 c.p.s. $-5.0 \mathrm{Mc} . \pm 3 \mathrm{db}$. useful gain beyond 10 Mc . | Output Voltage | with virtually no overshoot |
| Gain | greater than 15 volis r.m.s. per stage. | or ring even with severe overload; accepts positive |
| stage at mid-frequencies. | Equivalent Nois | Inpui Impedance |
| Output Impedance | Input | 1.0 megohm in parallel with |
| less than 200 ohms in series | $30 \mu \mathrm{~V}$ of grid. | 8 unf. |

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| SYSTEMS | MOST ADVANCED | NEW PRDGRAM | IN | IN |
| ENGNEERING | ELECTRONIC DATA | IN | IN | AVIATION | ELECTRON TUBE

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\begin{tabular}{|c|c|c|c|c|c|c|c|c|c|c|c|c|}
\hline \multirow{3}{*}{FIELDS OF ENGINEERING ACTIVITY} \& \multicolumn{12}{|c|}{TYPE OF DEGREE AND YEARS OF EXPRRIENCE PREFERREO} \\
\hline \& \multicolumn{3}{|c|}{Electrical Engizeers} \& \multicolumn{3}{|l|}{Mechanacal Engineers} \& \multicolumn{3}{|c|}{Physical Scuace} \& \multicolumn{3}{|l|}{\begin{tabular}{c} 
Chemistry \\
Ceramics \\
6iass Technology \\
Metallurgy \\
\hline 1.2 2.3
\end{tabular}} \\
\hline \& 1.2 \& 2.3 \& \(4+\) \& 1.2 \& 2.3 \& \(4+\) \& 1.2 \& 2.3 \& \(4+\) \& 1.2 \& 2.3 \& \(4+\) \\
\hline \begin{tabular}{l}
SYSTEMS \\
(Integration of theory, equipments, and environment to create and optimize major clectronic concepts.)
\end{tabular} \& \& \& \& \& \& \& \& \& \& \& \& \\
\hline AIRBORNE FIGE CONTROL \& \& \& w \& \& \& \& \& \& w \& \& \& \\
\hline digital data handling devices \& \& \& C \& \& \& C \& \& \& C \& \& \& \\
\hline missile and radar \& \& \& M \& \& \& M \& \& \& M \& \& \& \\
\hline InERTIAL NAVIGATION \& \& \& M \& \& \& M \& \& \& M \& \& \& \\
\hline COMMUNICATIONS \& \& \& \[
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\hline \begin{tabular}{l}
DESIGN • DEYELOPMENT \\
COIOR TV TUBES-Electron Optics-Instrumental Analysis \\
Solid States (Phosphors, High Temperature Phenomena, Photo Sensitive Materials and Glass to Metal Sealing)
\end{tabular} \& L \& L \& L \& \(i\) \& L \& L \& L \& \(t\) \& L \& L \& L \& L \\
\hline receiving tubes-Circuitry-Life Test and Rating-Tube Testing-Thermionic Emission \& H \& H \& H \& \& H \& H \& \& H \& H \& \& H \& H \\
\hline SEmI-Conductors-Transistors Semi-Conductor Devices \& H \& H \& H \& \& \& \& H \& H \& H \& \& \& \\
\hline microwave tubes-Tube Development and Manufacture (Traveling Wave-Backward Wave) \& \& H \& H \& \& H \& H \& \& H \& H \& \& H \& H \\
\hline gas, power and photo tubes-Photo Sensitive DevicesGlass to Metal Sealing \& L \& L \& L \& L \& L \& 1 \& L \& L \& L \& L \& 1 \& 1 \\
\hline \begin{tabular}{l}
aviation electronics-Radar-Computers Servo Mech-anisms-Shock and Vibration-Circuitry-Remote Control -Heat Transfer-Sub-Miniaturization-Automatic Flight \\
-Design for Automation-Transistorization
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\hline radar-Circuitry-Antenna Design-Servo Systems-Gear Trains-Intricate Mechanisms-Fire Control \& \(x\) \& \[
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\hline | compurers-Systems-Advanced Development-Circuitry |
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| -.Assembly Design-Mechanisms-Programmiag | \& C \& \[

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\hline communications-Microwave-Aviation-Specialized M litary Systems \& \& F \& M
C
F \& \& F \& C ${ }_{\text {M }}$ \& \& F \& M
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f \& \& \& <br>
\hline RADIO SYSTEMS—HF-VHF-Microwave-Propagation Analysis-Telephone, Telegraph Terminal Equipment \& \& 0 \& \% \& \& 0 \& \% \& \& 0 \& F \& \& \& <br>
\hline misile guidance Systems Planning and Design-Radar -Fire Control-Shock Problems-Servo Mechanisms \& \& $F$ \& M \& \& f \& M \& \& F \& $\stackrel{M}{\text { F }}$ \& \& \& <br>
\hline COMPONENTS-Transformers-Coils-TV Defection Yokes (Color or Monochrome)-Resistors \& \& C \& C \& \& c \& C \& \& C \& C \& \& \& <br>
\hline Mech. and Elec.-Automatic or Semi-Automatic Machines \& \& H \& H \& \& H \& H \& \& H \& H \& \& \& <br>
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\end{tabular}

|  | C-Camden, K. J.-i? Greate: Philade phia riear many suburban |
| :---: | :--- |
| Lommunities. |  |

[^66]Please send resume of education and experience, with location preferred, to:

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（Comtinted on Fuge 126．1）

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(Continued from pase 12+A)

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# Systems Development and <br> <br> The Ramo-Wooldridge Corporation 

 <br> <br> The Ramo-Wooldridge Corporation}

The Ramo-Wooldridge Corporation (except for the specialized activities of our subsidiary, Pacific Semiconductors, Incorporated) is engaged primarily in developing - and will soon start to manufacture-systems rather than components. For military customers our weapons systems responsibilities are in the fields of guided missiles, fire control, communications, and computers. Our non-military systems activities are in the general area of automation and data-processing.

Emphasis on systems development has consequences that profoundly affect all aspects of an organization. First, it demands an unusual variety of scientific and engineering talent. A single systems development project often requires concurrent solutions of challenging problems in the fields of electronics, aerodynamics, propulsion, random phenomena, structures, and analytic mechanics. In addition, the purely technical aspects of a systems problem are ofter associated with equally important nontechnical problems of operational, tactical, or human relations character.

Therefore, competent systems development requires that a company contain an unusually large proportion of mature, experienced scientists and engineers who have
a wide range of technical understanding and an unusual breadth of judgment. Further, all aspects of company operations must be designed so as to maximize the effectiveness of these key men, not only in the conduct of development work but in the choice of projects as well.

At Ramo-Wooldridge we are engaged in building such a company. Today our staff of professional scientists and engineers comprises $40 \%$ of the entire organization. Of these men, $40 \%$ possess Ph.D. degrees and another $30 \%$ possess M.S. degrees. The average experience of this group, past the B.S. degree, is more than eleven years.

We believe the continuing rapid growth of our professional staff is due, in part, to the desire of scientists and engineers to associate with a large group of their contemporaries possessing a wide variety of specialties and backgrounds. It is also an indication that such professional men feel that the Ramo-Wooldridge approach to systems development is an appropriate one.

We plan to continue to maintain the environmental and organizational conditions that scientists and engineers find conducive to effective systems development. It is on these factors that we base our expectation of considerable further company growth.


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Our income and benefit advantages scored high on this test, too. Finally, there were many "extras," like the Westinghouse Patent Award Program, that make investigation of the current openings worthwhile for all electronic engineers.

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## (Continued from page 126.A)

PHYSICIST OR ELECTRONICS ENGINEER
Physicist or electronics engineer to design, construct and install set-ups to obtain lata on engine ignition, periormance. Diversfied projects might require mechanical or electronic instrumentation, also design of auxiliary control circuits. Electric Auto-lite Company, Toledo 1, Ohio.

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(Cominted on page 130.1)

## ENGINEERS

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*

1. Professional Engineers Conference Board for Industry survey, "How to Train Engineers in Industry"
2. Universlty of Chicago Survey of Employee Attitudes
3. National Society of Professional Engineers,
"A Professional Look at Engineers in Industry"
 ELECTRONICS
CORPORATION
OF AMERICA
77 Broadway
Cambridge 42, Mass.

## (Continued from puge 128A)

## ELECTRONIC EMGINEER

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(Continued on page 132A)

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- Quality Control \& Test Engineers

(Continued from page 130.A)


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The Institute publishes sree of charge notices of positions wanted by I.R.E. members who are now in the Service or have received an honorable discharge. Such notices should not have more than five lines. They may be inserted only after a lapse of one month or more following a previous insertion and the maximum number of insertions is three per year. The Institute necessarily reserves the right to decline any announcement without assignment of reason.

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BSEE 1950. Experience in petroleum geophysics, broadcast, communications radio, mili. tary electronics. Interested in electronic design, development, test or field engineering. Location preference southwestern or southern area. Will carefully consider any offer. Box 827 W .
(Continued on page 136.A)


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Positions Ihamed

## By Armed Forces Veterans

(Continued from page 13tA)

## ENGINEER

BSEE 1951, T.IT. Eta Kappa Nu. Age 25. 2 years experience in adapting power system analysis to digital computers. 2 years experience at White Sands in programming and in trouble shooting a large scale digital computer. Desires position involving logical design of computers. Box 838 W .

## ELECTRONIC ENGINEER

IBEE 1953. Age 25. 5 years electronic experience, including 2 years shop and testing, 1 year assistant project engineer on classified Navy project. 2 years as radio officer USAF. Desires responsible position in production or development. Box 839 W:

## INSTRUCTOR

BSEE, MSEE. Age 29, married, 1 child. Craduate work in advanced electron tube circuits, network analysis and synthesis, and feedback systems. Mathematics minor. Licensed radio annateur. 1 year communications (R.F.) design and development; 1 year analogue computer circuit research and design; 3 years applied transistor research. Excellent references. Desires a full time teaching position in an institution that has an E.E. graduate school (Ph.D.) with privilege of engaging in six semester hours per semester of graduate study. Available Sestember 1, 1955. Box 840 W .

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## ELECTRICAL ENGINEER

BS in wath. 1950, 1BEE 1954. IRE, AIEE, EIT (Ohio). Age 29, married, 1 child. 1 year experience with resistor components. Desires position in power or electronics field. Location Florida. Box 843 W .

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(Continued on page 139A)

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> Personnel Director UNIVERSITY OF DENVER
> Denver 10, Colorado

## By Armed Forces Veterans

(Continued from page 136A)

## SENIOR ELECTRONIC DEVELOPMENT ENGINEER

BS 1950. Age 25, married. 3 years video, pulse and ultrasonic systems development, project level, (patents) some guided missile system development, 2 years Army electronics instructor. Desires responsible R \& D position. Box 854 W.

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BSEE January 1951. Age 27, married, 2 children. $31 / 2$ years experience in circuit development, instrument and communications systems planning, and ordnance testing. 1 year administrative and supervisory experience. Desires position offering responsibility and advancement in San Francisco area. Box 855 W.

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(Continued on page 141A)

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120

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(Continused from page 139A)
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Virginia Polyiechnic Institute (IRE-AIEE Branch)
Film on the Manufacturing of Paper and General Meeting; February 22, 1955.

Demonstration of Microwave Communications by representatives of Bivins and Caldwell; March 1, 1955.

Election of Officers and General Meeting; March 29. 1955

Washington University (IRE-AIEE Branch)
Election of Officerz and Business Meeting: March 31, 1955, and 'A Thickness Gage I'sing Radioactive Isotopes," by Bush and Wallscheidt and "Analog Computer" by Rojko Also, March 31, 1955.
Uninersity or Washincton (IRE-AIEE Branch)
"Civil Service Employment Opportunities," by Mr. Wihrow, Seattle City Light; April 6, 1955.

Wayne University (IRE-AIEE Branch)
"Engineering and Management," by Ray Plourde, Detroit Edison Co.; April 28, 1955.

University of Wisconsin (IRE-AIEE Branch)
"Patent Law," by John Leib, Patent Attorney for Allis-Chalmers Co.; April 21, 1955.
"The Network Anelyzer," by James Skiles, Instructor at U. of Wisconsin and Election of Officers; May 9. 1955.

Worcester Polytechnic Institute (Ire-Aiee Branch)
"Hi Fi Demonstration and Lecture," by D. R. Von Recklinghausen and E. G. Dyett, Jr., of H. H. Scott Co.; April 12, 1955.
University of Wyoming (IRE-AIEE Brancit)
"Success Story," by Robert L. Hudelson, Student Member and General Meeting; April 26, 1955.

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Engineering Employment Manager
(Contimasd from prose 20A)


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(Continucd on page 147A)
(Continued on page 147A)
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MR. JAMES STARK

Electronics Park
Syracuse, N.Y.
(Continued from page 145.1)
The new device, known as a Qui-klip, was developed in conjunction with the Tinnerman Products Corporation of Cleveland, manufacturers of speed-type fasteners. It does not require special sockets for mounting, only needing, according to Kadio Receptor, two round holes to be snapped into place. Irı addition, solderless connecters are available for making electrical contact to the rectifier.

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Internationally known research organization seeks engineers and mathematicians for challenging research and development programs in the following fields:
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(Continued from page 1+7.1)

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(Continued on prage 151A)

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(Continued from page 104A)
Los Angeles Chapter-Feloruary 17
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(Continued on puge 15:A)

## Electronics Engineer

(MAGNETIC AMPLIFIER EXPERIENCE)

## THE APPLIED PHYSICS

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Los Angeles Chapter-January 20
"The Bendix Combination (;.P. and Differential Analyze:" by Max l'alevsky, Bendix Computer IDivision.
"Transistor Flip-Flops" by Chris Wanlass, Ramo-W'ooldridge Corporation.

New York Chapter-February It
"The Naval Orlnance Research (alculator (DORC)" by Byron L. Haters, Engincer in Charge of Di)RC Project.

- New York Chapter-January 20
"Bi-Lateral Mannetic Selection Systems for Large Scale Computer Memory" by Amir H. Septhban, Monrobot lah, oratories.
"High Speed Core Memory" b. E. J. Otis, Air Force Cambridge Research Station.

New York Chapter-I)ecember 16,
"Feasibility of an All Magnetic Computer" by Isaac Auerbach, Burroughs Research Center.

New York Chapter-November 2.3
"The Logical l'riaciples Employed in Underwood Electronic Computers" by Evelyn Berezin, Unlerwood Corporation.
"The Type of Circuitry Employed in Inderwood Electronic Computers" by Albert Auerbach, I'nderwood Corporation.
New York Chapter-October 26 and 27
"Project Cyclone" by Leo Batuer, Reeves Instrument Company.

## Information Theiory

Albuquerque-Ios Alamos ChapterMarch 9
"An Example in Statistical Communication Theory" by Walter E. Brown.

## Instrumentation

Houston Chapter-March 22
"A Multi-Loop Self Balancing Power Amplifier" by J. Ross MacDonald, Texas Instruments, Incorporated.

## Medical Elefctronics

San Francisco Chapter-April 7
"Electronic Instrumentation in Surgery," by Bertram Feinstein, Mt. Zion Hospital.
"The New Operating Room at Mt. Zion Hospital" by Mr. Carter Collins, Consultant, Research and Development Laboratories, U. C. Medical Center.

Microwave Theory and
Techniques
Northern New Jersey Chapter-
Jebruary 16
"Microwave Applications of Ferrites" by J. 1I. Rowen, Beli Tolephone Laboratories.
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Northern New Jersey ChapterJanuary 19
". 1 Display of X -Rand Impedance on an Oscilloscope" by Henry L. Bachman, Wheeler Labortories.

## Nuclear Science

Connecticut Valley Chapter-
February 23
The following three films were presented: "Bikini-Radiological Laboratory," "Nuclear Reactors for Research," "Operation Ivy."

## Chicago Chapter-February 18

"Trends in Reactor Development" by I.loyd V. Berkner, Associated U'niversities, Incorporated.

Chicago Chapter-December 17
"Radio Carbon Dating" by Jaines R. Arnold, Institute of Nuclear Studies, ITniversity of Chicago.
Washington, D. C. Chapter-March 25
"An Accounting of the Benefits of Nuclear Energy" by Clifford K. Beck, North Carolina State College.

## Reliability \& Quality Control

Washington, D. C. Chapter-April 1.3
"Statistical Methods for Engineers" by Leon Bass, Quality Control Div., Jet Engine Div., General Electric Company.
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"The Design of Receivers for Split Channel Operation" by R. I.. Cassellerry, General Electric Company.

Houston Chapter-.ipril 26
"Progress in Mobile Communications Equipment Design" by J. A. McCormick, Mobile Communications Equipment, General Electric Company, Syracuse.

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[^15]:    * Original manuscript received by the IRL, January 14, 1955; revised manuscript received, March 1, 1955.
    + Wheeler l aboratories, Inc., Great Neck, N. Y.
    $\ddagger$ This theorem was discovered by the author shortly after publication of his 1942 paper on the skin effect (Bibliographical reference 2). He has presented it at various meetings, including a seminar at New York Unive-sity, New York, N. Y., on March 29, 1950, and a staff neeting at Wheeler I aboratories on December 12, 1951.

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    1729
    A 1-100- $\Omega$ Build-Up Resistor for the Calibration of Standard Resistors-B. V. Hamon. (Jour. Sci. Instr., vol. 31, pp. 450-453; December, 1954.) A build-up resistor circuit developed at the Australian National Standards Laboratory comprises $10-\Omega$ manganin resistors in series-parallel. The estimated accuracy when used as a ratio device is of the order of 1 part in $10^{8}$. Test results are tabulated and the construction is illustrated.

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[^66]:    M-Moorestowr. N. J.-quiet, attractive community close to Phila.
    $\mathbf{0}$-Overseas-domestie and overseas locations.
    W-Waltham, Mass - near :he cultural center of Boston.
    $X$-Los Angeles. Callf. - west coast elactronics center.

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